

# **Decay constants and semileptonic form factors of $B_c$ mesons**

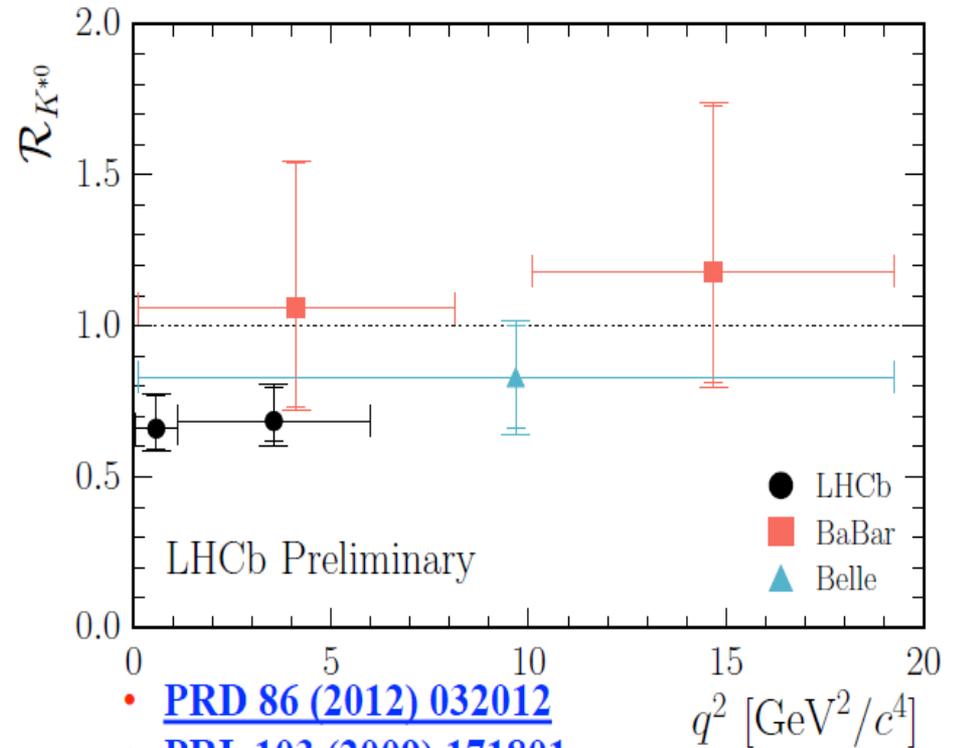
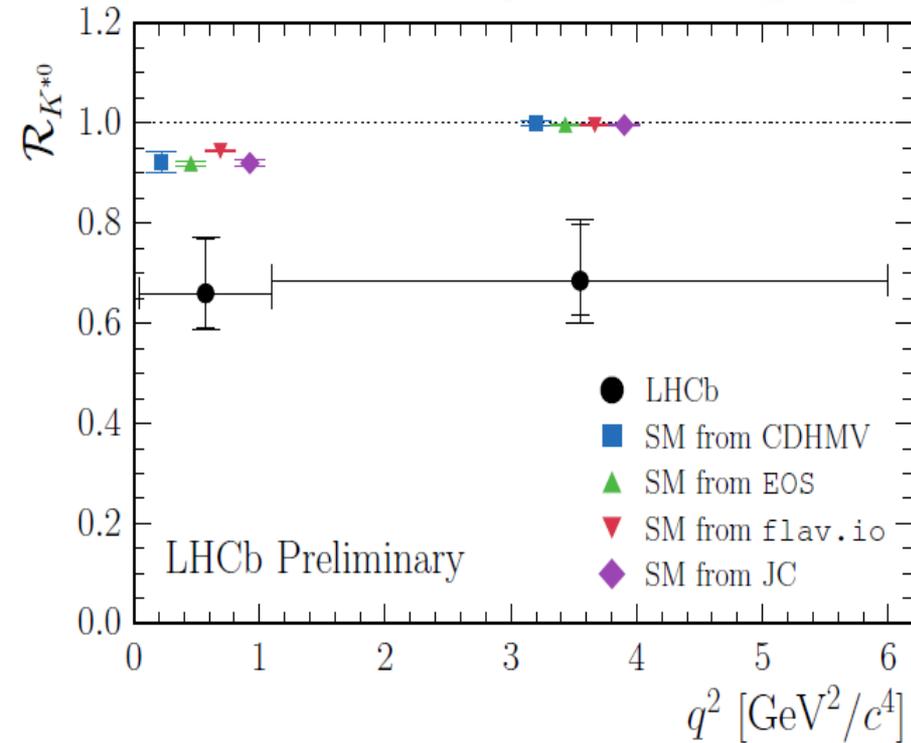
*Nilmani Mathur*

**Department of Theoretical Physics  
Tata Institute of Fundamental Research, India**

**Collaborator : M. Padmanath**

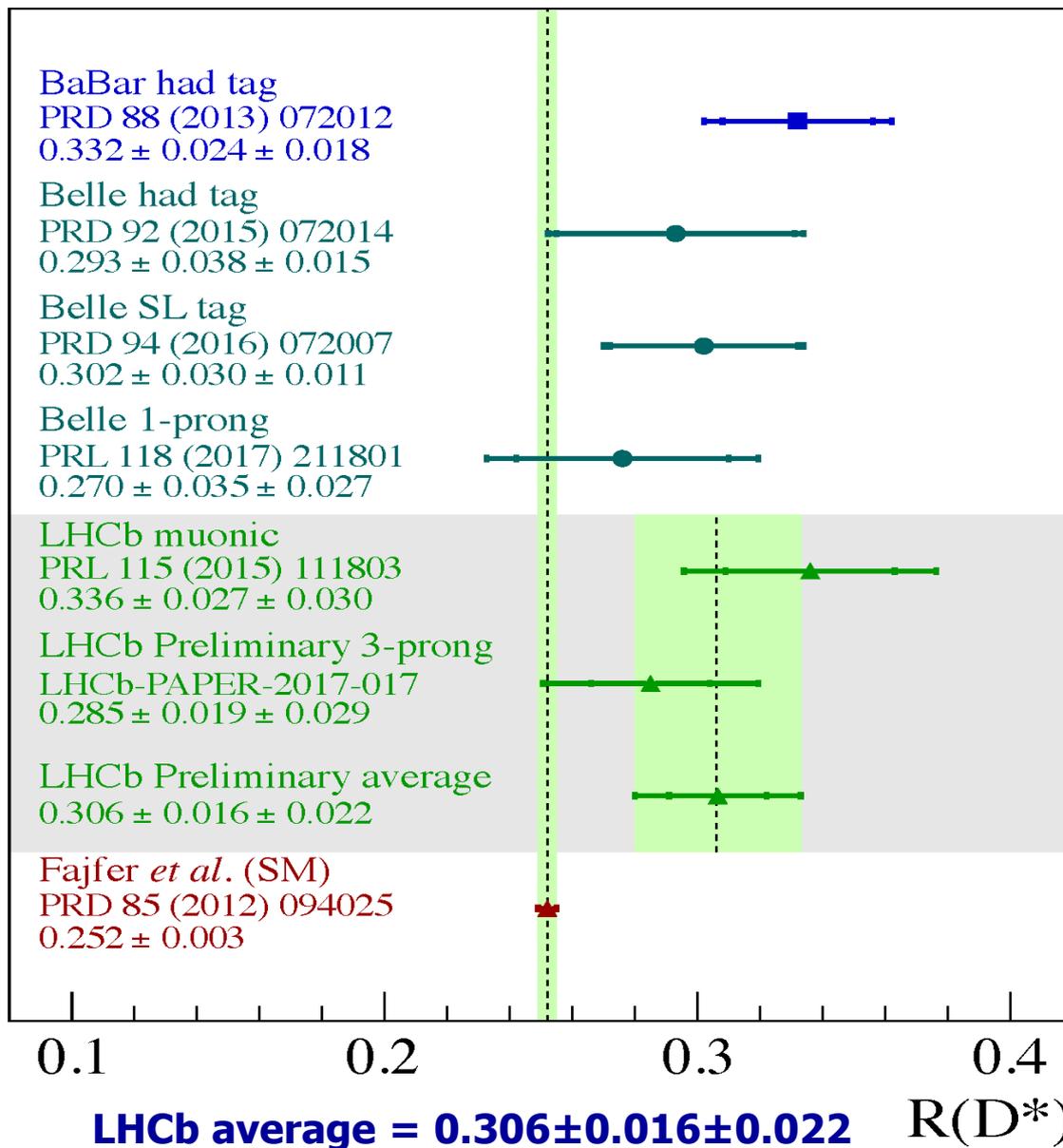
**Lattice 2017, Granada, Spain**

$$\mathcal{R}_{K^{*0}} = \frac{\mathcal{B}(B^0 \rightarrow K^{*0} \mu^+ \mu^-)}{\mathcal{B}(B^0 \rightarrow K^{*0} J/\psi (\rightarrow \mu^+ \mu^-))} \bigg/ \frac{\mathcal{B}(B^0 \rightarrow K^{*0} e^+ e^-)}{\mathcal{B}(B^0 \rightarrow K^{*0} J/\psi (\rightarrow e^+ e^-))}$$



- [PRD 86 \(2012\) 032012](#)
- [PRL 103 \(2009\) 171801](#)

LHCb Preliminary	low- $q^2$	central- $q^2$
$\mathcal{R}_{K^{*0}}$	$0.660 \pm_{-0.070}^{+0.110} \pm 0.024$	$0.685 \pm_{-0.069}^{+0.113} \pm 0.047$
95% CL	[0.517–0.891]	[0.530–0.935]
99.7% CL	[0.454–1.042]	[0.462–1.100]



# Motivation

- $|V_{cb}|$  from  $b \rightarrow c$  transition in semileptonic decay
- Compare  $R(B_c)$  and  $R(B_c^*)$  with standard model prediction
- Expt :
  - ✓  $B_c \rightarrow J/\psi \mu \nu$  R. Aaij et al. [LHCb Collaboration],  
Eur. Phys. J. C 74, 2839 (2014)
  - $B_c \rightarrow J/\psi \tau \nu$  Soon?
- Lattice :
  - $B_c \rightarrow \eta_c l \nu$  Colquhoun et al **HPQCD : 1611.01987**
  - $B_c \rightarrow J/\psi l \nu$  A. Lytle : CKM2016

# Set up

- Gauge Configurations: HISQ 2+1+1
- Valence quark propagators :
  - Light to charm : Overlap
  - Bottom : NRQCD with improved coefficients

# Overlap fermions on 2+1+1 Flavors HISQ Configurations

➤ Lattices used for this study :

HISQ gauge configurations from MILC

$24^3 \times 64$  ,  $a = 0.12$  fm,  $m_l/m_s = 1/5$ ,  $m_\pi L = 4.54$ ,  $m_\pi = 305$  MeV

$32^3 \times 96$  ,  $a = 0.089$  fm,  $m_l/m_s = 1/5$ ,  $m_\pi L = 4.5$ ,  $m_\pi = 312$  MeV

$48^3 \times 144$  ,  $a = 0.058$  fm,  $m_l/m_s = 1/5$ ,  $m_\pi L = 4.51$ ,  $m_\pi = 319$  MeV

PHYSICAL REVIEW D 87, 054505 (2013) (MILC)

➤ HYP smearing on gauge fields

➤ Both point source and coulomb gauge fixed wall source are used

➤ No of eigenvectors projected : 350 ( $a = 0.012$  fm)

: 350 ( $a = 0.09$  fm)

: 75 ( $a = 0.058$  fm)

# Overlap Fermions

## ➤ Some desirable features:

– No  $O(a)$  error.

$$(1 - \frac{1}{2}D)D(m)^{-1} = (D_c + ma)^{-1}$$

– The effective propagator :

$D_c = D/(1 - D/2)$  is chirally symmetric, i.e.,  $\{\gamma_5, D_c\} = 0$ .

–  $D_c + m$  is like in the continuum formalism.

– Multi-mass algorithm (more than 20 masses  
-10-15% overhead)

– Renormalization may be relatively simple (e.g. with chiral Ward identity).

## ➤ Undesirable feature:

-- Cost

# Rest mass Vs Kinetic mass

Charm mass is tuned by meson kinetic mass  
and not from rest mass  
.....a la FermiLab formulation

El-khadra et al,  
PRD55, 3933(1997)

Expanding the energy momentum relation in powers of  $pa$

$$E(p)^2 = M_1^2 + \frac{M_1}{M_2} \mathbf{p}^2 + O(\mathbf{p}^4)$$

$$= \mathbf{M}_1^2 + \mathbf{c}^2 \mathbf{p}^2$$

$$|\mathbf{p}| \ll m_0, 1/a$$

Rest mass :  $M_1 = E(\mathbf{0})$

Kinetic mass :  $\mathbf{M}_2 = \mathbf{M}_1/\mathbf{c}^2$

# NRQCD Action

$$aH = aH_0 + a\delta H; \quad aH_0 = -\frac{\Delta^{(2)}}{2am_b},$$

$$a\delta H = -c_1 \frac{(\Delta^{(2)})^2}{8(am_b)^3} + c_2 \frac{i}{8(am_b)^2} (\nabla \cdot \tilde{\mathbf{E}} - \tilde{\mathbf{E}} \cdot \nabla)$$

$$- c_3 \frac{1}{8(am_b)^2} \sigma \cdot (\tilde{\nabla} \times \tilde{\mathbf{E}} - \tilde{\mathbf{E}} \times \tilde{\nabla})$$

$$- c_4 \frac{1}{2am_b} \sigma \cdot \tilde{\mathbf{B}} + c_5 \frac{\Delta^{(4)}}{24am_b} - c_6 \frac{(\Delta^{(2)})^2}{16n(am_b)^2}.$$

$\nabla$  is the symmetric lattice derivative

$$\Delta^{(2)} \text{ and } \Delta^{(4)} \rightarrow \sum_i D_i^2 \text{ and } \sum_i D_i^4$$

**Tadpole improvement.**

$$H_0 \rightarrow \mathcal{O}(v^2)$$

**Coefficients are improved**  $\mathcal{O}(\alpha_s^2 a^2 v^4)$

$$\delta H \rightarrow \mathcal{O}(v^4).$$

**Dowdall et al, PRD 85,054509,2012**

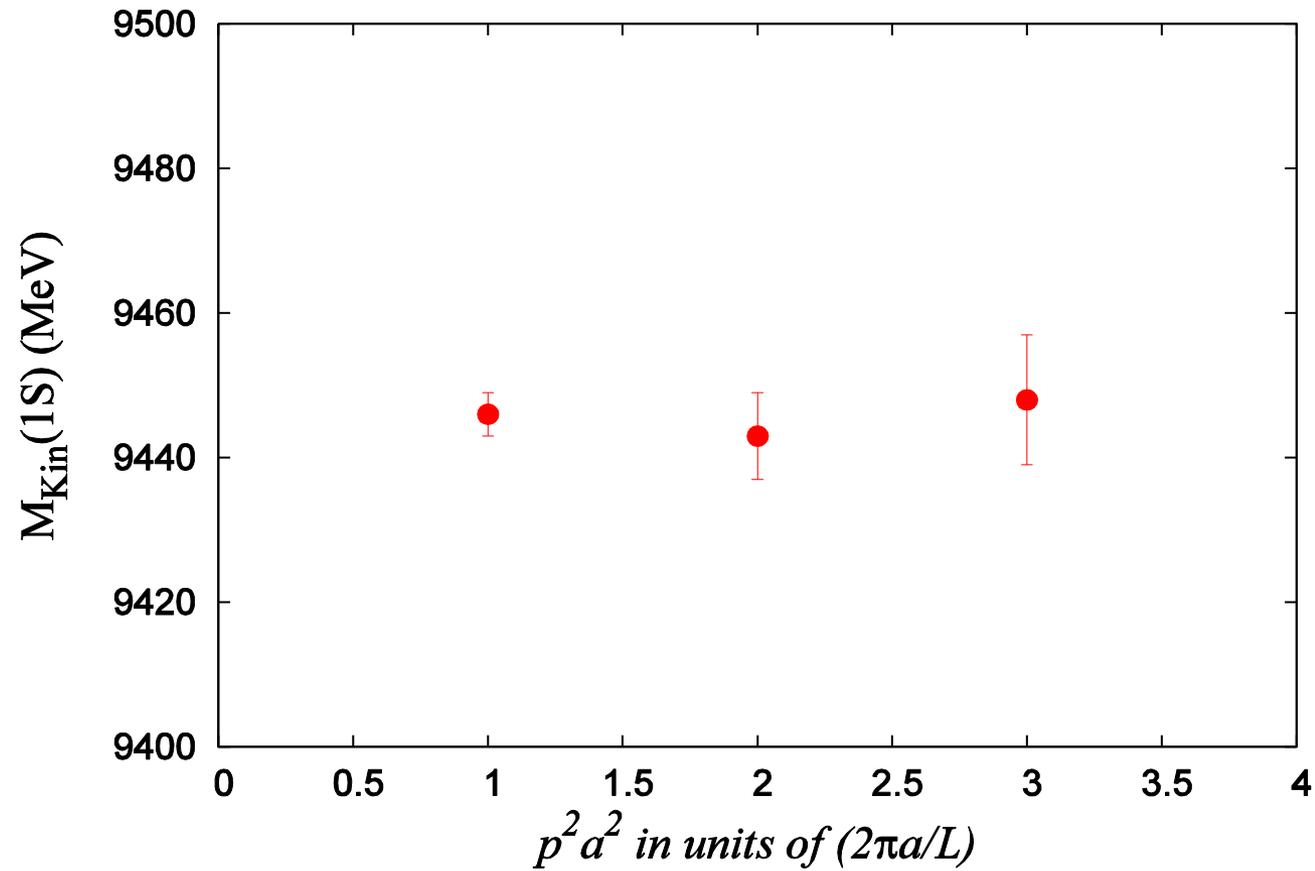
**A. Hart et al, PRD 79, 074008 (2009)**

$$aE(P) = aE(0) + \sqrt{a^2 P^2 + a^2 M_{\text{Kin}}^2}$$

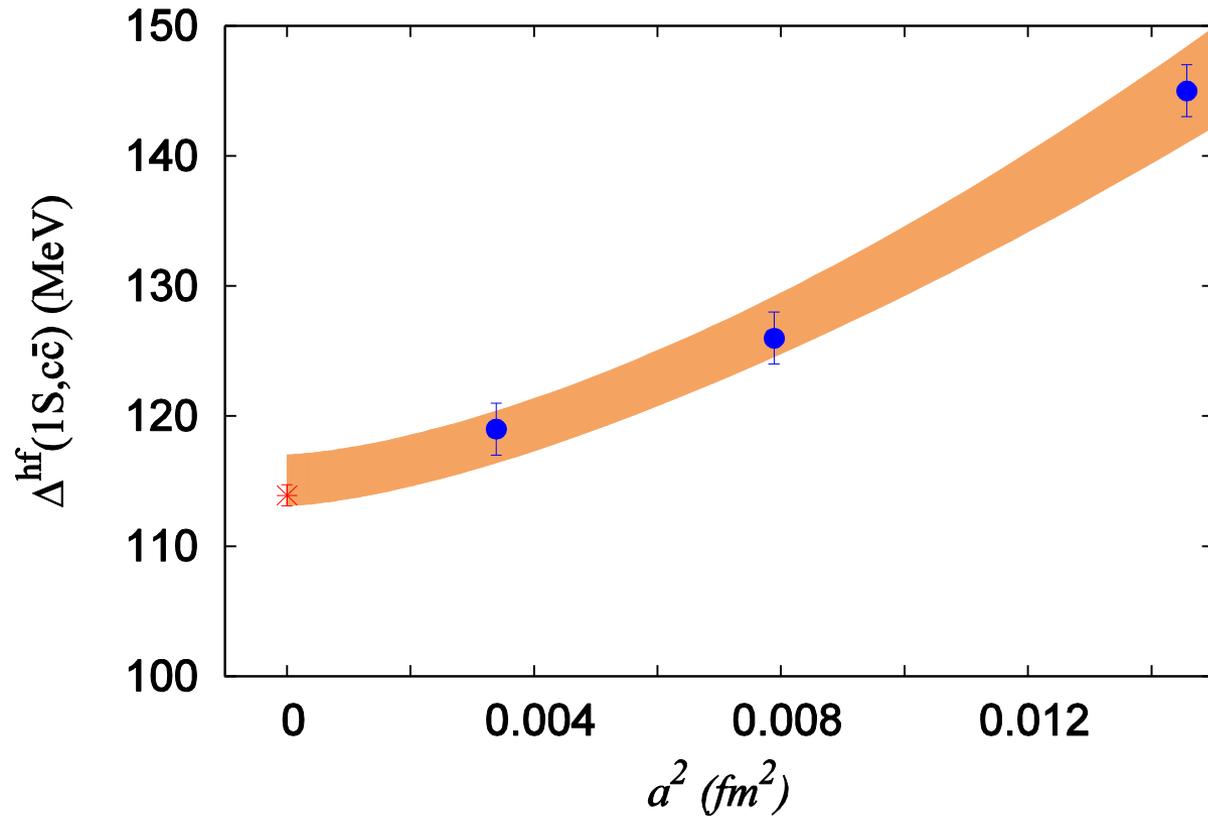
$$aM_{\text{Kin}} = \frac{a^2 P^2 - (a\Delta E)^2}{2a\Delta E}$$

$$\bar{M}_{\text{Kin}}(1S) = \frac{(3M_{\text{Kin}}(\Upsilon) + M_{\text{Kin}}(\eta_b))}{4}$$

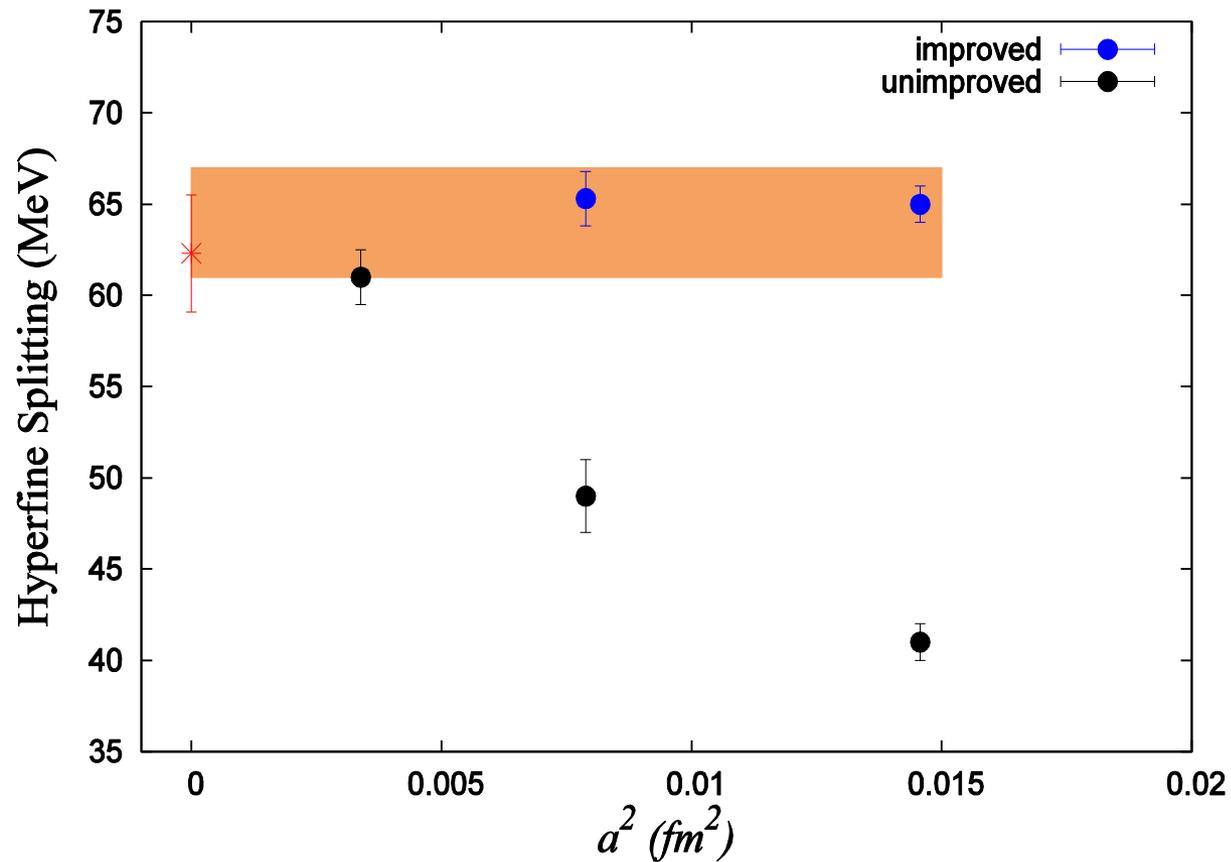
# b-quark mass tuning



# Hyperfine splitting in charmonia

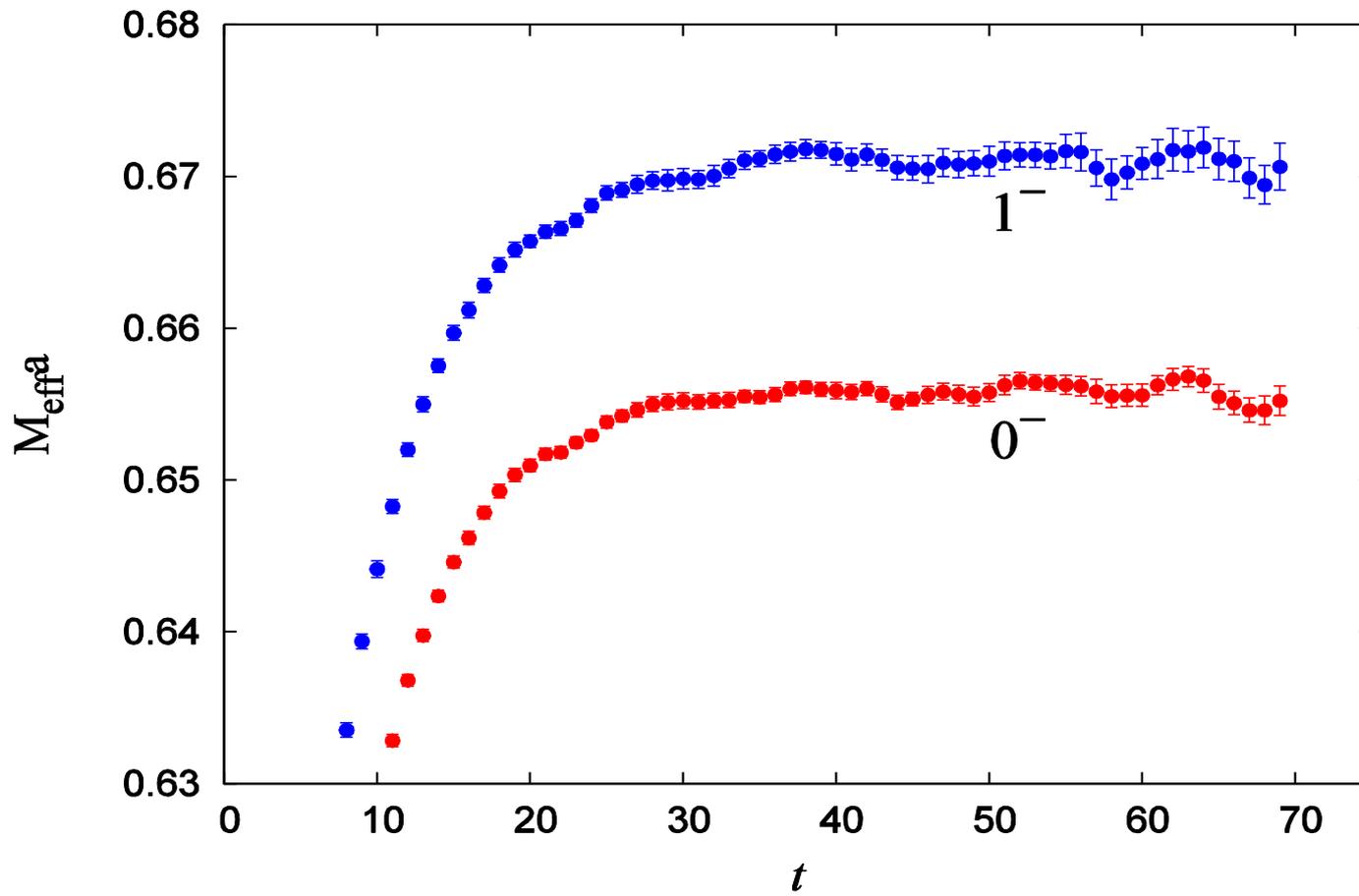


# Hyperfine splitting in bottomonia

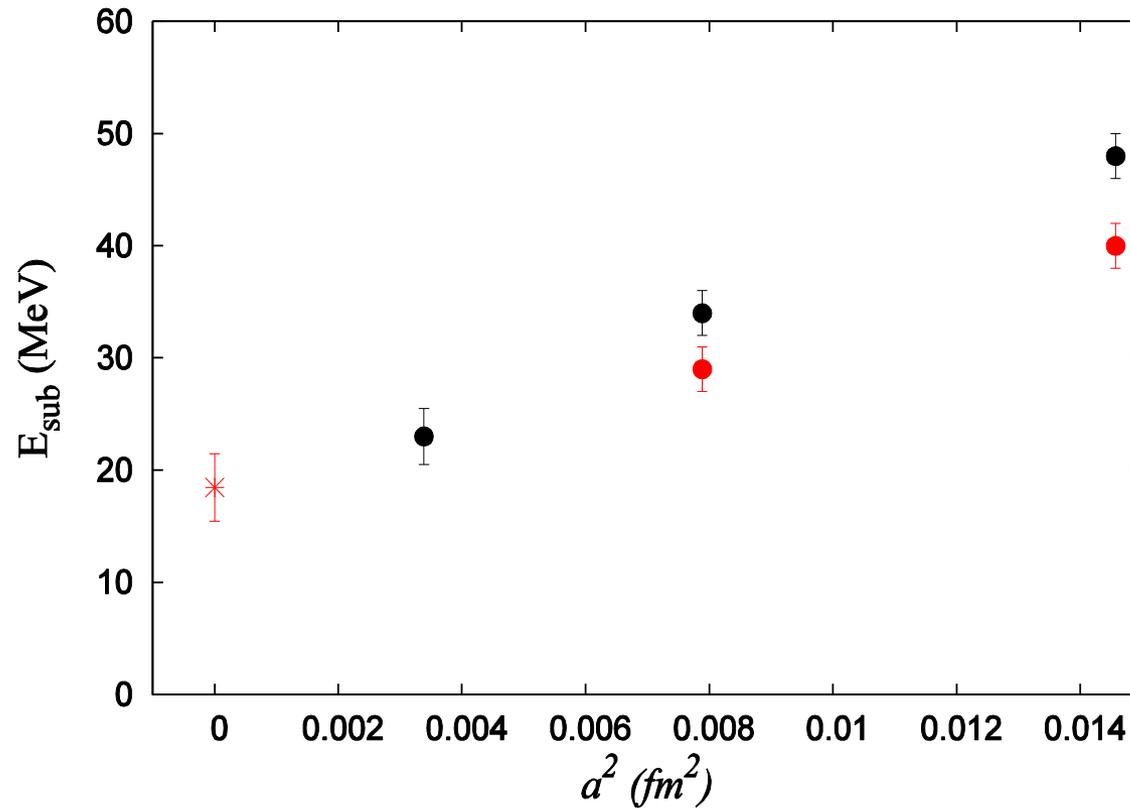


- We have calculated ground state energy spectra of mesons and baryons for all possible quantum numbers with  $l$ ,  $s$ ,  $c$  and  $b$  quarks (S.Mondal, talk Wednesday)
- Here we discuss decay constants and semileptonic form factors and present **preliminary** results.

# Effective masses for $B_c$ mesons

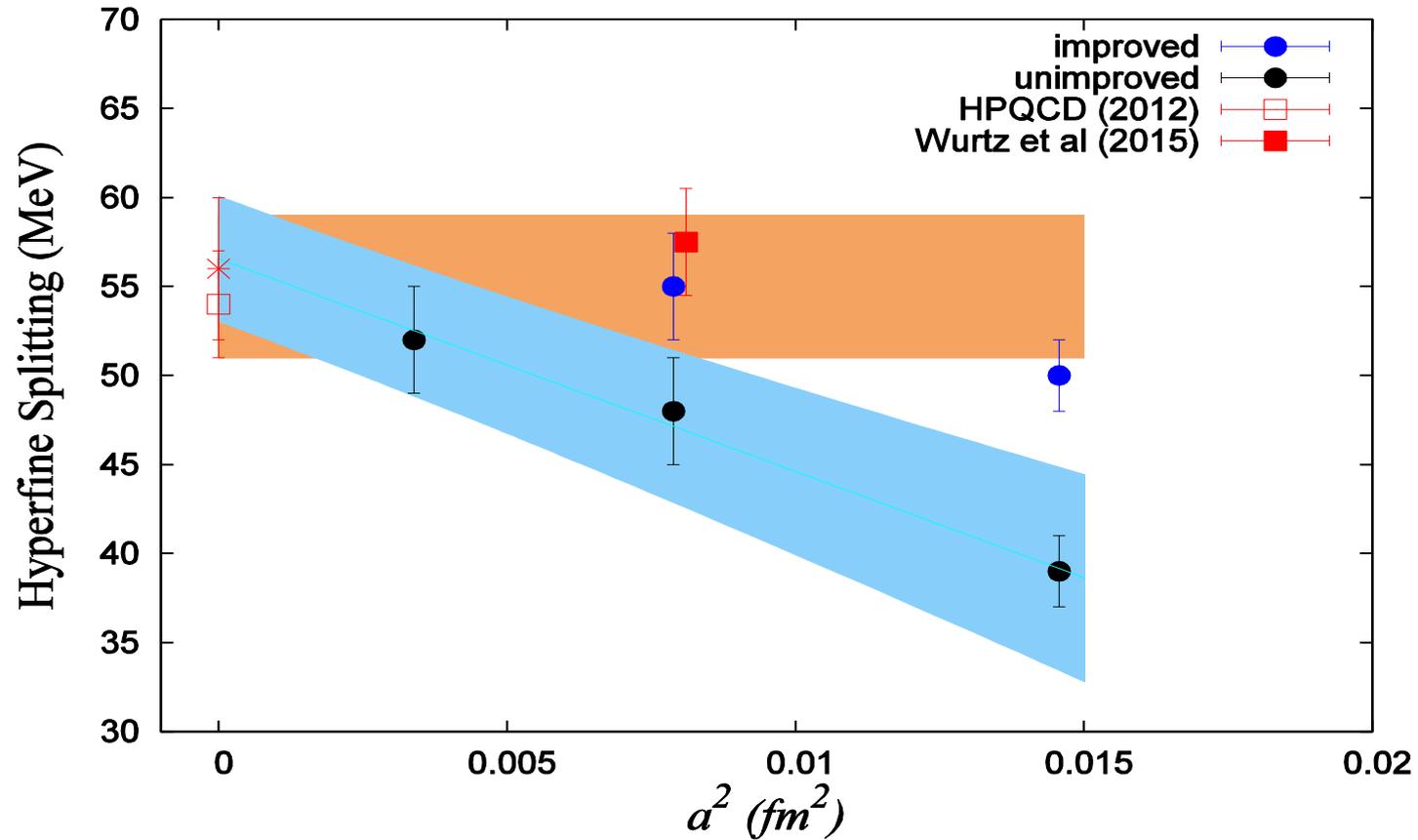


# $B_c(0^-)$



$$E_X^{(\text{sub})} = E_X - \frac{n_c}{2} \bar{E}_{c\bar{c}} - \frac{n_b}{2} \bar{E}_{b\bar{b}}$$

# Prediction : $B_c^*$



# Currents and decay constants

Colquhoun *et al* HPQCD : Phys. Rev. D91 (2015), 114509

**PS**

$$\langle 0 | J_{A_0} | B_q \rangle = f_{B_q} M_{B_q}$$

$$J_{A_0} = (1 + z_{A_0} \alpha_s) (J_{A_0}^{(0)} + J_{A_0}^{(1)})$$

$$J_{A_0}^{(0)} = \bar{\Psi}_q \gamma_5 \gamma_0 \Psi_b$$

$$J_{A_0}^{(1)} = -\frac{1}{2m_b} \bar{\Psi}_q \gamma_5 \gamma_0 \vec{\gamma} \cdot \vec{\nabla} \Psi_b$$

**V**

$$\langle 0 | J_{V_i} | B_q^j \rangle = f_{B_q^*} M_{B_q^*} \delta_{ij}$$

$$J_{V_i} = (1 + z_V \alpha_s) (J_{V_i}^{(0)} + J_{V_i}^{(1)})$$

$$J_{V_i}^{(0)} = \bar{\Psi}_q \gamma_i \Psi_b$$

$$J_{V_i}^{(1)} = -\frac{1}{2m_b} \bar{\Psi}_q \gamma_i \vec{\gamma} \cdot \vec{\nabla} \Psi_b$$

$$C_J^i(t) = \sum_i |a_J^i|^2 e^{-E_i t}$$

$$a_J^0 = \frac{\langle 0 | J | B_c \rangle}{\sqrt{2M_{B_c}}}$$

# Two point functions → Decay Constants

$$C_J^i(t) = \sum_i |a_J^i|^2 e^{-E_i t} \quad a_J^0 = \frac{\langle 0 | J | B_c \rangle}{\sqrt{2M_{B_c}}}$$

Colquhoun *et al*

**HPQCD :**

Phys. Rev. D91 (2015), 114509

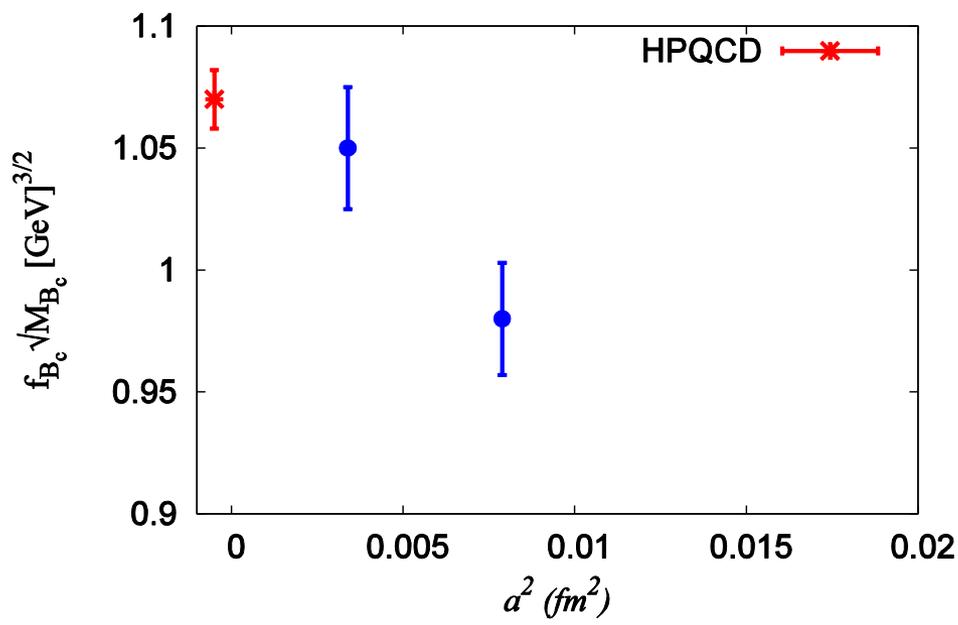
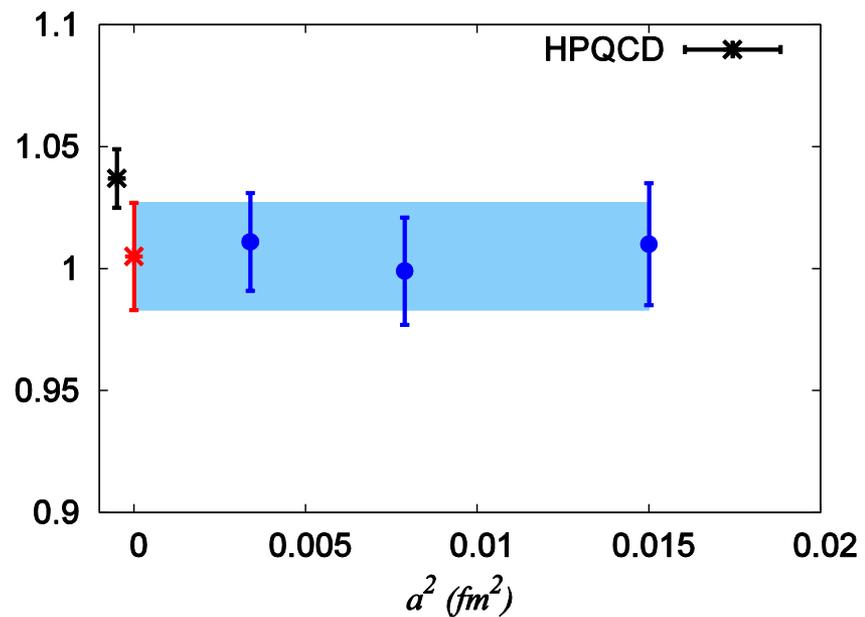
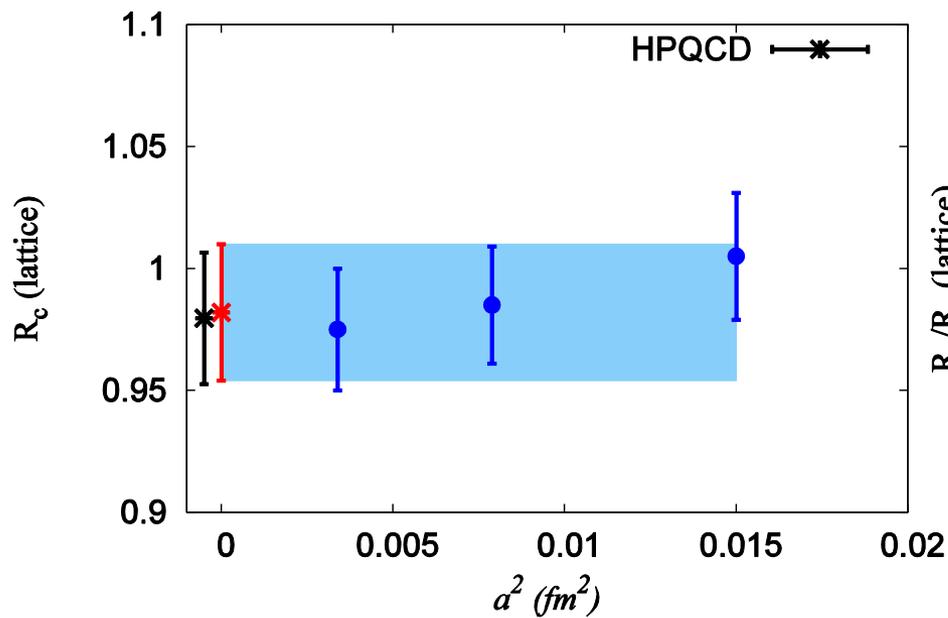
$$\Phi_J^{(0)} = \sqrt{2} a_{J(0)}^0 \quad \Phi_J^{(1)} = \sqrt{2} a_{J(1)}^0$$

$$f_{B_c} \sqrt{M_{B_c}} = (1 + \mathcal{Z}_{A_0} \alpha_s) \left( \Phi_{A_0}^{(0)} + \Phi_{A_0}^{(1)} \right)$$

$$f_{B_c^*} \sqrt{M_{B_c^*}} = (1 + \mathcal{Z}_V \alpha_s) \left( \Phi_{V_i}^{(0)} + \Phi_{V_i}^{(1)} \right)$$

$$R_c \equiv R_{B_c} \equiv \frac{f_{B_c^*} \sqrt{M_{B_c^*}}}{f_{B_c} \sqrt{M_{B_c}}} = (1 + \delta \mathcal{Z} \alpha_s) \frac{(\Phi_{V_i}^{(0)} + \Phi_{V_i}^{(1)})}{(\Phi_{A_0}^{(0)} + \Phi_{A_0}^{(1)})}$$

$$\frac{R_c}{R_s} = (1 + [\delta \mathcal{Z}_c - \delta \mathcal{Z}_s] \alpha_s) \frac{(\Phi_{B_c^*}^{(0)} + \Phi_{B_c^*}^{(1)})}{(\Phi_{B_c}^{(0)} + \Phi_{B_c}^{(1)})} \frac{(\Phi_{B_s^*}^{(0)} + \Phi_{B_s^*}^{(1)})}{(\Phi_{B_s}^{(0)} + \Phi_{B_s}^{(1)})}$$



# Form factors

Colquhoun *et al* **HPQCD : 1611.01987**

A. Lytle : CKM2016

## ■ $B_c \rightarrow \eta_c l \nu$

$$\langle \eta_c(p) | V^\mu | B_c(P) \rangle = f_+(q^2) \left[ P^\mu + p^\mu - \frac{M^2 - m^2}{q^2} q^\mu \right] + f_0(q^2) \frac{M^2 - m^2}{q^2} q^\mu$$

## ■ $B_c \rightarrow J/\psi l \nu$

$$\begin{aligned} \langle J/\psi(p, \varepsilon) | V^\mu - A^\mu | B_c(P) \rangle = & \frac{2i\varepsilon^{\mu\nu\rho\sigma}}{M+m} \varepsilon_\nu^* p_\rho P_\sigma V(q^2) - (M+m) \varepsilon^{*\mu} A_1(q^2) + \\ & \frac{\varepsilon^* \cdot q}{M+m} (p+P)^\mu A_2(q^2) + 2m \frac{\varepsilon^* \cdot q}{q^2} q^\mu A_3(q^2) - 2m \frac{\varepsilon^* \cdot q}{q^2} q^\mu A_0(q^2) \end{aligned}$$

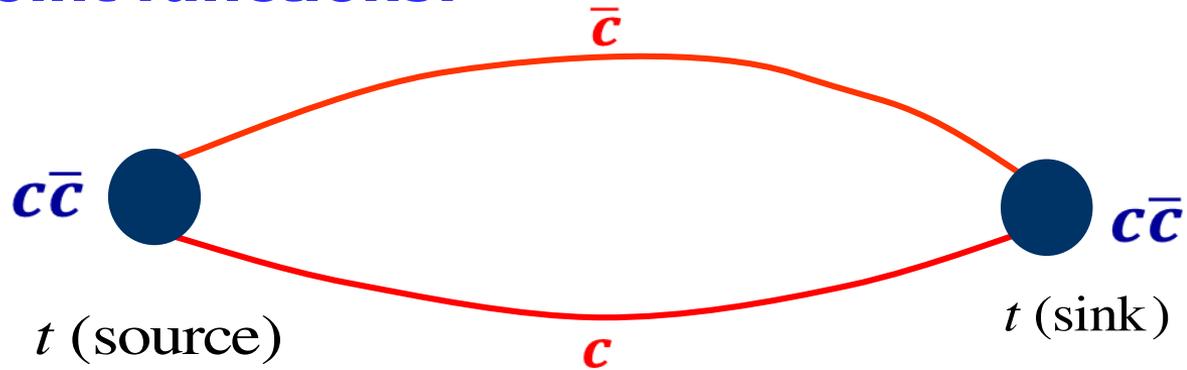
$$q = P - p$$

$q_{\max}^2$  : **Outgoing hadron at rest**

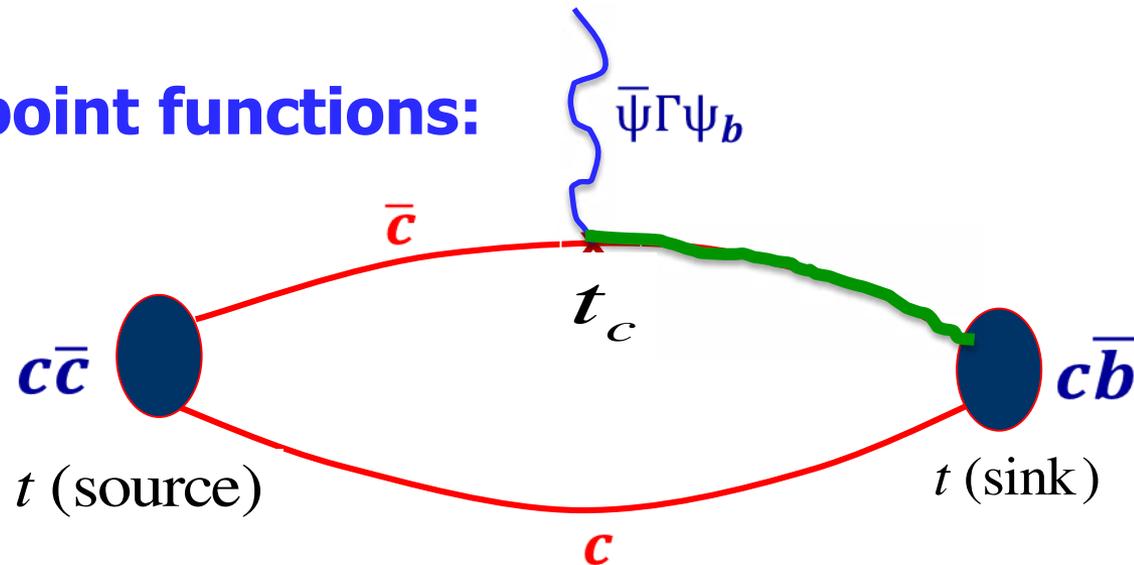
$q^2 = 0$  : **Maximum recoil**

# Semileptonic Decay

□ Two point functions:



□ Three point functions:



# Two and three point functions

$$C_2^{\eta_c}(t) = \sum_i (A_{\eta_c}^i)^2 e^{-E_{\eta_c}^i t}$$

$$C_2^{B_c}(\tau) = \sum_i (A_{B_c}^i)^2 e^{-E_{B_c}^i \tau}$$

$$C_{3,m}^{B_c \rightarrow \eta_c}(t, \tau) = \sum_{i,j} A_{\eta_c}^i \varphi^m A_{B_c}^j e^{-E_{\eta_c}^i t} e^{-E_{B_c}^j \tau}$$

- $\varphi^m$ : Can be obtained by
- fitting these two and three point function simultaneously
  - constructing appropriate ratios

# Ratios

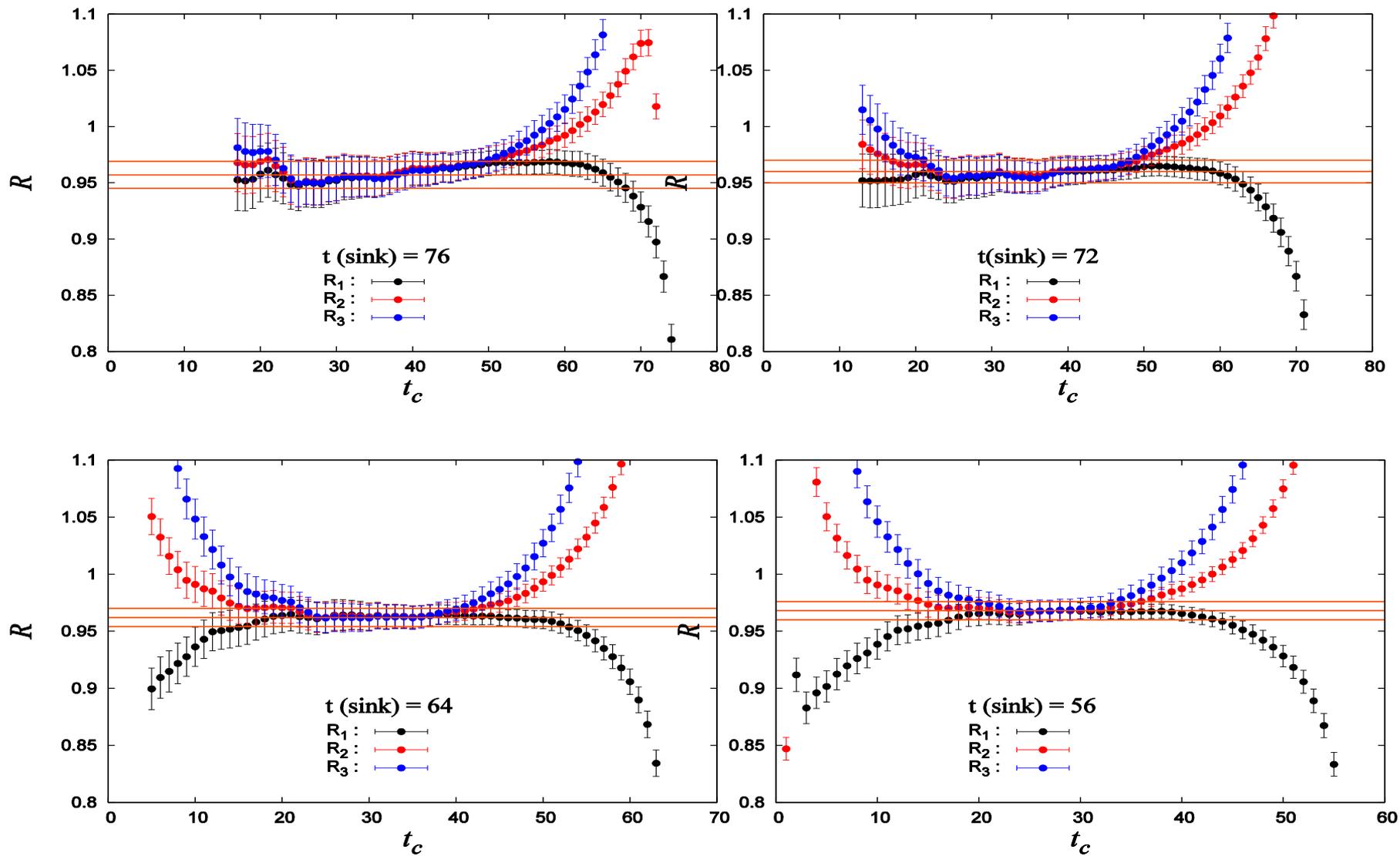
$$R_1(t_i, \tau) = \frac{C_3(t, \tau)}{[C_2^{\eta_c}(t_i) C_2^{B_c}(t_i)]^{1/2}} e^{E_{\eta_c}(t-t_i/2)} e^{E_{B_c}(\tau-t_i/2)}$$

$$R_2(t, \tau) = \frac{C_3(t, \tau)}{[C_2^{\eta_c}(t) C_2^{B_c}(\tau)]^{1/2}} e^{E_{B_c}\tau/2} e^{E_{\eta_c}t/2}$$

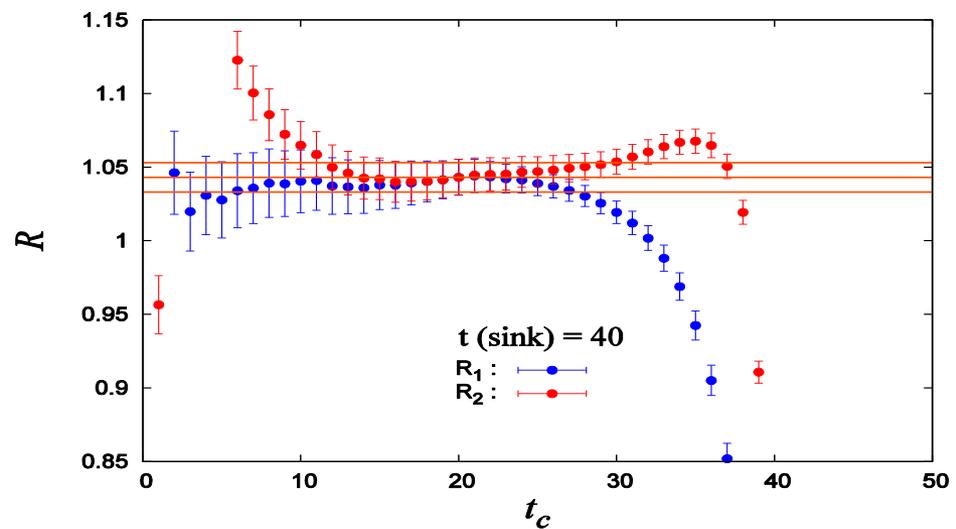
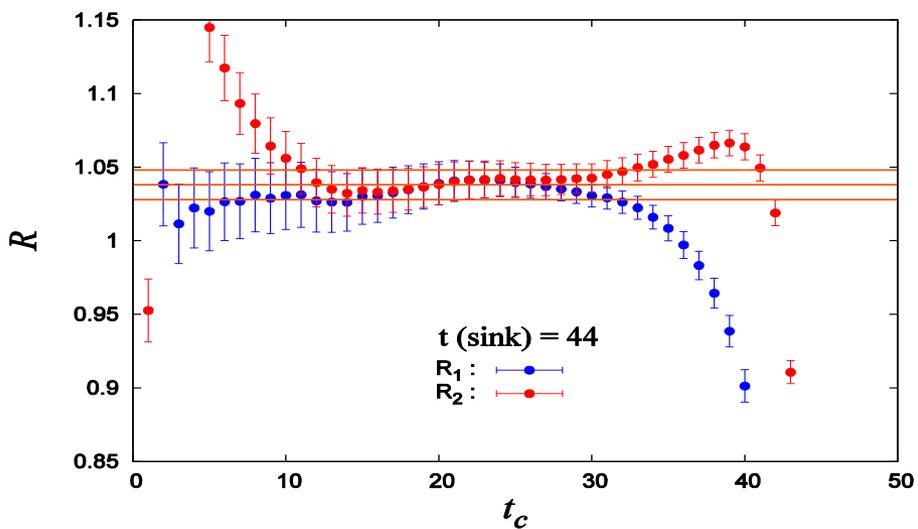
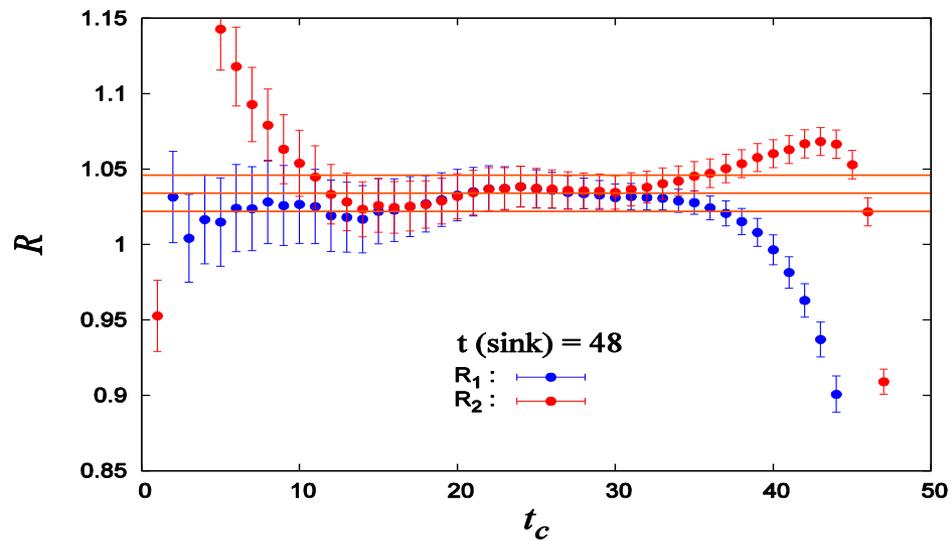
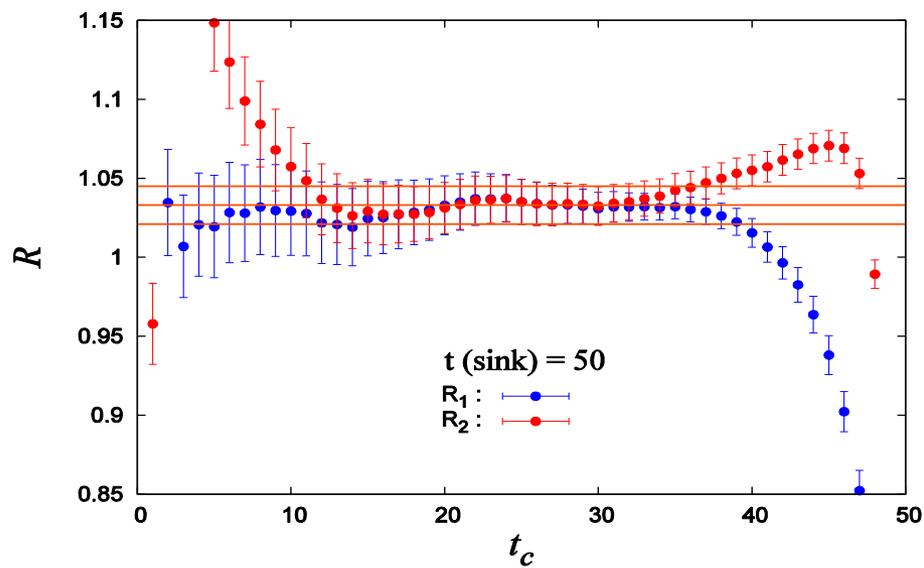
$$R_3(t, \tau) = \frac{C_3(t, \tau)}{\frac{1}{A_{B_c} A_{\eta_c}} C_2^{\eta_c}(t) C_2^{B_c}(\tau)}$$

Bahr et al : 1701.03299

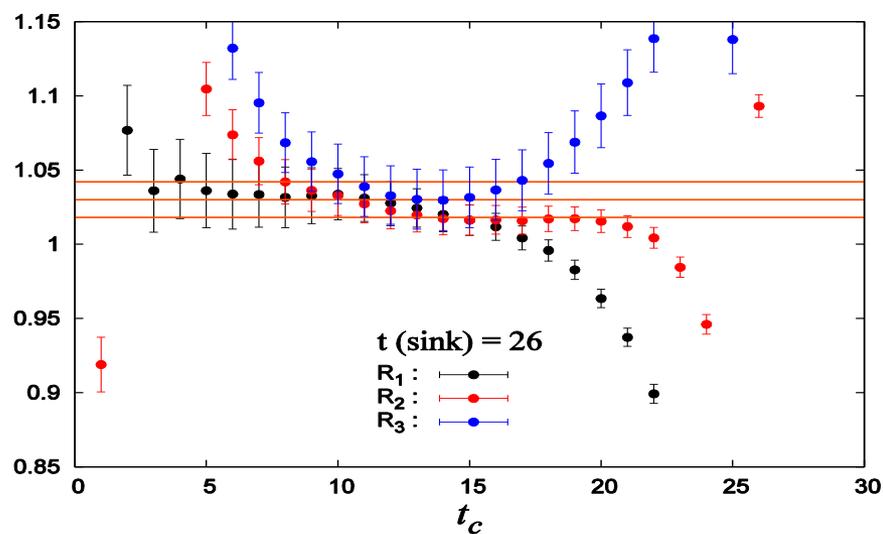
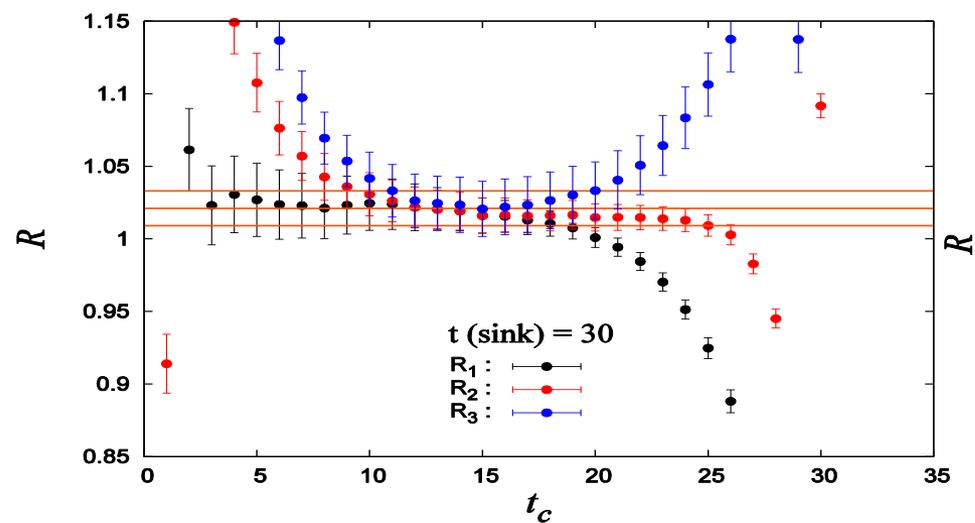
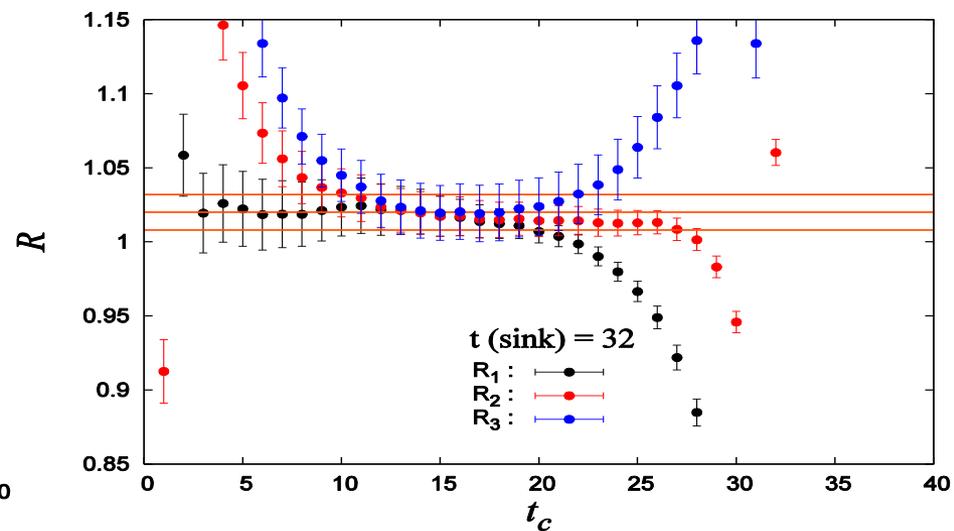
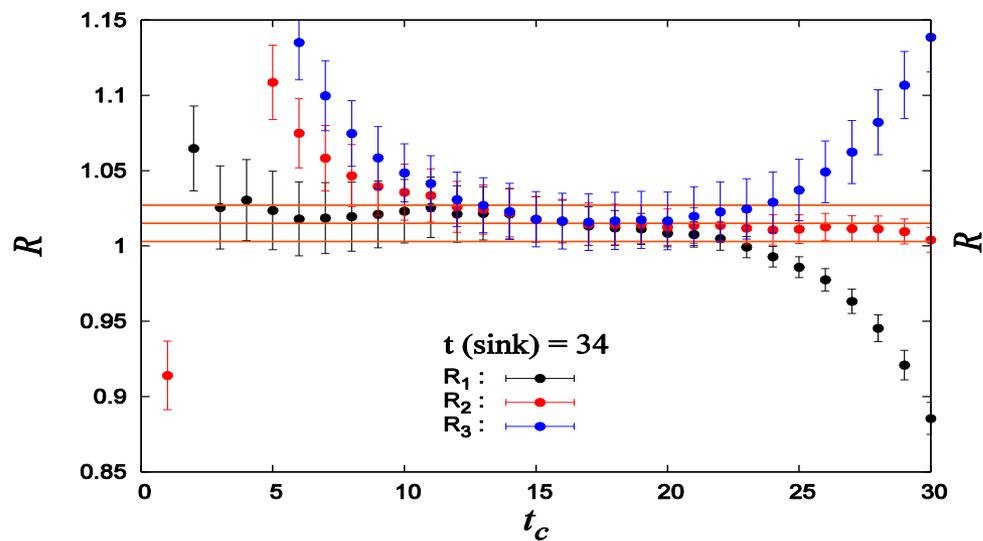
# $48^3 \times 144$ ( $a \sim 0.06 \text{ fm}$ )

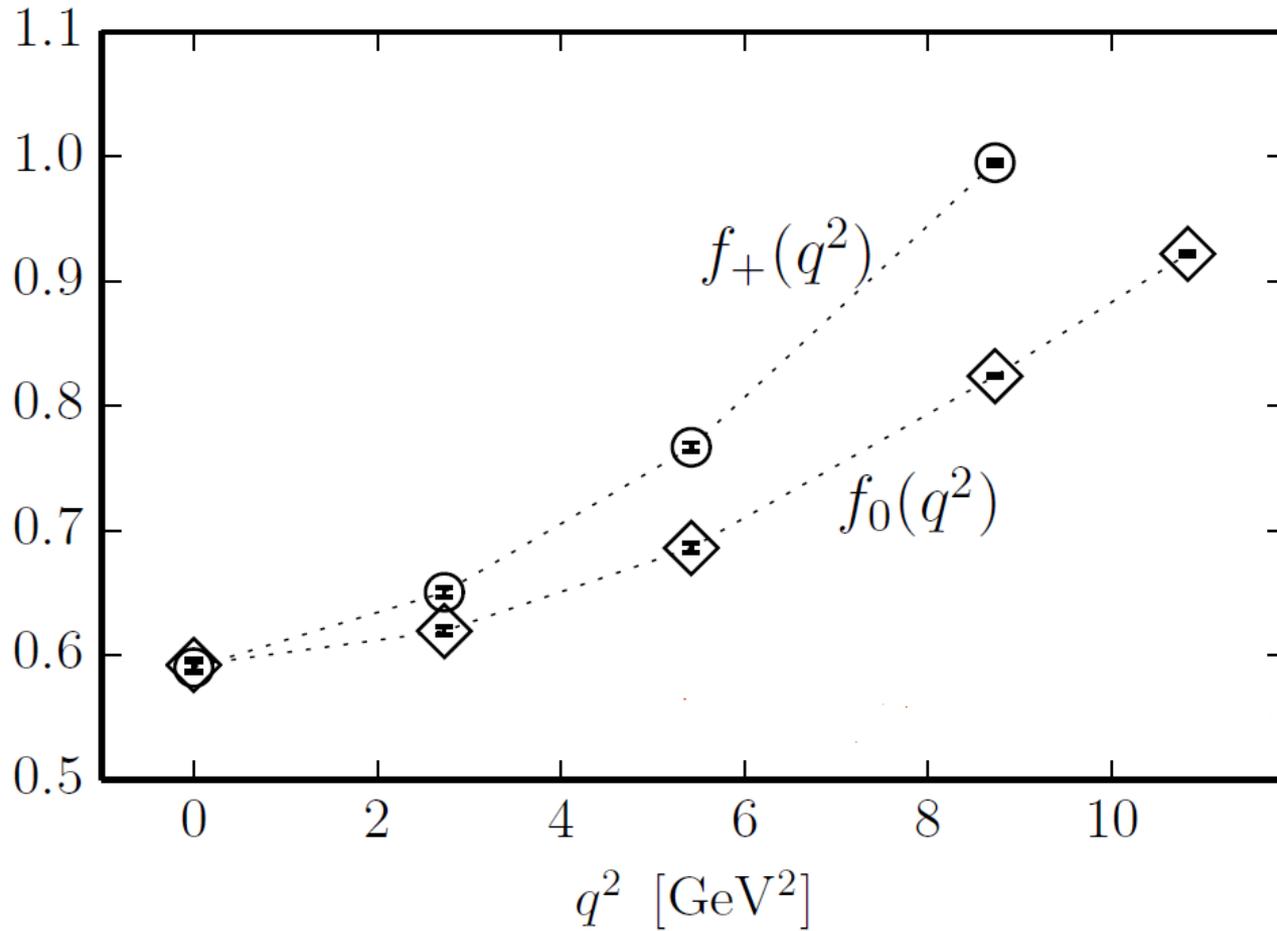


# $36^3 \times 96$ ( $a \sim 0.09 \text{ fm}$ )



# 24<sup>3</sup>x64 ( $a \sim 0.12\text{fm}$ )

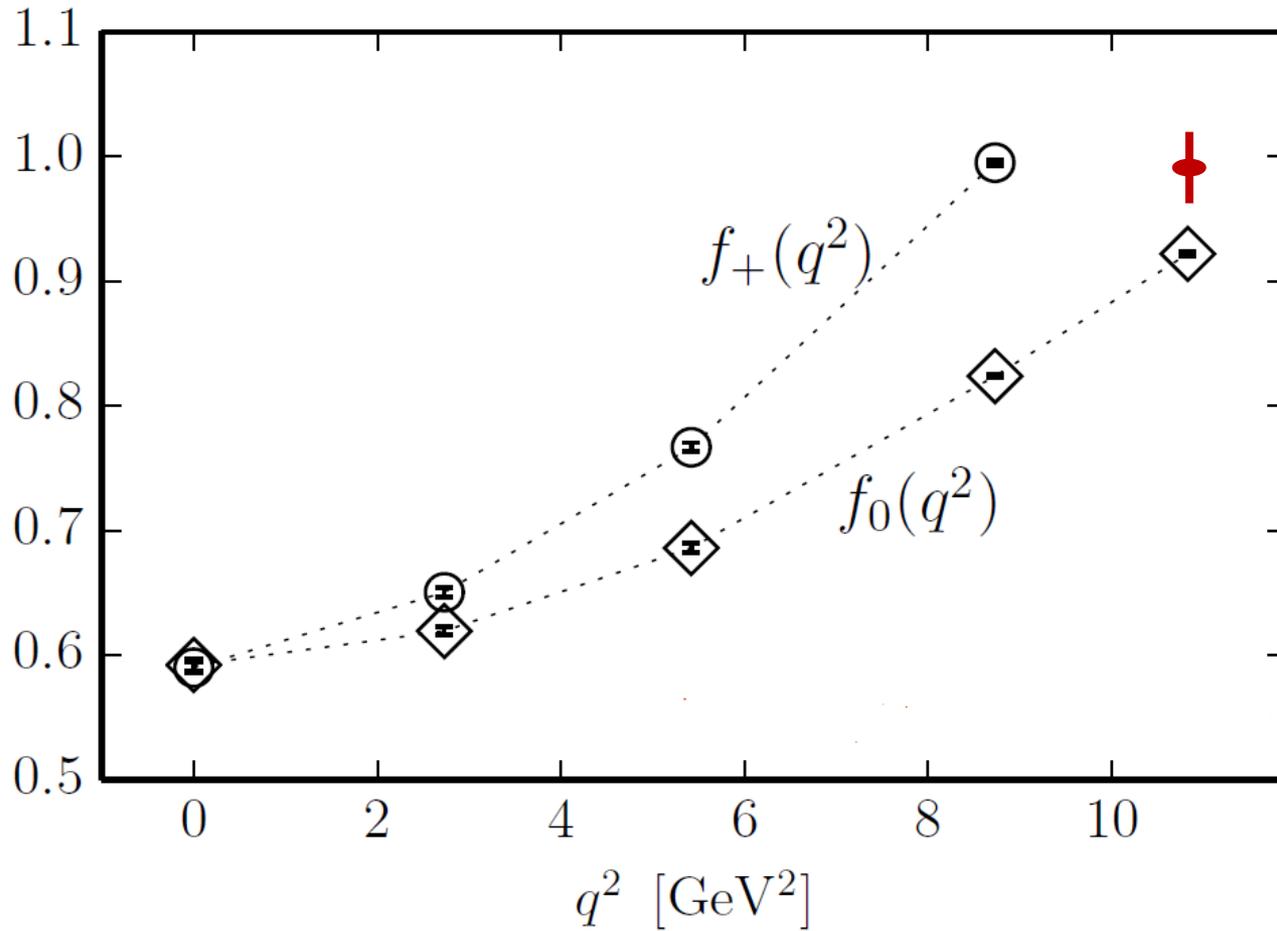




**Colquhoun et al HPQCD : 1611.01987**

**A. Lytle : CKM2016**

**A. Lytle : Lattice2017**



Colquhoun *et al* **HPQCD : 1611.01987**

A. Lytle : CKM2016

# Conclusions and Outlooks

- ✚ Study of decay constants and semileptonic decay of  $B_c$  mesons provides opportunity to obtain information about  $b \rightarrow c$  transition, in particular important quantities like :  $|V_{cb}|$  and  $R(B_c)$ .
- ✚ We have started to study these using NRQCD ( $b$ ) and overlap ( $l$ ,  $s$ , and  $c$ ) valence quarks on HISQ gauge configurations.
- ✚ Preliminary results are presented here.
- ✚ Need to calculate renormalization constants to obtain physical results
- ✚ Use of fully relativistic bottom quarks on finer lattices will also be followed.
- ✚ Acknowledgements : ILGTI, MILC (S. Gottlieb), R. Lewis, B. Colquhoun, A. Lytle and G. Bali

# Lattice spacings and tuning of charm and strange masses

Lattice spacings are calculated by  $\Omega(sss)$  mass = 1672 GeV

$48^3 \times 144$  : 0.0582(5) fm

$32^3 \times 96$  : 0.0877(10) fm

$24^3 \times 64$  : 0.1192(14) fm

which are quite consistent with lattice spacings determined by MILC

- Strange mass is tuned by setting pseudoscalar  $\underline{ss}$  mass at 685 MeV

$$\begin{aligned} m_s a &= 0.0738 \quad (a = 0.0118\text{fm}) \\ &= 0.0485 \quad (a = 0.0888\text{ fm}) \\ &= 0.028 \quad (a = 0.0582\text{fm}) \end{aligned}$$

Taking  $m_s = 100$  MeV

$$\begin{aligned} m_s a &= 0.0450 \quad (a = 0.0888\text{fm}), \\ &= 0.0295 \quad (a = 0.0582\text{fm}) \end{aligned}$$

- Charm mass is tuned by  $\frac{1}{4} (m_{\eta_c} + 3 m_{J/\psi})$

$$\begin{aligned} m_c a &= 0.528 \quad (a = 0.1192\text{fm}) \\ &= 0.428 \quad (a = 0.0888\text{fm}), \\ &= 0.29 \quad (a = 0.0582\text{ fm}) \end{aligned}$$

Considering kinetic masses of mesons  
(a la Fermilab formulation)