

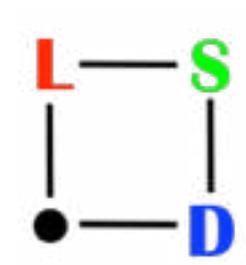
# Near Conformal Dynamics in $SU(3)$ $N_f=8$ Theory

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**(for the LSD Collaboration)**

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# Why $SU(3)$ $N_f=8$ ?

- The LSD Collaboration has been studying  $SU(3)$   $N_f=8$  for nearly a decade. Why is it so interesting?
- Early on, it was clear that this theory had near conformal dynamics due to very slowly running coupling.
- As Enrico already showed, this theory also has a light isoscalar meson:  $m_\sigma \sim F_\pi$  in chiral limit. This is a potential **composite Higgs** candidate.
- Are near conformal dynamics and the light scalar related?
- Andy will next discuss our efforts to find a good EFT for light pions and scalars.



# Lattice Strong Dynamics Collaboration



James Osborn  
Xiao-Yong Jin



Joe Kiskis



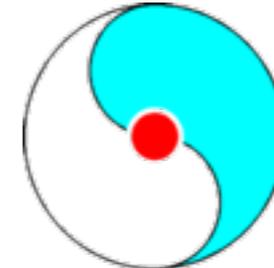
Graham Kribs



David Schaich



Anna Hasenfratz  
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Richard Brower  
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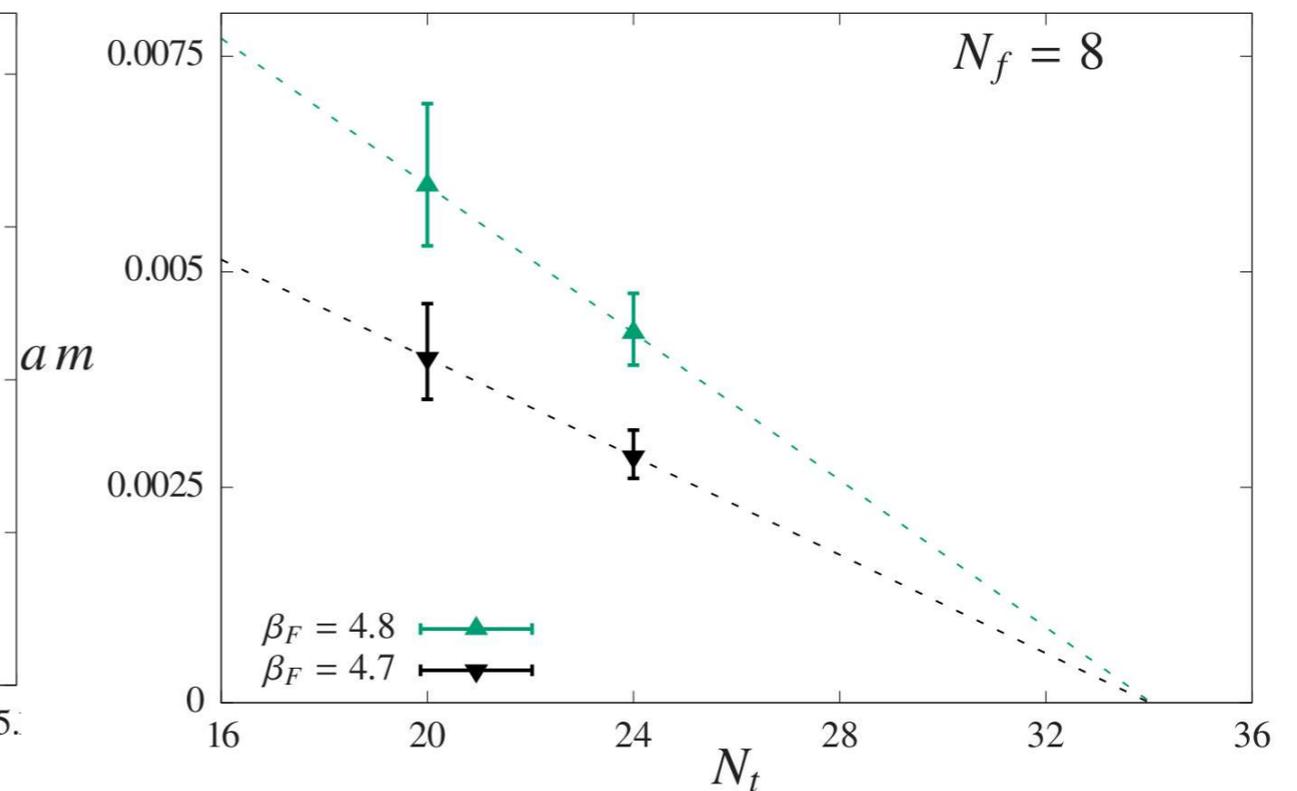
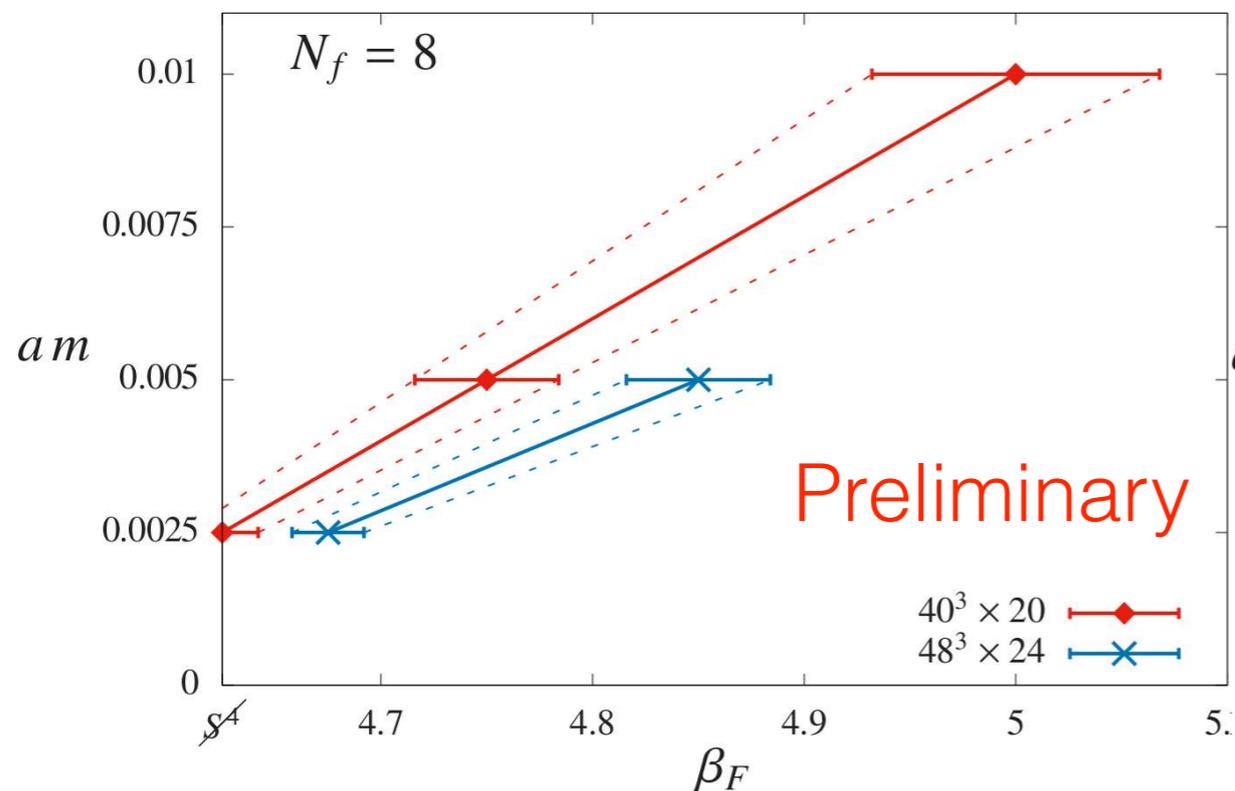
Meifeng Lin



Evan Berkowitz  
Michael Buchhoff  
Chris Schroeder  
Pavlos Vranas

# LSD SU(3) $N_f=8$ Stag

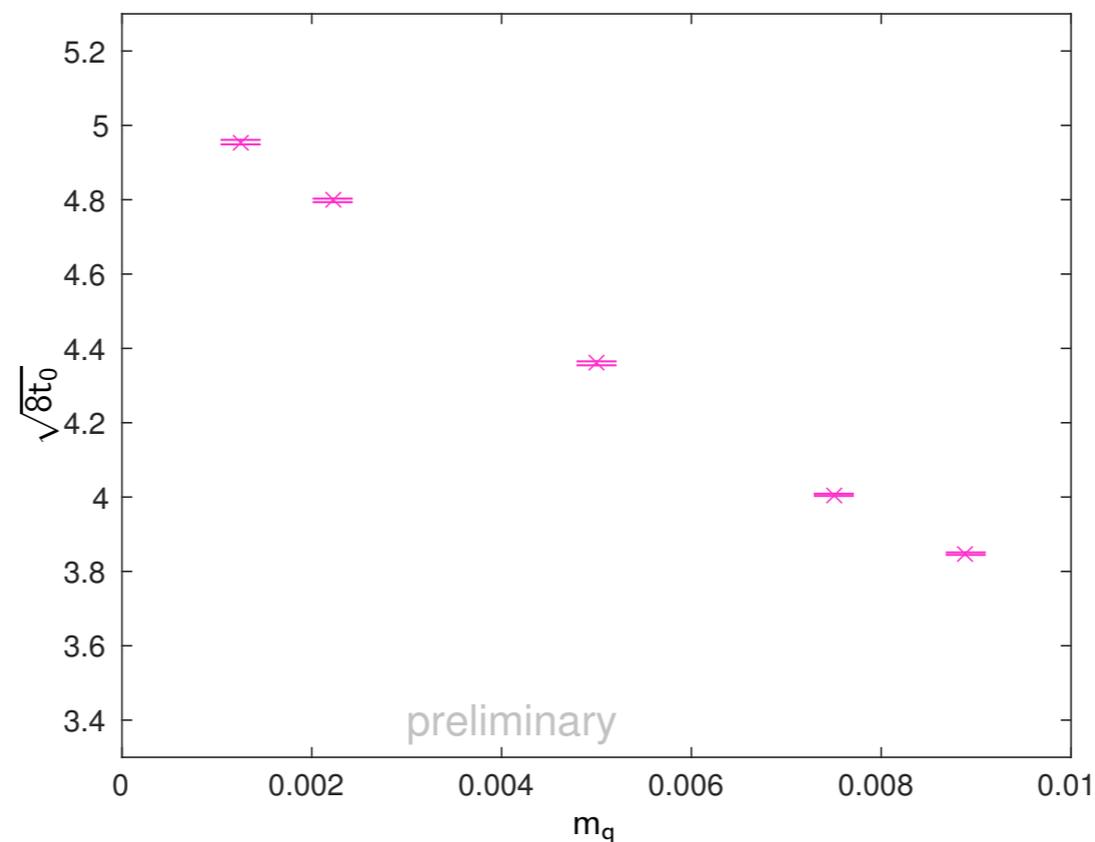
- Earlier USBSM studies (and LatKMI) used HISQ fermions which become prohibitively expensive for  $N_f=8$  on coarse lattices.
- Using nHYP stag fermions and fund+adj gauge action pioneered by Boulder group to get to somewhat coarser (but still very fine) lattices.



$T_c$  and bulk phase

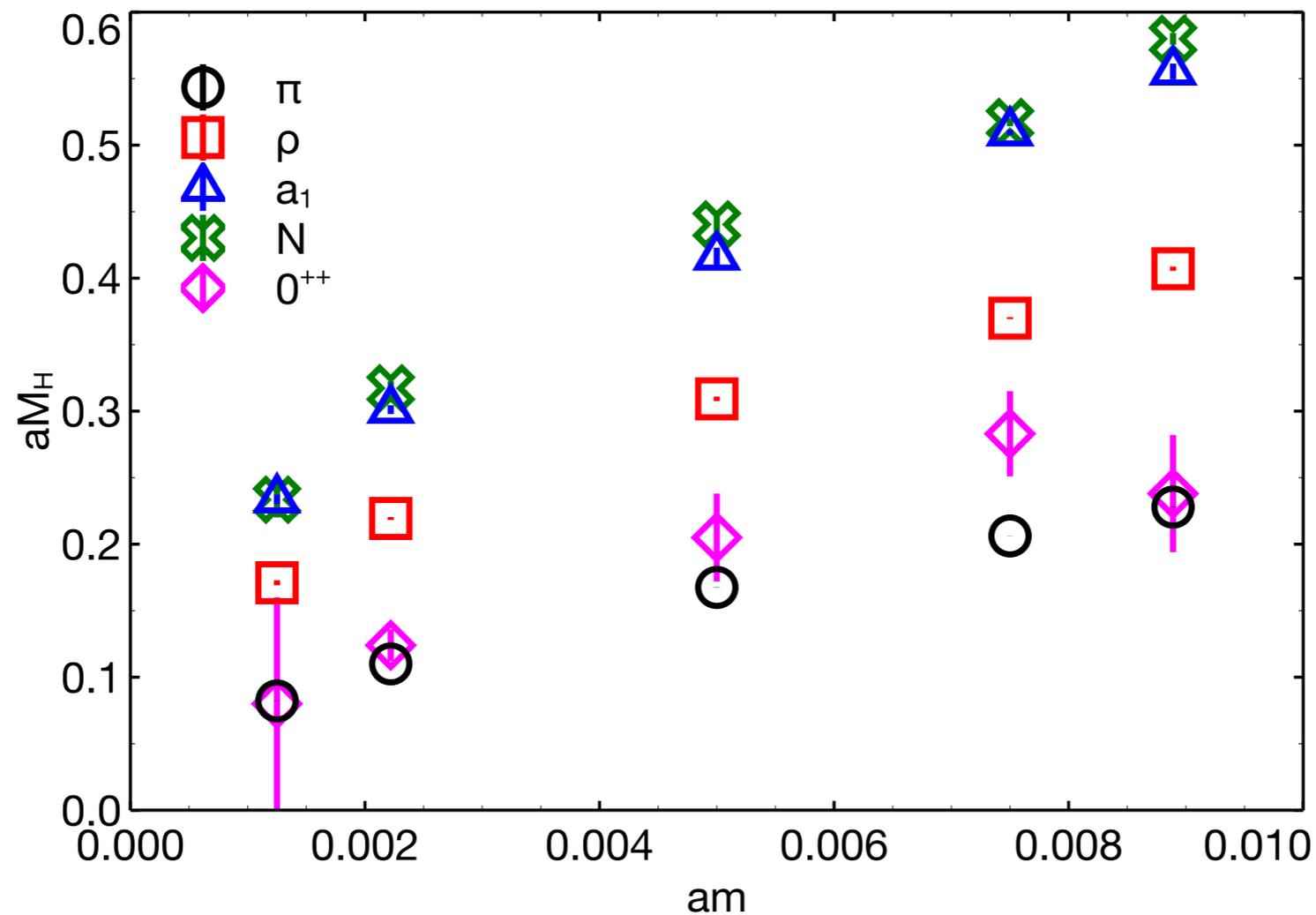
# Near Conformal Scale Setting

- In QCD, the dynamical scale  $\Lambda^{-1}$  is strongly affected by gauge coupling and weakly affected by quark mass.
- This reflects strong gluonic anti-screening and weak fermionic screening in vacuum.
- Near conformal dynamics has balanced anti-screening and screening, leading to strong mass dependence in setting the scale



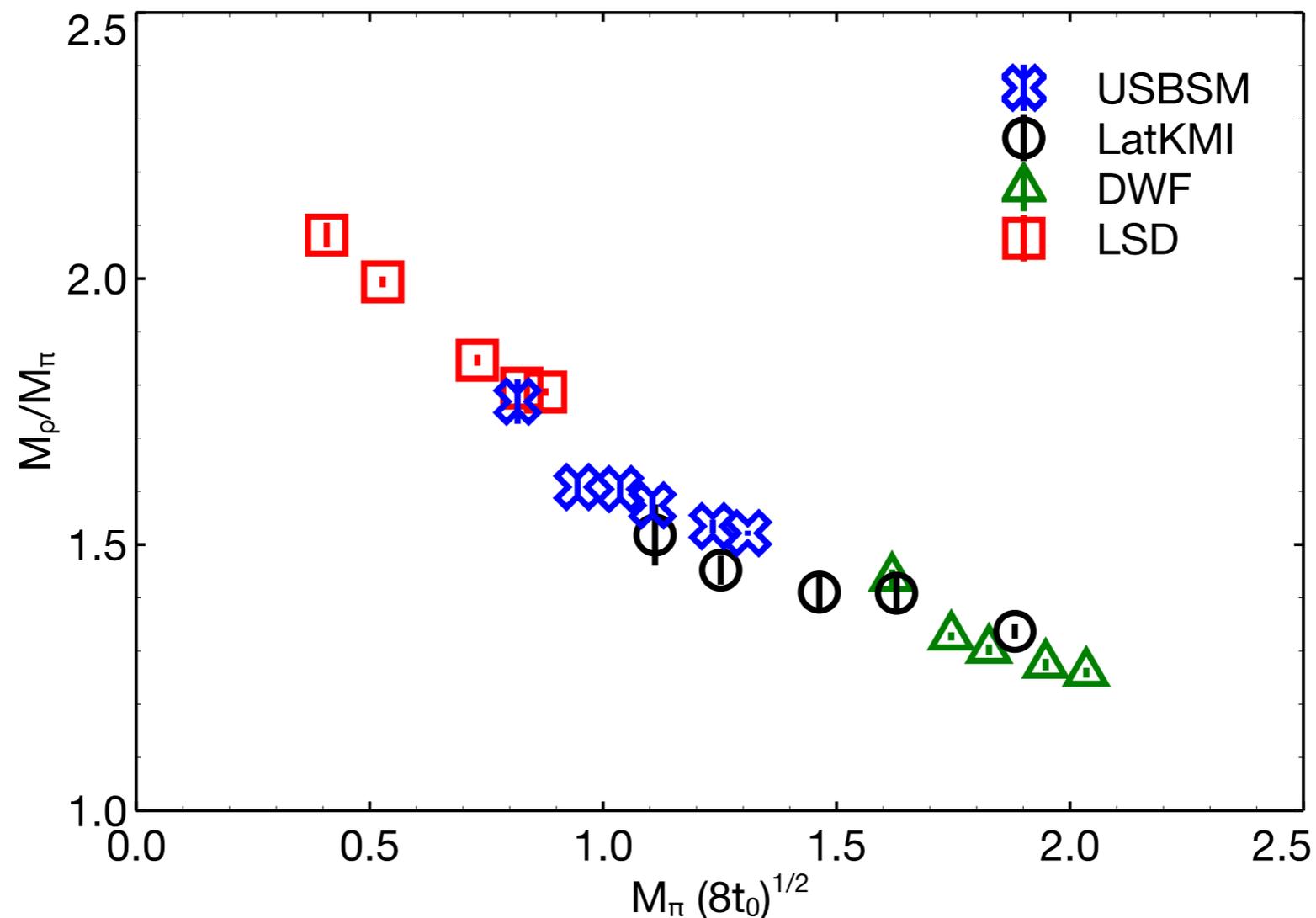
# Light hadron spectrum

- Spectrum consistent with earlier LSD  $N_f=8$  results but at lighter quark mass.
- Very strong quark mass dependence for quantities expressed in lattice units, as expected from near conformal dynamics.



# Not hyperscaling

- Mass-deformed IRFP theories have hadron masses which scale in constant ratios in approach to conformality:  $M_\rho/M_\pi \sim \text{const}$  as  $M_\pi \rightarrow 0$ .
- Pretty clear evidence that  $N_f=8$  is outside conformal window since pion is becoming light relative to rho meson. Very different from  $N_f=12$ .



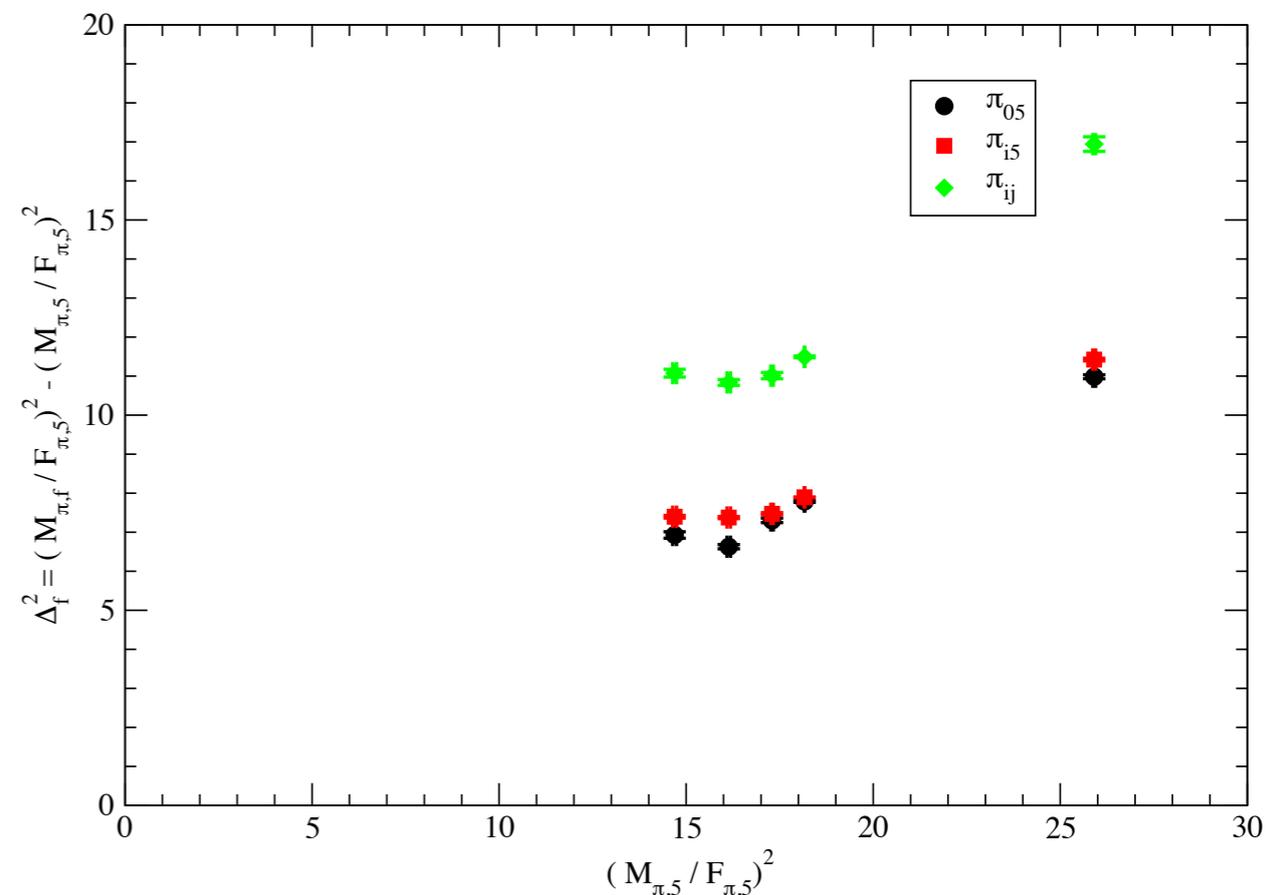
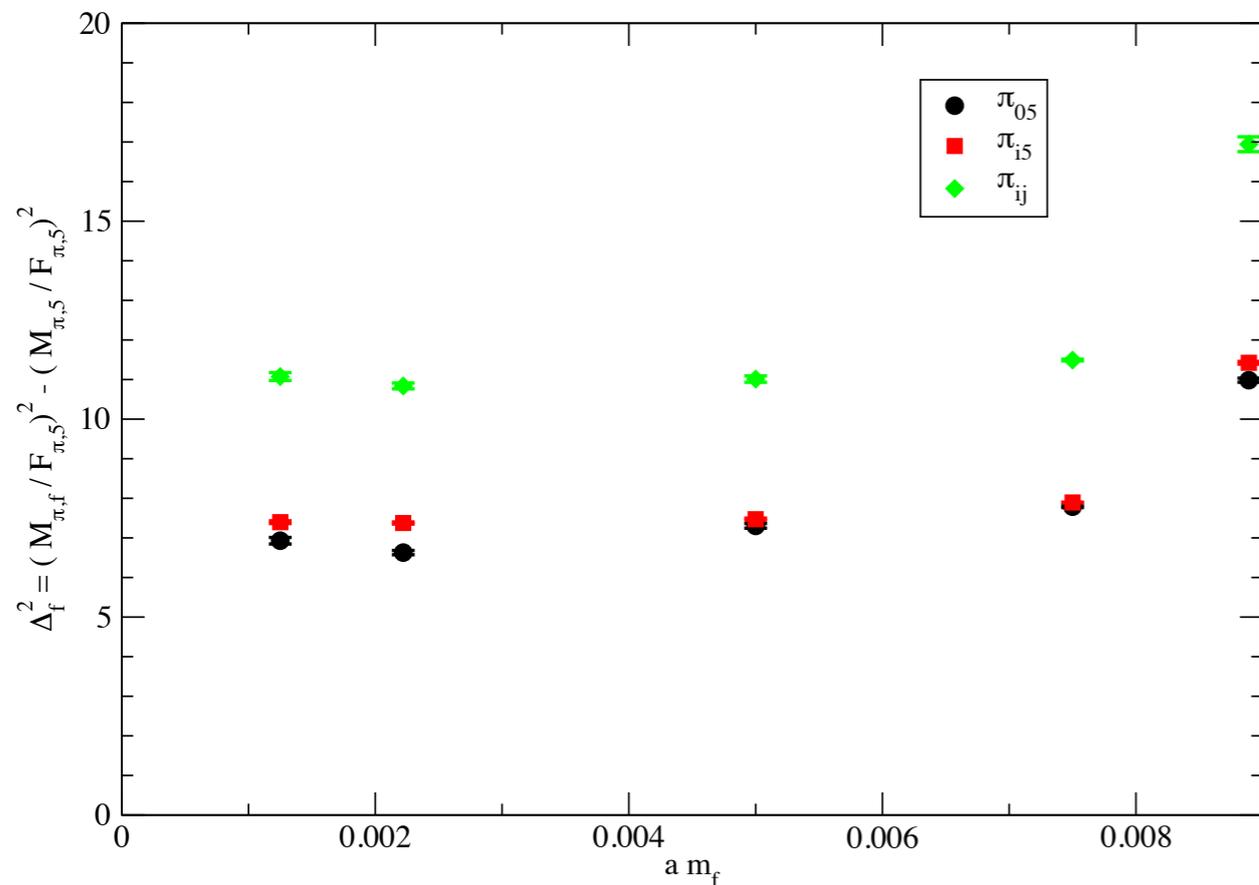
# Finite Volume Effects

- Finite volume effects are small, order 1% in sensitive quantities and order 2% in ratios at  $M_\pi L \sim 5.3$ .
- Note how rapidly quantities change in lattice units with quark mass while ratios change slowly.

| $m_f$   | volume            | $M_\pi$     | $F_\pi$     | $M_\pi / F_\pi$ | $M_\pi L$ |
|---------|-------------------|-------------|-------------|-----------------|-----------|
| 0.0075  | $24^3 \times 48$  | 0.21067(19) | 0.04746(8)  | 4.439(9)        | 5.1       |
|         | $32^3 \times 64$  | 0.20630(6)  | 0.04823(3)  | 4.228(3)        | 6.6       |
|         | $48^3 \times 96$  | 0.20575(3)  | 0.04827(1)  | 4.262(1)        | 9.9       |
| 0.005   | $32^3 \times 64$  | 0.16795(10) | 0.03939(4)  | 4.264(6)        | 5.4       |
|         | $48^3 \times 96$  | 0.16619(7)  | 0.03996(3)  | 4.159(4)        | 8.0       |
| 0.00222 | $48^3 \times 96$  | 0.11017(6)  | 0.02742(3)  | 4.017(6)        | 5.3       |
| 0.00125 | $64^3 \times 128$ | 0.08273(13) | 0.02111(11) | 3.918(24)       | 5.3       |

# Staggered flavor breaking

- Even though  $a^{-1} \sim 50 F_\pi$  the taste breaking is still rather large. Another consequence of near conformal physics.



# Wrap around effects?

- If  $M_\sigma \sim M_\pi$  wrap around effects might be noticeable for  $t \sim N_t / 2$ .
- But we choose  $M_\pi L > 5$  and  $N_t = 2L$ , so the effect will be highly suppressed. *Unless...*
  - $M_\sigma < M_\pi$  (not the case for  $N_f=8$ )
  - or  $A_{ij0} \ll A_{ij\sigma}$  (not true in general but maybe special cases)
- Looking for possible wrap around effect in  $a_0$  channel using psi-bar-psi like operator, stay tuned for details in writeup.

$$C_{I,ij}(t) = \frac{e^{-M_I t} (A_{ij0} + A_{ij\sigma} e^{-M_\sigma(N_t-t)}) + e^{-M_I(N_t-t)} (A_{ij0} + A_{ij\sigma} e^{-M_\sigma t})}{1 + e^{-M_\sigma N_t}}$$

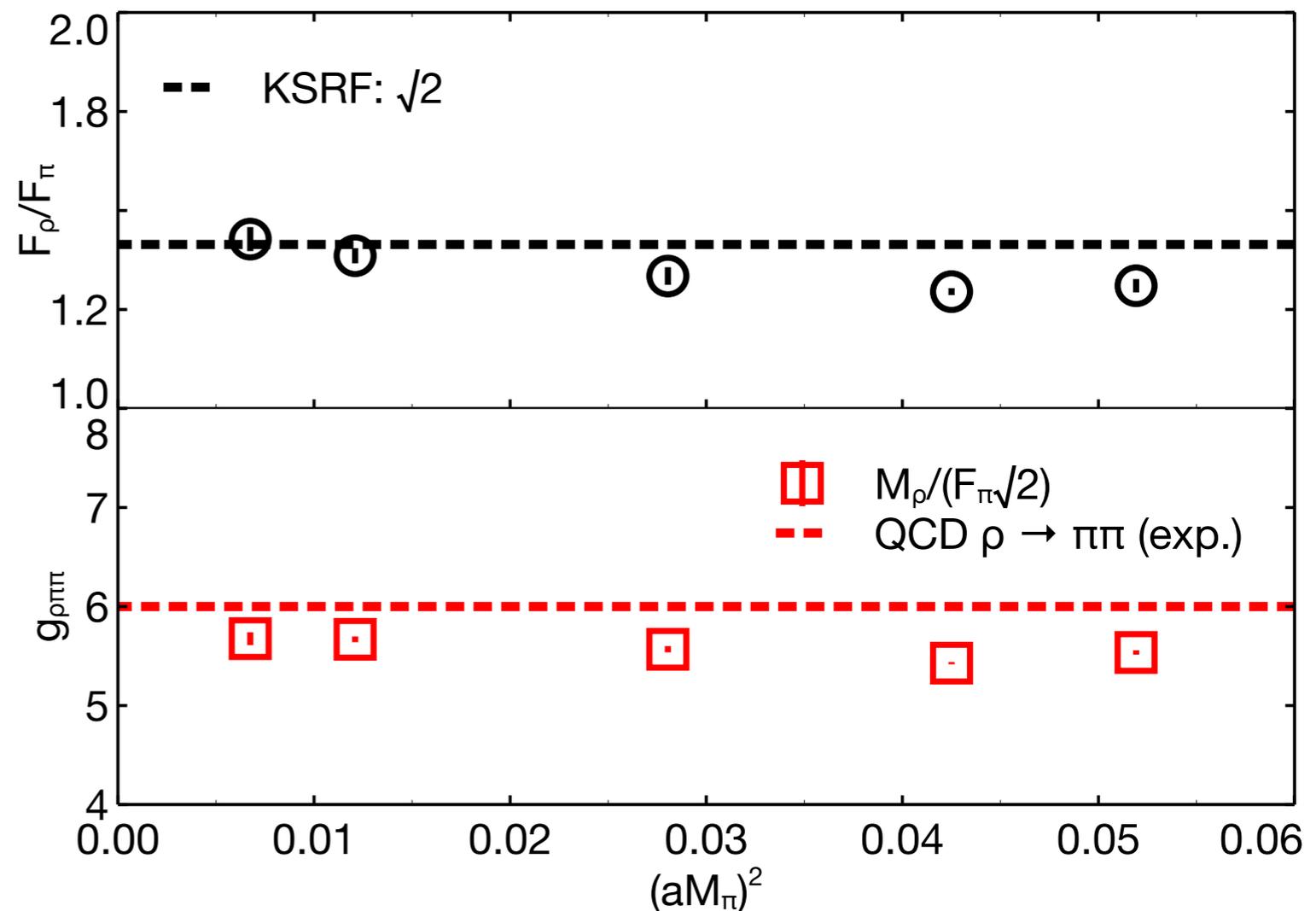
$$A_{ij0} = \langle 0 | \mathcal{O}_{I,i}^\dagger | M_I \rangle \langle M_I | \mathcal{O}_{I,j} | 0 \rangle \quad A_{ij\sigma} = \langle M_\sigma | \mathcal{O}_{I,i}^\dagger | M_I \rangle \langle M_I | \mathcal{O}_{I,j} | M_\sigma \rangle$$

# KSUF Relation

- Dynamical origin of vector meson dominance (VMD) not well understood in QCD. Is it also true in  $N_f=8$ ?
- Seems to be true, so in LHC might expect 2 TeV vector resonances with  $\sim 25\%$  width.

$$F_\rho = \sqrt{2} F_\pi, \quad g_{\rho\pi\pi} = \frac{M_\rho}{\sqrt{2} F_\pi},$$

$$\Gamma_\rho \approx \frac{g_{\rho\pi\pi}^2 M_\rho}{48\pi} \approx \frac{M_\rho^3}{96\pi F_\pi^2}$$



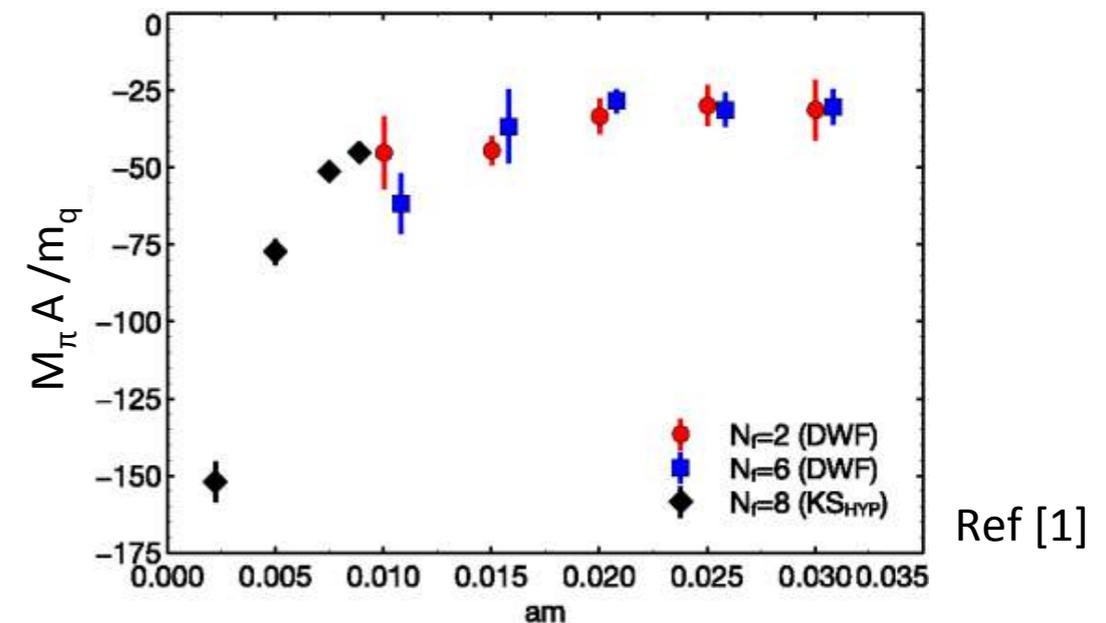
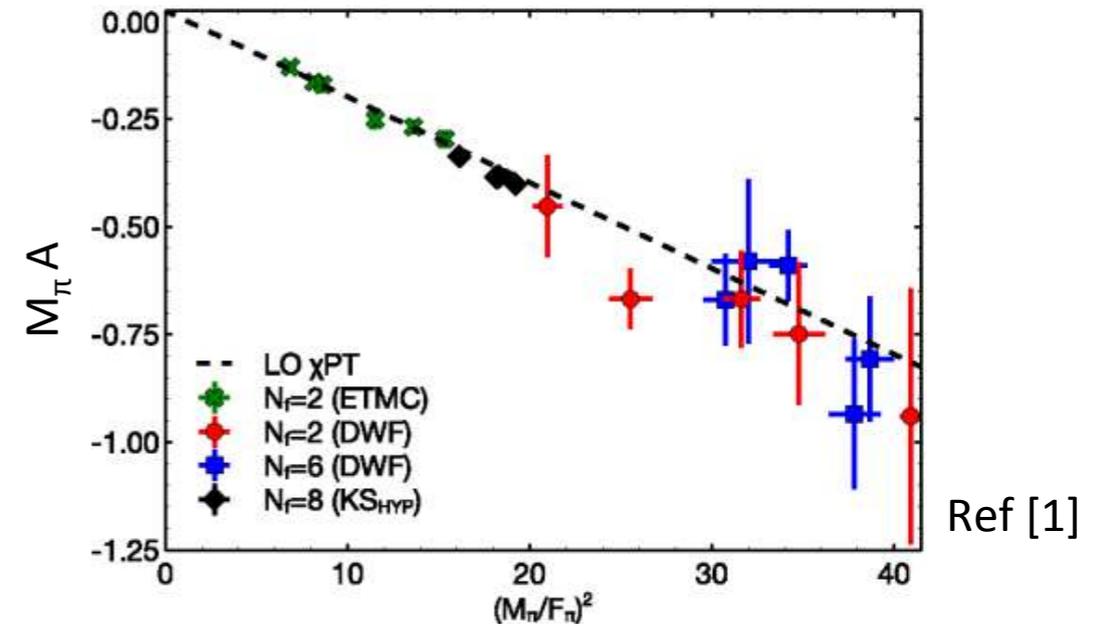
# $I=2$ $\pi\pi$ scattering

## A Clue from Pion Scattering

- Scattering length agrees well with LO XPT when plotted against physical (computed) values of  $M_\pi/F_\pi$ .

$$(M_\pi A)_{LO} = \frac{-1}{(16\pi)^2} \left( \frac{M_\pi}{F_\pi} \right)^2 = \frac{-1}{(16\pi)^2} \frac{2Bm_q}{F^2}$$

- Plotted against bare quark mass, very poor agreement with LO XPT
- Again, suggests  $F$  has significant dependence on chiral breaking at tree level



# Summary

- We now have clear examples of gauge theories with light scalars which might eventually be useful for composite Higgs phenomenology.
- Computing at masses  $m_\pi \leq f_\pi$ , where  $\chi$ PT might work, seems prohibitively expensive. So it's not clear how to extrapolate lattice results to chiral limit.
- Despite many technical difficulties due to near-conformal dynamics, useful physical results can be extracted. Can we identify a useful low-energy EFT? Listen to next talk...

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