



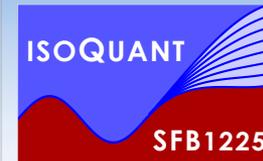
# Improved real-time dynamics from imaginary frequency lattice simulations

**Alexander Rothkopf**  
Institute for Theoretical Physics  
Heidelberg University

in collaboration with Jan Pawłowski

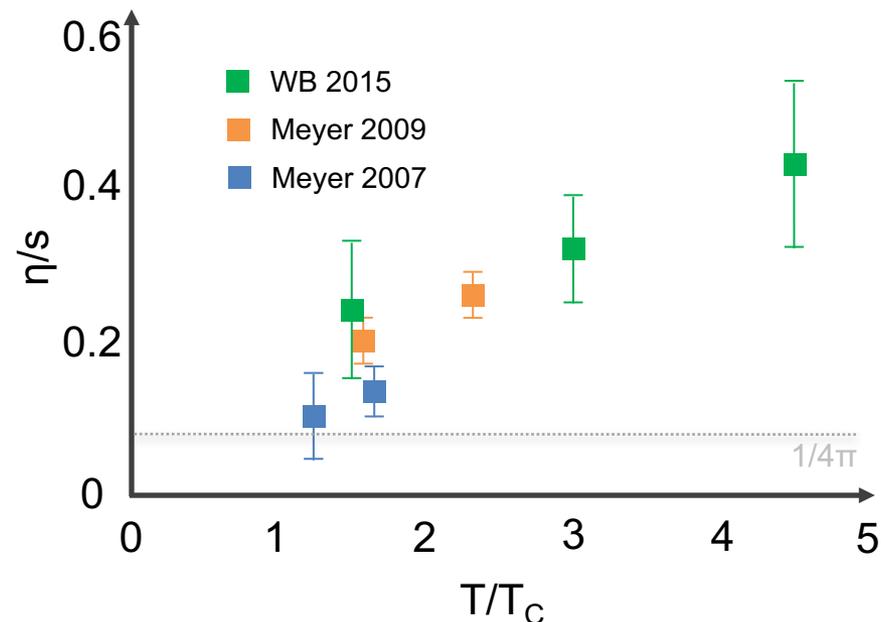
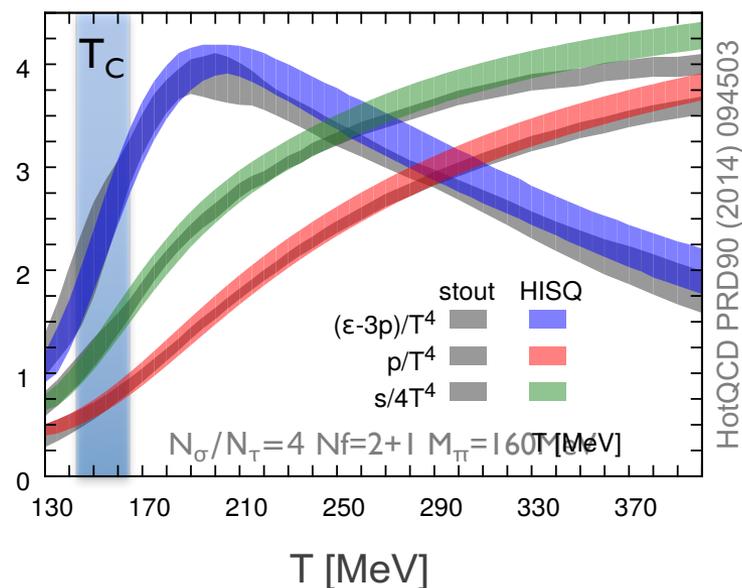
**Reference:**

J.M. Pawłowski, A.R. arXiv:1610.09531



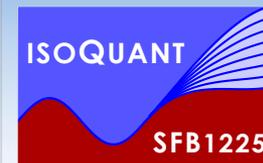
# Static vs. Real-time Observables

- LQCD highly successful for QCD E.O.S., challenged for e.g. transport coefficients



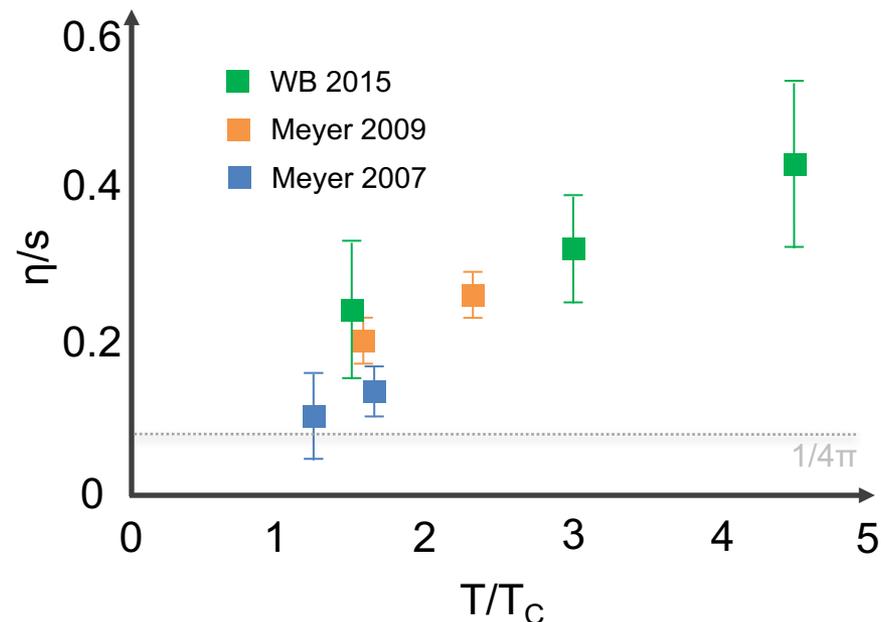
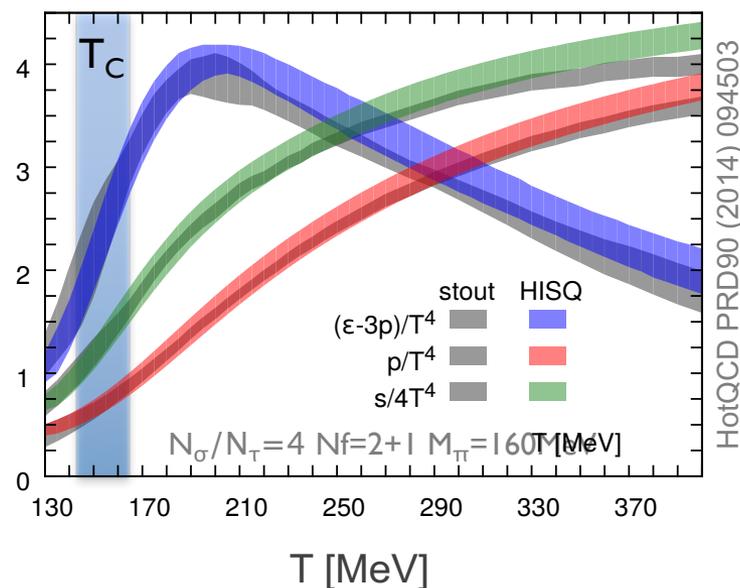
- Example: transport coefficients only accessible in LQCD via spectral functions

$$\eta = \frac{1}{20} \lim_{\mathbf{p} \rightarrow 0} \lim_{\omega \rightarrow 0} \frac{1}{\omega} \int d^4x e^{i\mathbf{p}\cdot\mathbf{x}} \langle [T_{12}(0), T^{12}(x)] \rangle$$



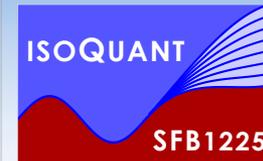
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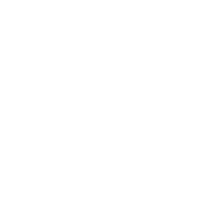
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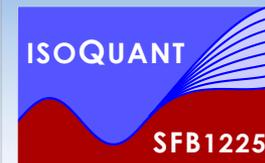


# The underlying difficulty

- Intrinsic problem of standard spectral reconstruction: exponential information loss

$$D(\tau) = \int_0^{\infty} d\omega \frac{\cosh[\omega(\tau - \beta/2)]}{\sinh[\omega\beta/2]} \rho(\omega)$$



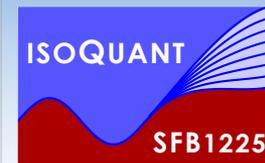


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- 1<sup>st</sup> part of the remedy: go over to imaginary frequencies (hint of possible exp. improvement)

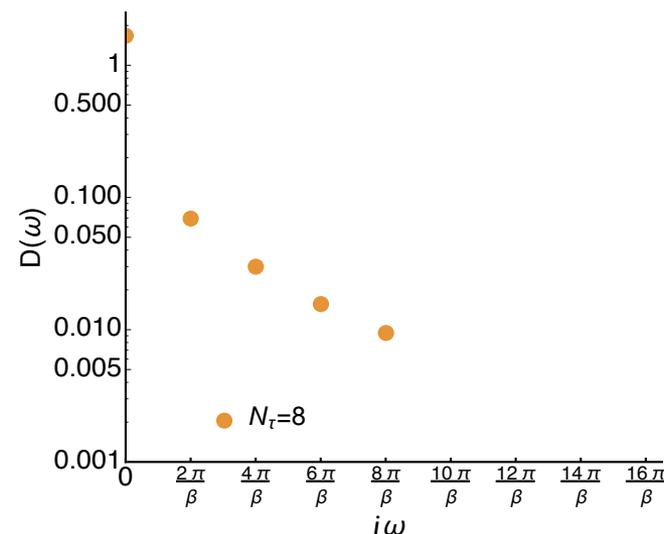
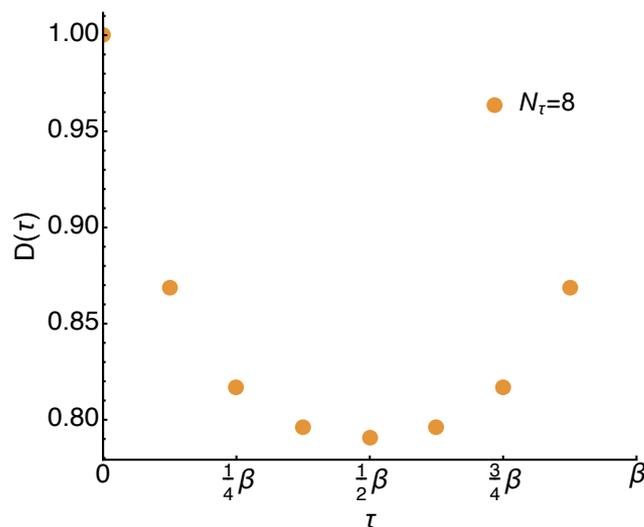


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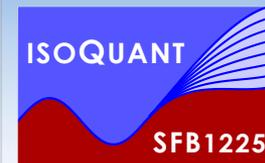
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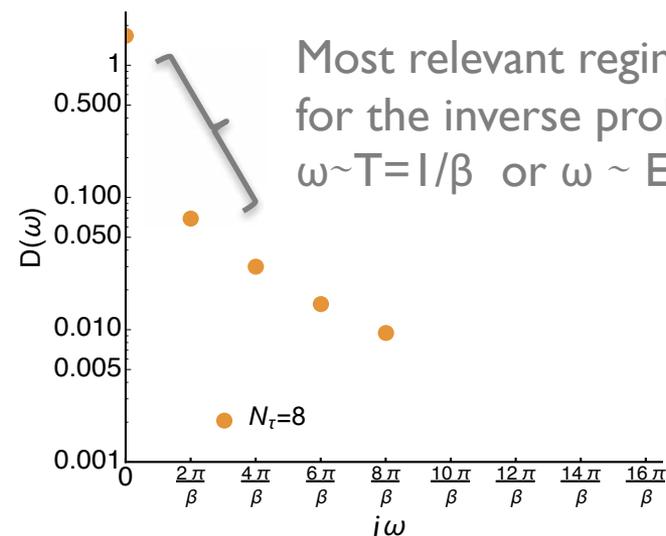
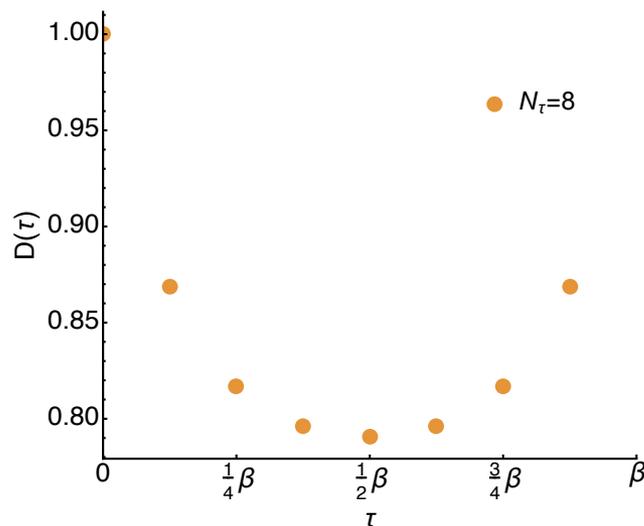


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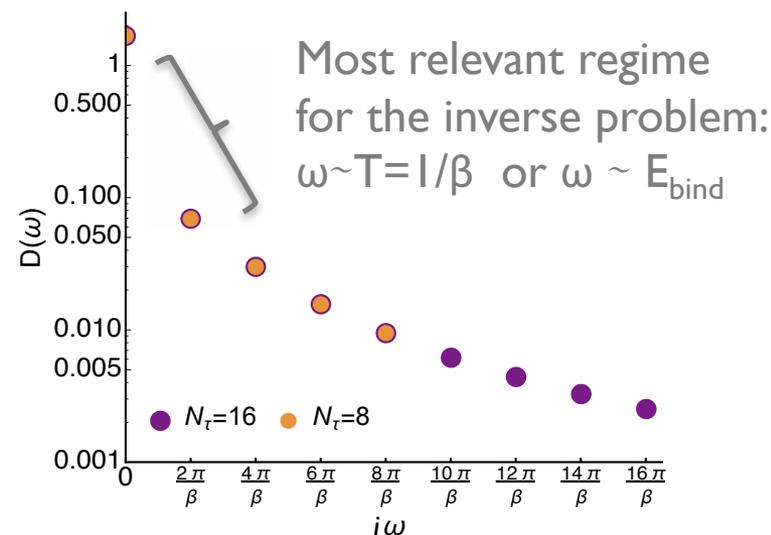
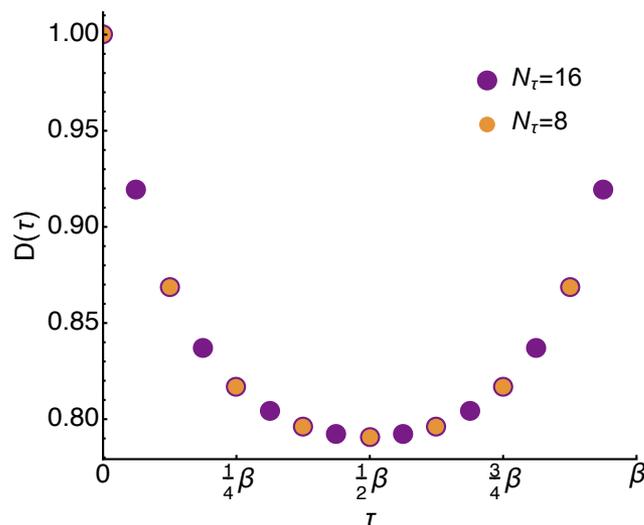


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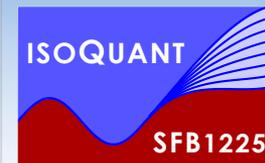
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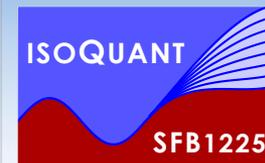


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## Our proposal



# Simulating thermal fields in imaginary frequencies

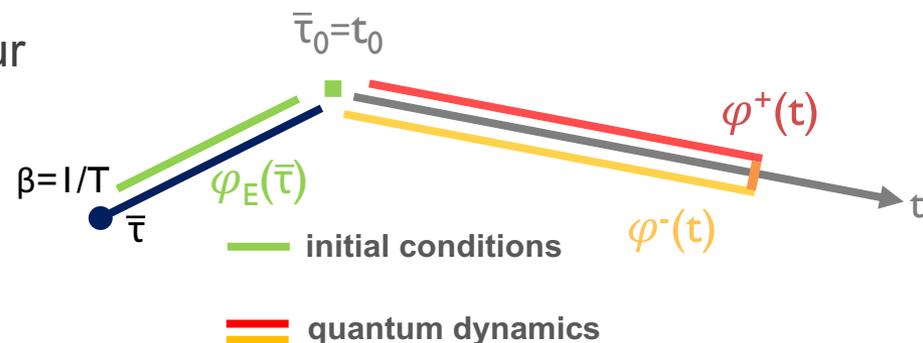


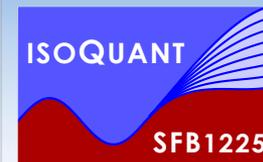
# A thermal initial value problem

$$\mathcal{Z} = \int_{\varphi_E(0)=\varphi_E(\beta)} \mathcal{D}\varphi_E e^{-S_E[\varphi_E]} \int_{\varphi^+(t_0, \mathbf{x})=\varphi_E(0)}^{\varphi^-(t_0, \mathbf{x})=\varphi_E(\beta)} \mathcal{D}\varphi e^{iS_M[\varphi^+] - iS_M[\varphi^-]}$$

## Setup on the Schwinger-Keldysh contour

- For finite  $t$   $\varphi^+$  and  $\varphi^-$  path coupled at  $t$
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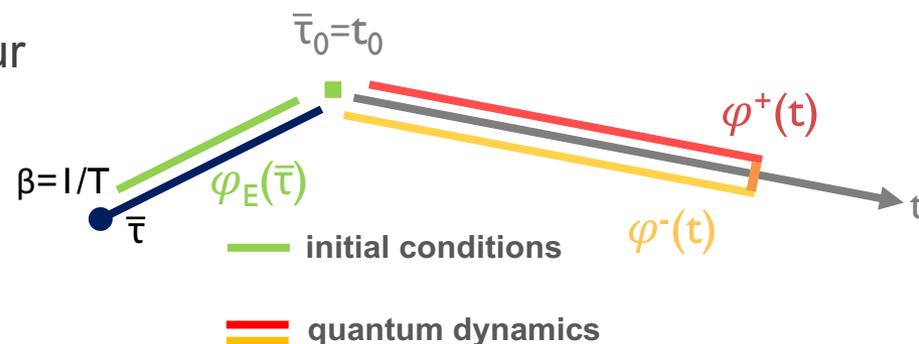


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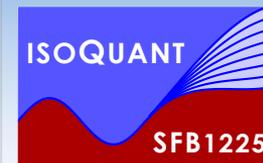
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$$G^{++}(p^0, \mathbf{p}) = \int \frac{dq^0}{2\pi i} \frac{\rho(q^0, \mathbf{p})}{p^0 - q^0 + i\epsilon} = n(p^0) \rho(p^0, \mathbf{p})$$

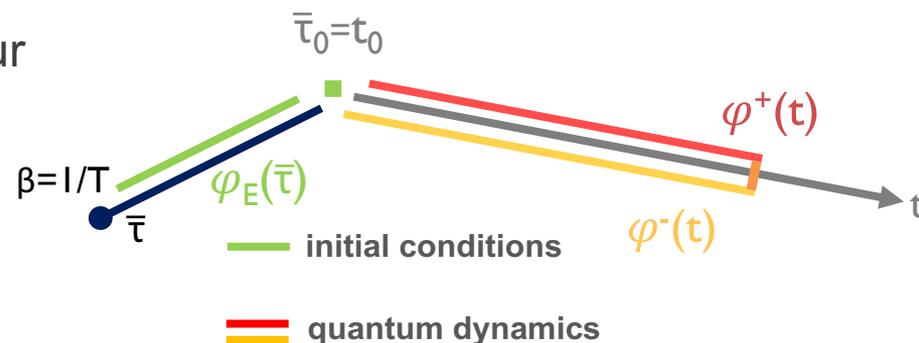


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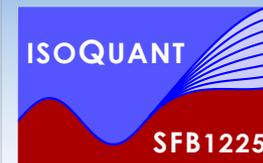
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- Correlations between any finite time  $t$  on forward branch and endpoint are damped

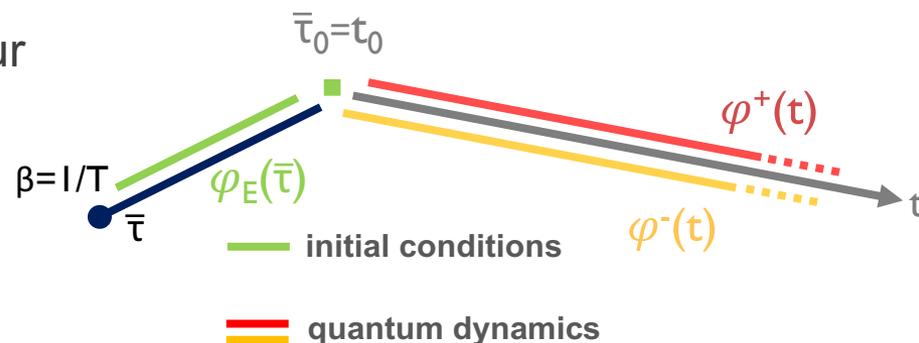


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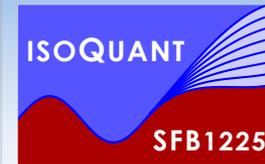
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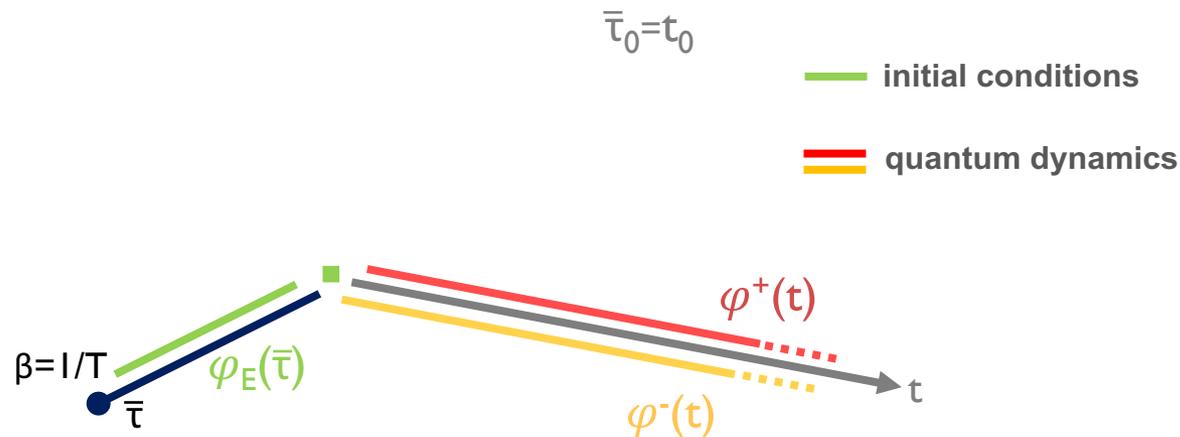
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# Analytic continuation

- Our idea: rotate branches of the real-time contour into noncompact Euclidean time

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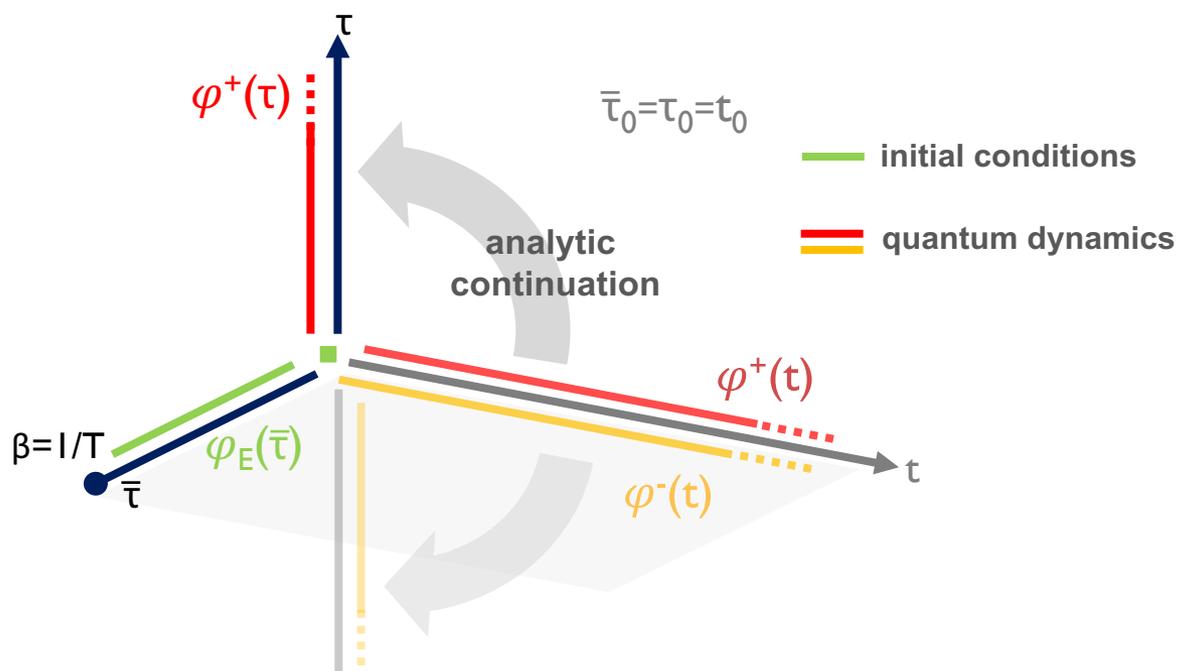


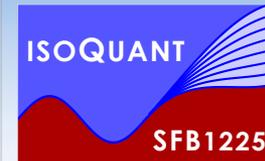


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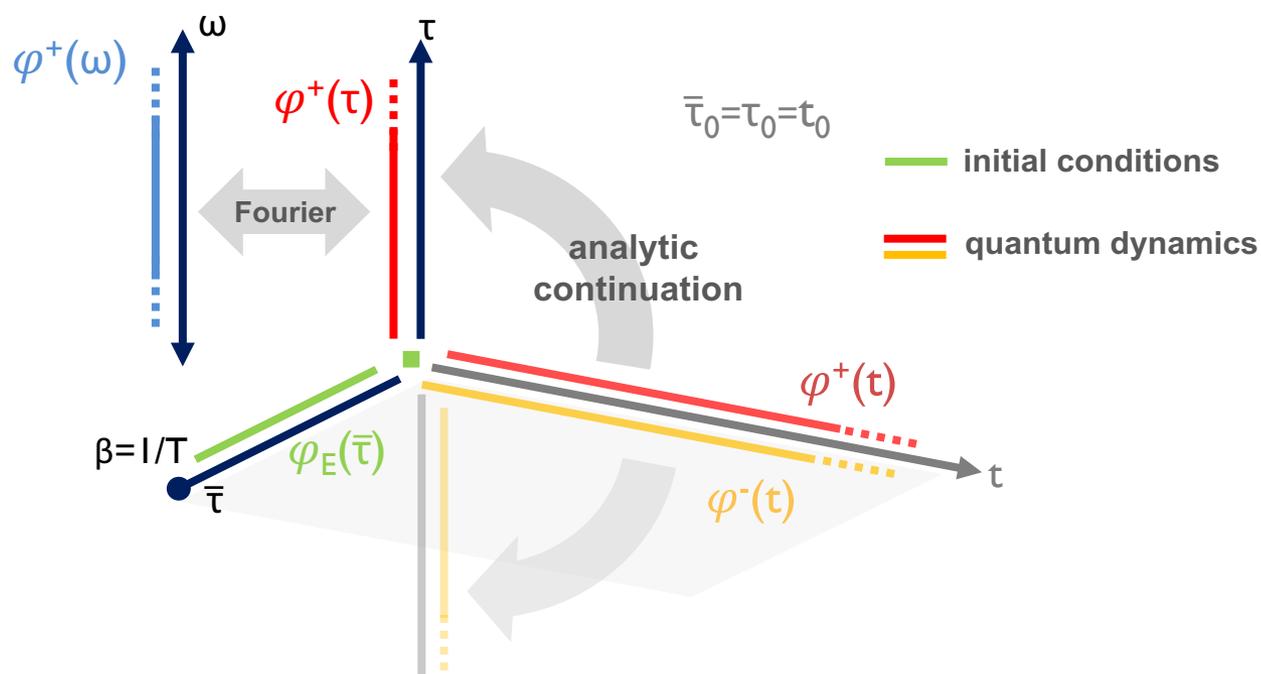




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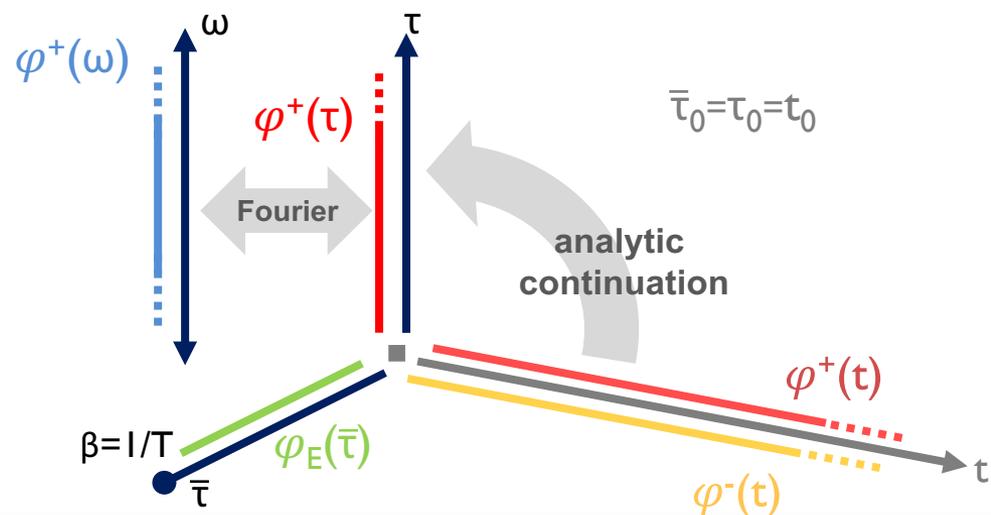
# Simulating $\varphi^+$ in $(0+1)$ dimensions

- Combine standard simulation of  $\varphi_E(\bar{\tau})$  with additional simulation of  $\varphi^+(\tau)$

$$S_E = \int d\tau \left( \underbrace{\frac{1}{2}(\partial_\tau \varphi_E)^2 + \frac{1}{2}m^2 \varphi_E^2}_{S_E^0} + \underbrace{\frac{\lambda}{4!} \varphi_E^4}_{S_E^{\text{int}}} \right)$$

$$\partial_{t_5} \varphi_E(\bar{\tau}) = - \frac{\delta S_E[\varphi_E]}{\delta \varphi_E(\bar{\tau})} + \eta(\bar{\tau}) \quad \bar{\tau} \in [0, \beta = 1/T]$$

$$\langle \eta(\tau) \eta(\tau') \rangle = 2\delta(\tau - \tau')$$





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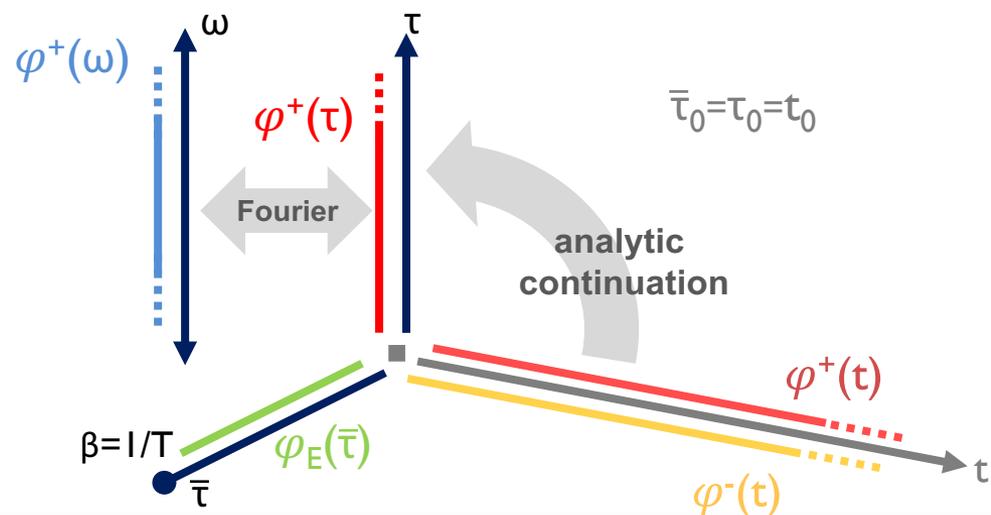
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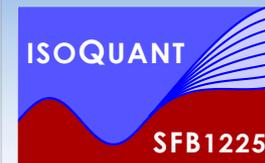
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$\langle \eta(\tau)\eta(\tau') \rangle = 2\delta(\tau - \tau')$

- Mixed representation: evaluate drift terms, where each of them is local

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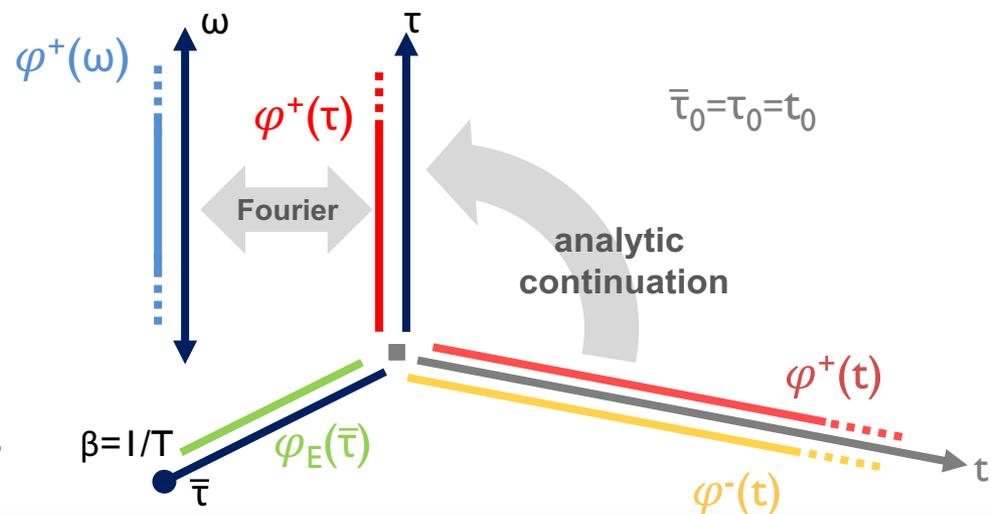
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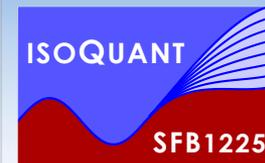
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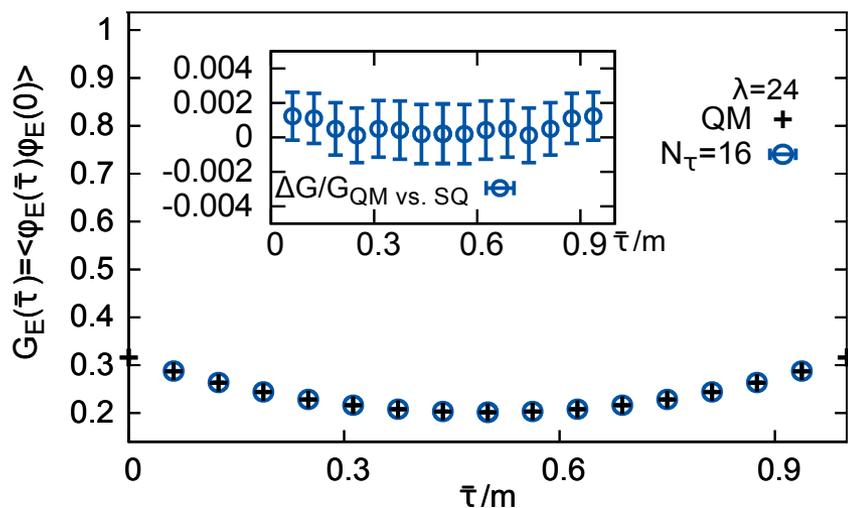
- $T$  enters only via interaction term, which is local in  $\tau$
- Strategy: when evaluating  $\delta S_E^{\text{int}}/\delta \varphi^+(\tau)$  replace  $\varphi^+(\tau=0) = \varphi_E(0)$
- Using DFT to change between  $\varphi^+(\omega)$  and  $\varphi^+(\tau)$  incurs finite volume artifacts





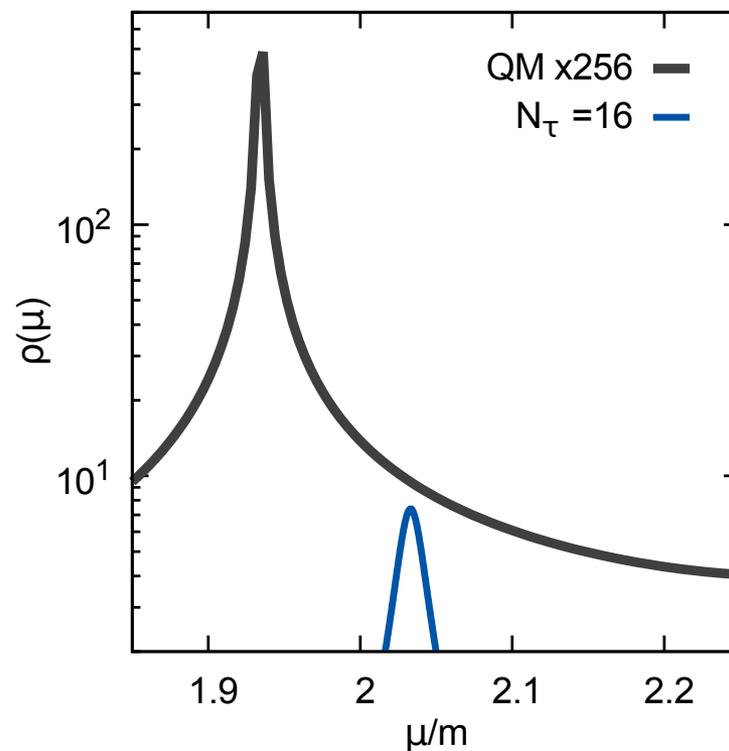
# Spectra from Euclidean times

- Standard update for  $\varphi_E$  with  $N_\tau=16$   $m=1$   $\lambda=24$   $N_\tau d\tau=1$  (compare to QM of A.H.O.)
- Spectrum from Euclidean time data with  $\Delta G_E/G_E=10^{-3}$ : peak position still off

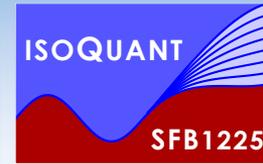


Spectral reconstruction via Bayesian BR method.

Y. Burnier, A.R.  
Phys.Rev.Lett. 111 (2013) 182003

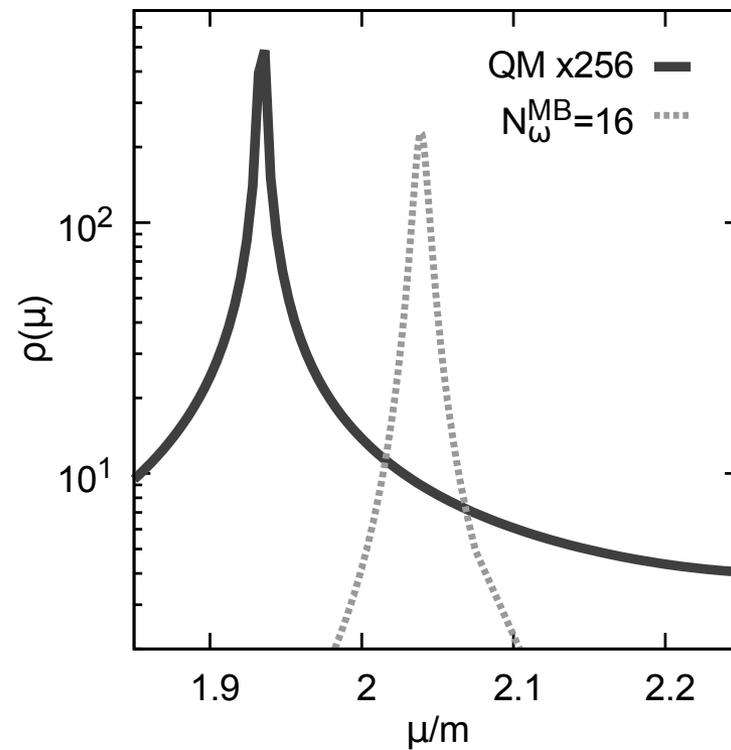
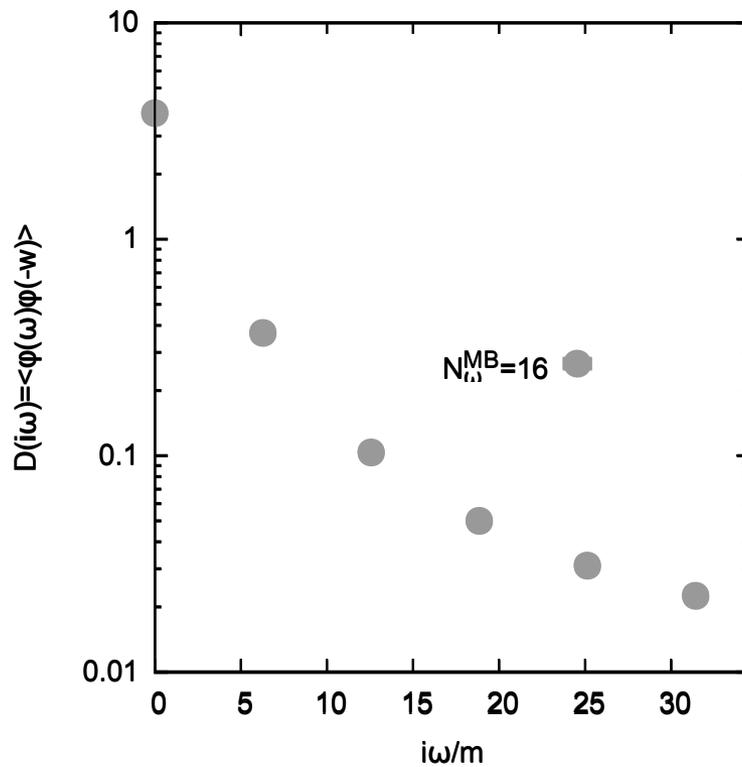


J.M. Pawłowski, A.R. arXiv:1610.09531



# Spectra from imaginary frequencies

- Just going over to Matsubara frequencies does not give better peak position



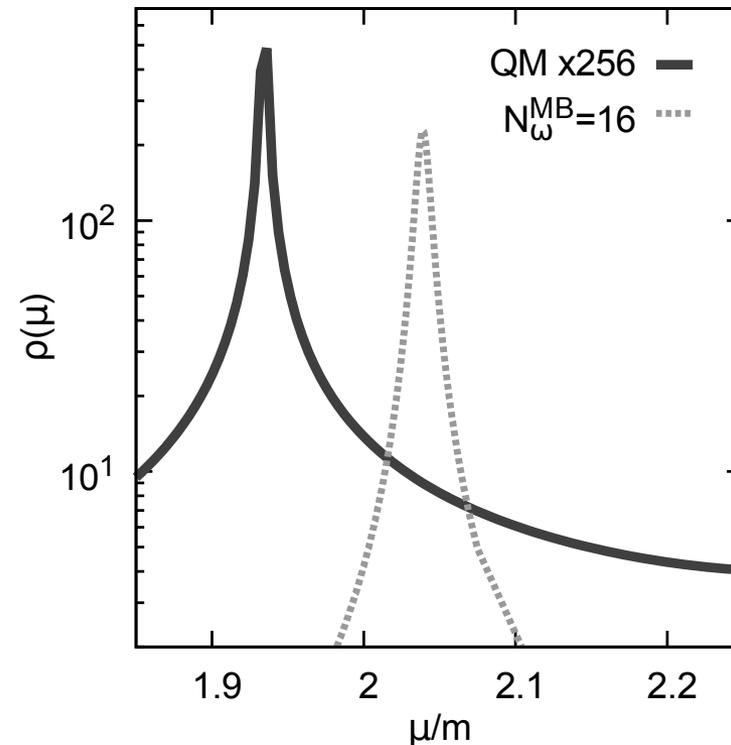
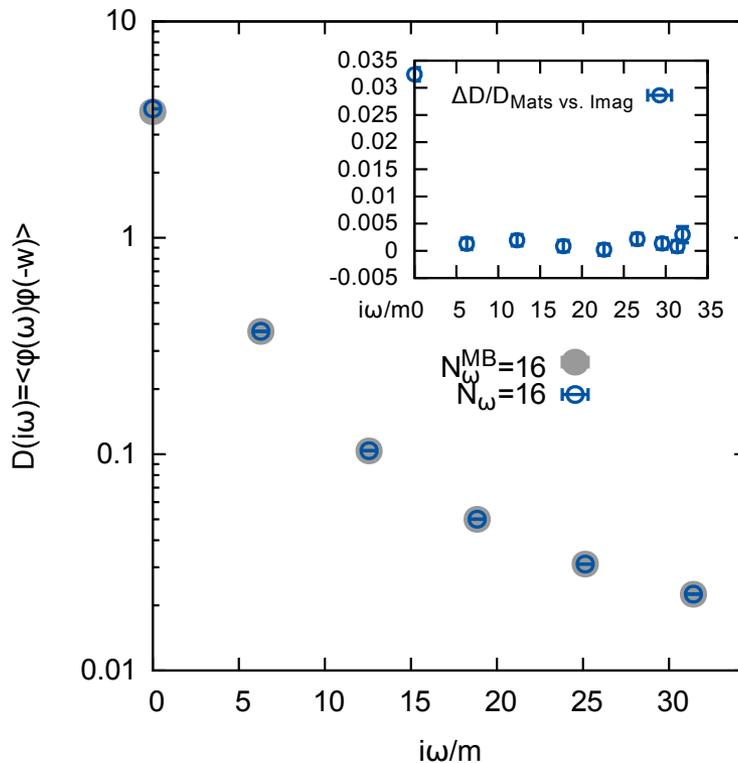
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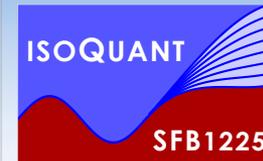


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- Just going over to Matsubara frequencies does not give better peak position
- Mixed update for  $\varphi^+$  in imaginary frequencies  $\omega$  and imaginary time  $\tau$

$$\partial_{t_5} \varphi^+(\omega_l) = -\frac{\delta S_0}{\delta \varphi^+(\omega_l)} - \frac{\delta S_E^{\text{int}}}{\delta \varphi^+(\tau_j)} \frac{\delta \varphi^+(\tau_j)}{\delta \varphi^+(\omega_l)} + \eta(\omega_l)$$

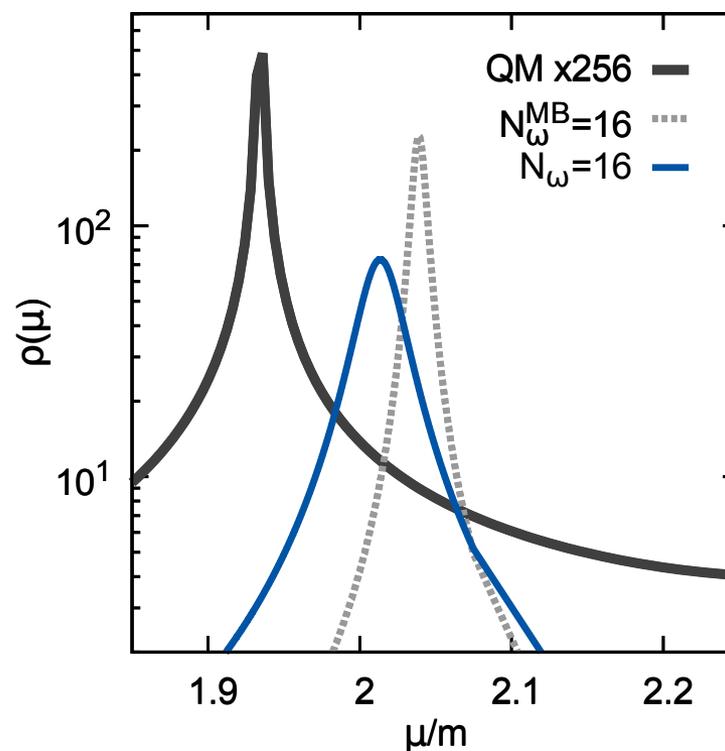
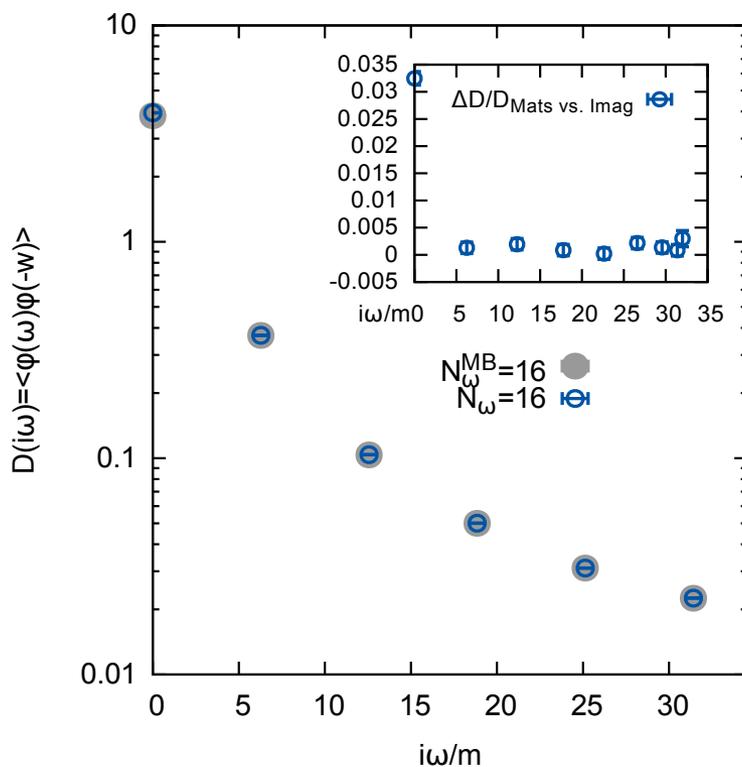




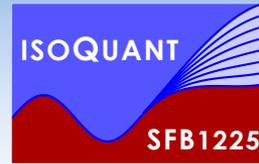
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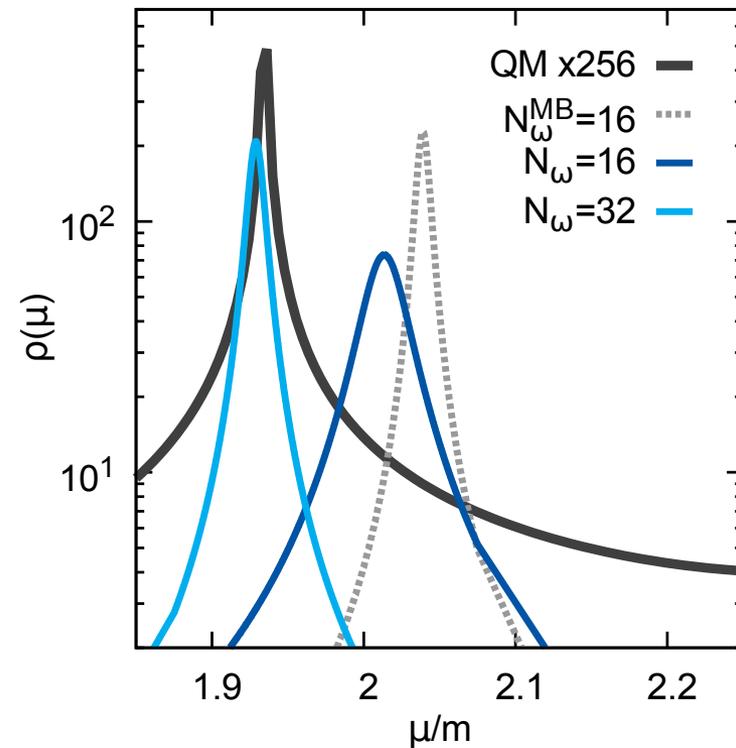
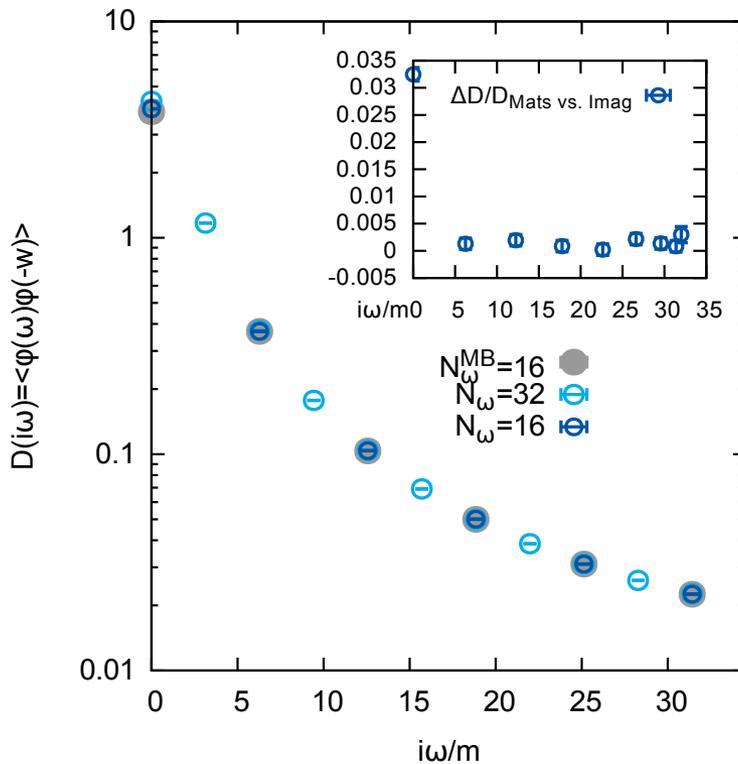
J.M. Pawłowski, A.R. arXiv:1610.09531



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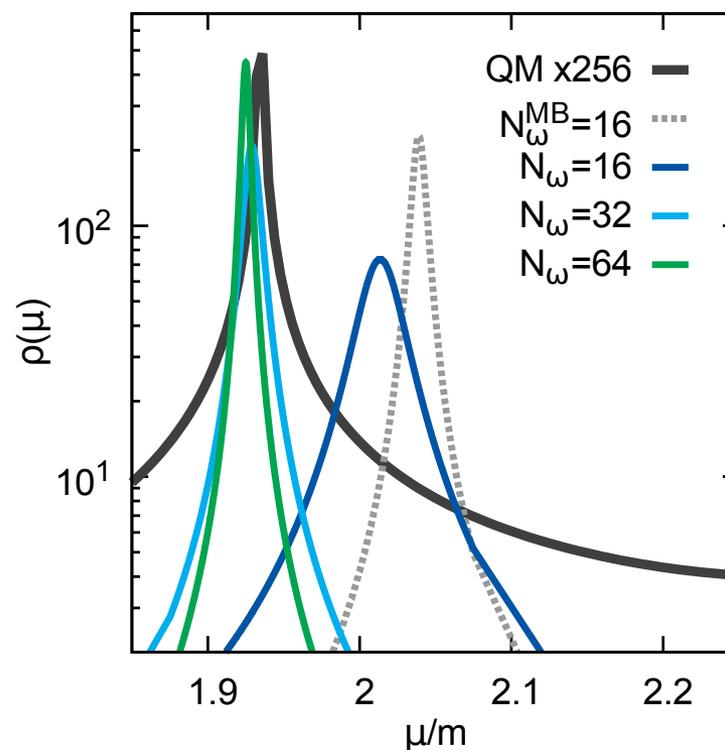
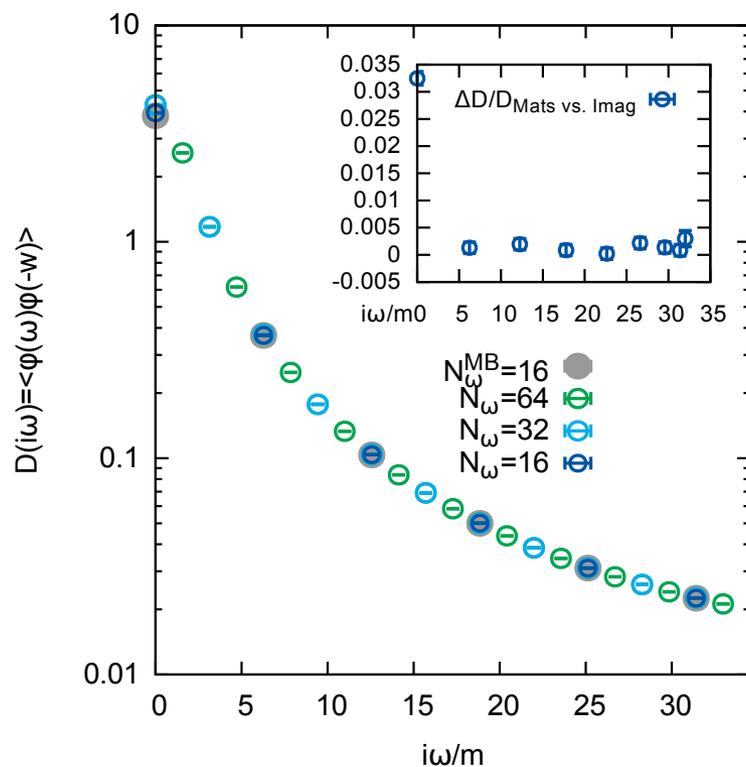
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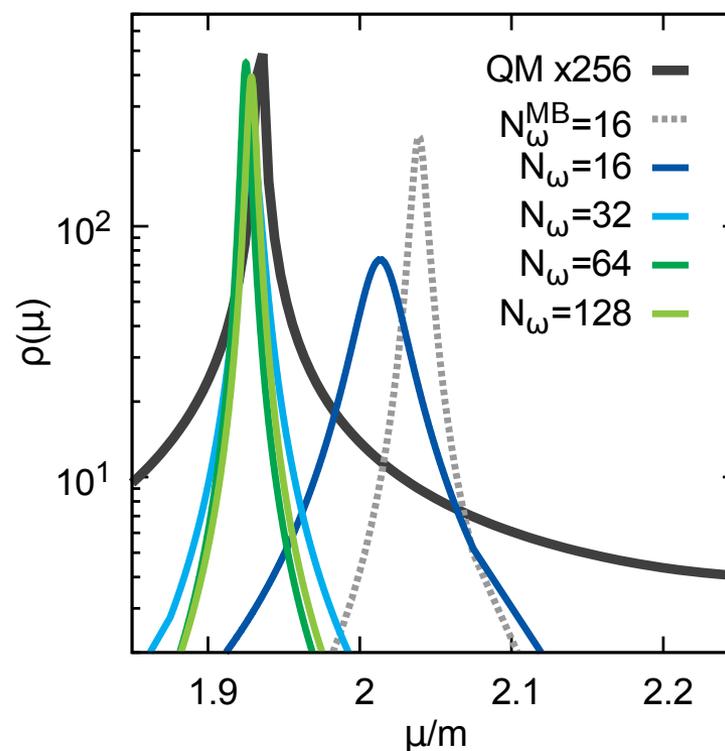
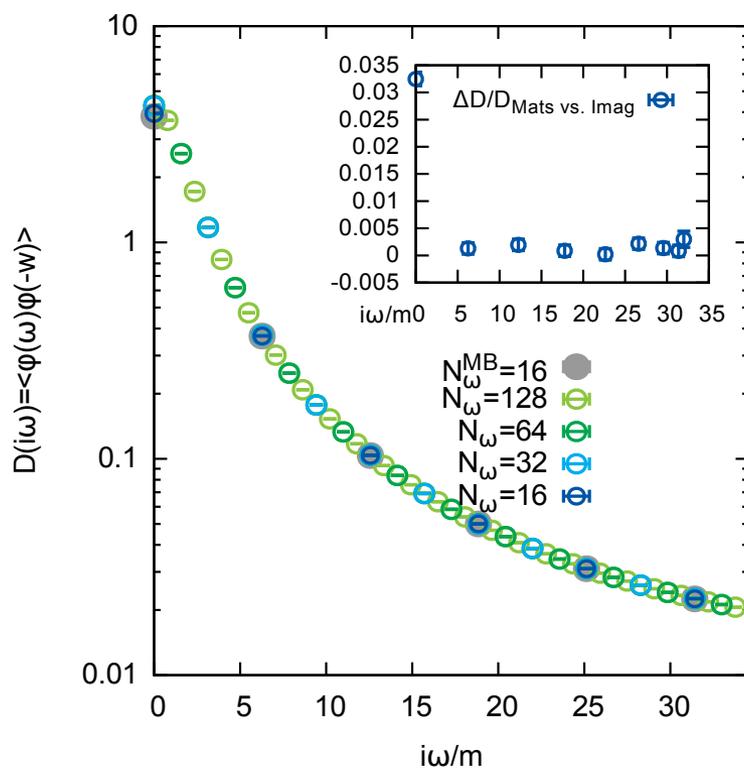
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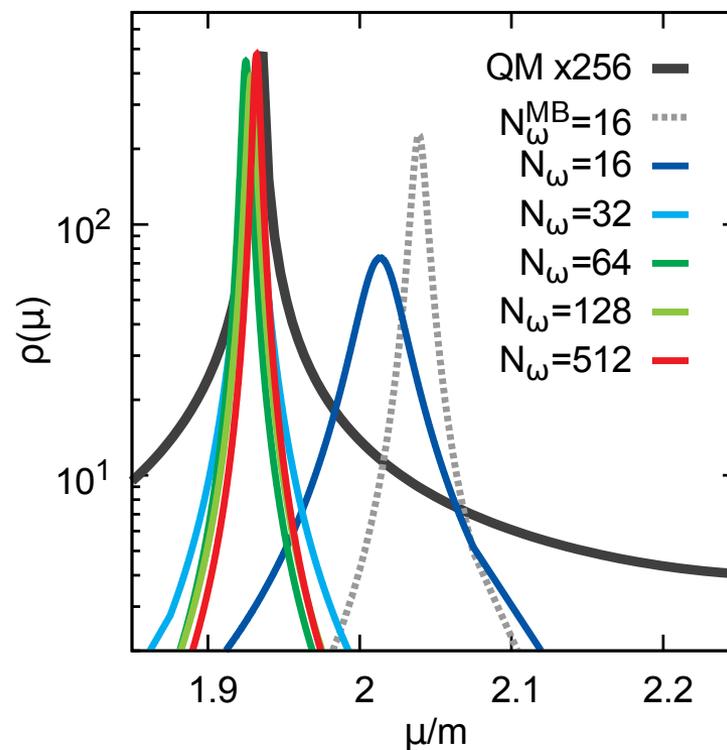
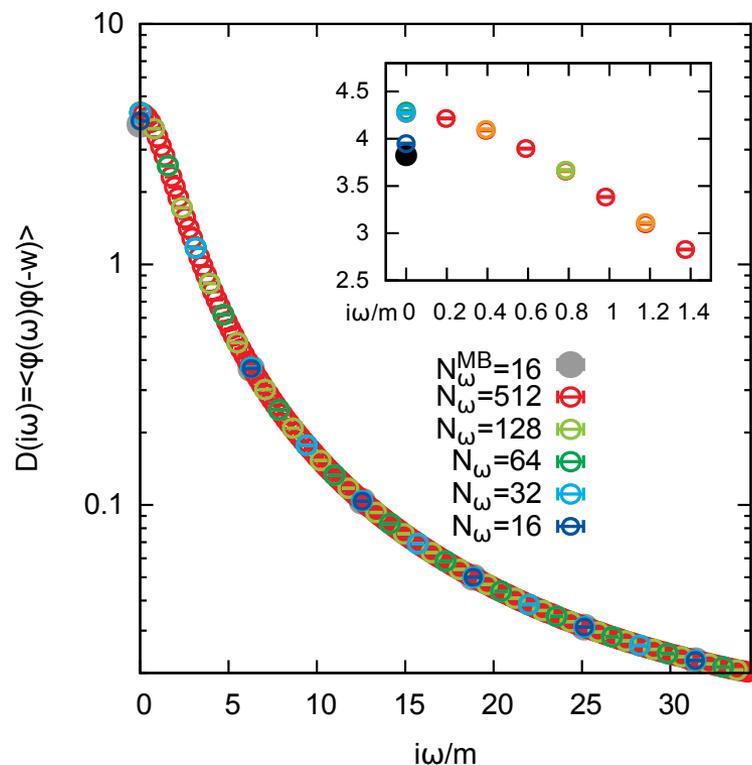




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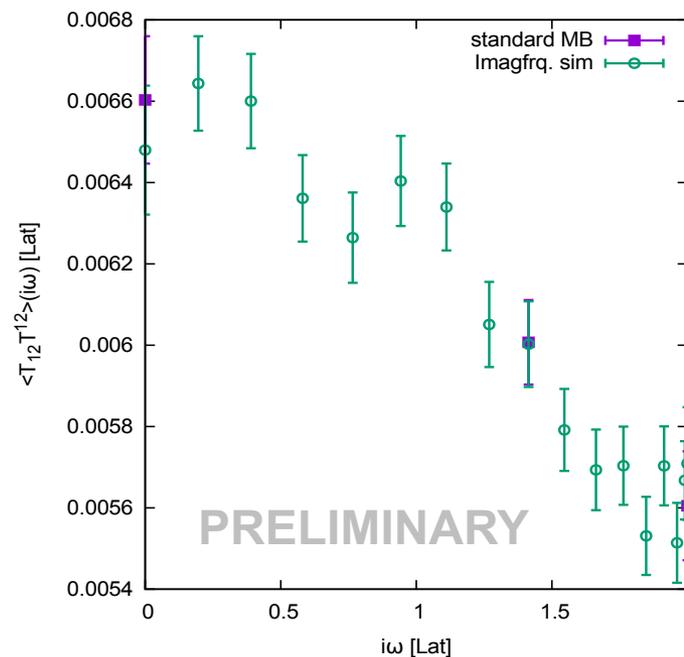
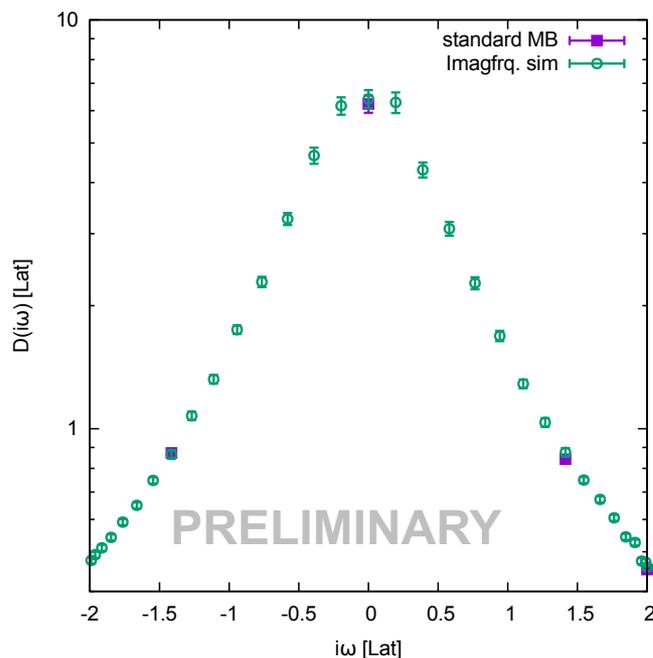


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# Towards genuine field theories

- Complex scalar field in 3+1d  $m^2=0.2$   $\lambda=1$   $d\tau=dx=1$   $N=8$   $N_\tau=4$   $N_\omega=32$

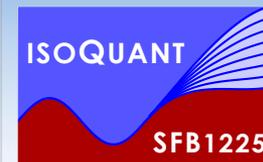


J.M. Pawłowski, A.R. in preparation

- Observables  $\langle \varphi(x) \varphi^*(y) \rangle$  and  $\langle T_{12} T^{12} \rangle \approx \langle (T_{11} - T_{22}) (T^{11} - T^{22}) \rangle$

$$(T_{11} - T_{22})(x) = \partial_1 \varphi(x) \partial_1 \varphi^*(x) - \partial_2 \varphi(x) \partial_2 \varphi^*(x)$$

- First low statistics runs promising but for gauge theory still conceptual challenges



# Summary

- Spectral reconstructions from standard Euclidean simulations exponentially hard
- Relevant physics encoded in correlators away from Matsubara frequencies
- Proposal for a different treatment of thermal fields
  - Consider initial-value problem on the Schwinger-Keldysh contour
  - Rotate the real-time branches into a non-compact Euclidean time
  - Setup mixed representation stochastic quantization for the forward branch since  $\rho(\omega)$  can be related to  $G^{++}(i\omega)$  via rational integral kernel
- First numerical test in a (0+1)d scalar toy model
  - Agreement of correlators on Matsubara frequencies (except at  $\omega=0$  due to finite volume)
  - Reconstructed spectra improve significantly if imaginary frequencies between Matsubara are resolved.
- **First test** for complex scalar in 3+1 dim. and **ongoing efforts** for gauge fields