

# Topological Susceptibility in $N_f=2$ QCD at Finite Temperature

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JLQCD collaboration:

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Kei Suzuki



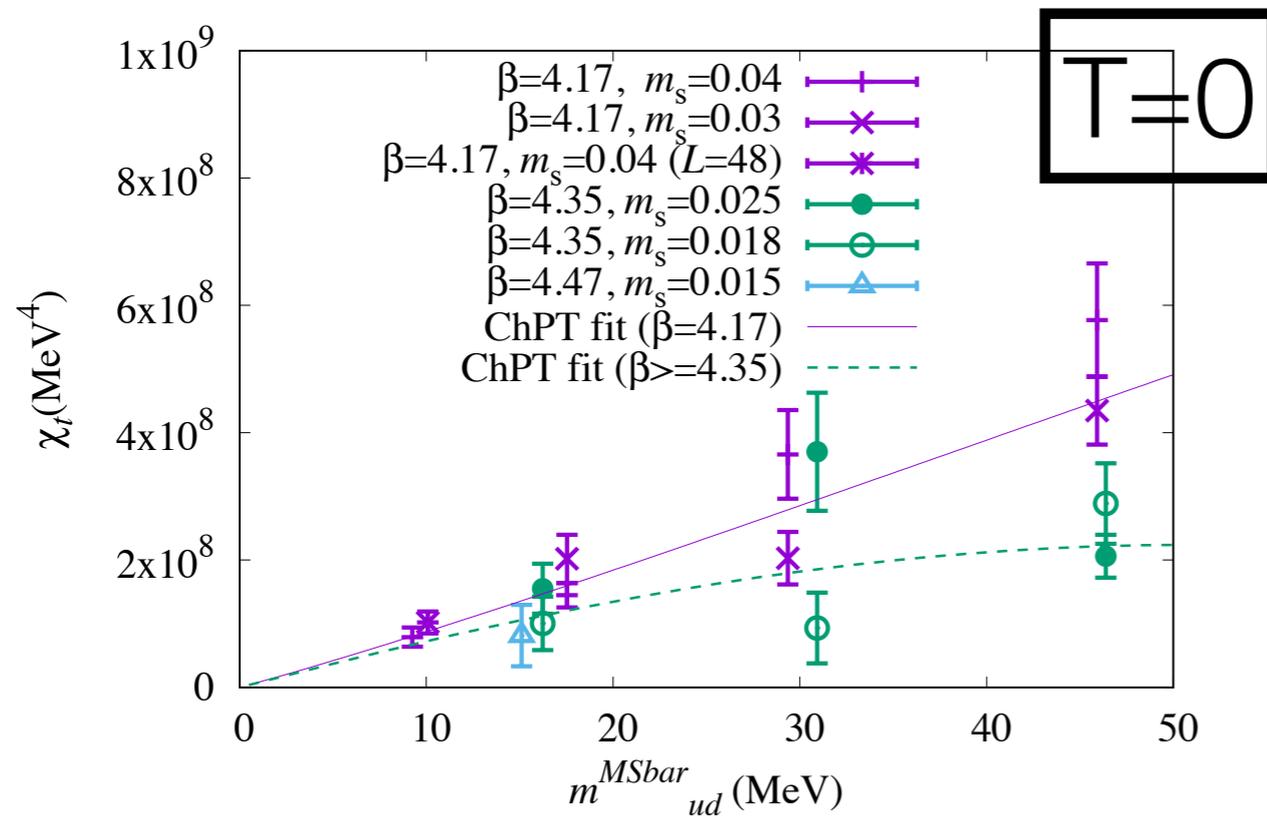
Lattice 2017 @ Granada  
June 19, 2017

# JLQCD finite temperature related

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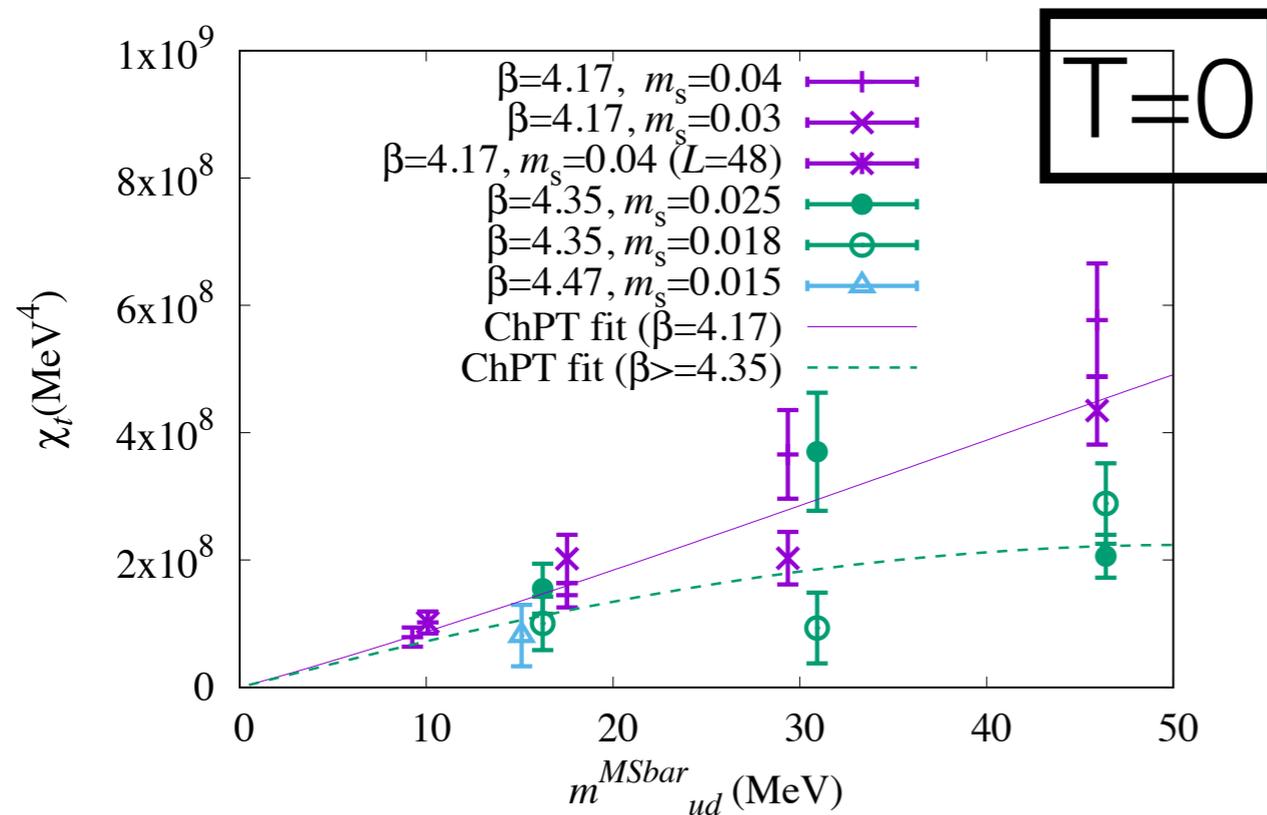
- Kei Suzuki:  $U_A(1)$
- YA: this talk about the topological susceptibility
- Hidenori Fukaya
  - plenary panel discussion on  $U_A(1)$
  - $\chi_t$  with slab method for  $T=0$  (parallel talk)
- Christian Rohrhofer: vector channel correlators

# topological susceptibility



[JLQCD: S.Aoki et al 2017,  $N_f=2+1$  DWF  
see talk by Fukaya]

# topological susceptibility



[JLQCD: S.Aoki et al 2017,  $N_f=2+1$  DWF  
see talk by Fukaya]

- $T > 0$  ?
  - $N_f=2$ 
    - relation with  $U(1)_A$  symm
    - relation with the order of PT
      - 2nd or 1st order ?
  - implication
    - axion :  $\chi_t(T)$  @ phys.pt.
    - $N_f=2+1$  phase diagram

# $U(1)_A$ and topology

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- global symmetry of QCD:  $U(N_f)_L \times U(N_f)_R$  @  $m_f \rightarrow 0$
- $T < T_c$ :  $\rightarrow SU(N_f)_V \times U(1)_V$
- $T > T_c$ :  $\rightarrow SU(N_f)_V \times U(1)_V \times SU(N_f)_A \times U(1)_A$  ?
  - $SU(N_f)_A$  is restored,  $U(1)_A$ ?
  - able to be investigated through order parameters
- relation through fermionic (near) zero mode:  $\rho(\lambda)$  : density of eigenvalue

$SU(N_f)_A$  order parameter

$$-\langle \bar{q}q \rangle = \pi \rho(0) = 0$$

$U(1)_A$  order parameter

$$\Delta_{\pi-\delta} \sim \rho'(0) = 0?$$



# $U(1)_A$ and topology

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[K.Suzuki]

# U(1)<sub>A</sub> and topology

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$$\begin{array}{ccc}
 \text{SU}(N_f)_A \text{ order parameter} & & \text{U}(1)_A \text{ order parameter} \\
 -\langle \bar{q}q \rangle = \pi \rho(0) = 0 & \longleftrightarrow & \Delta_{\pi-\delta} \sim \rho'(0) = 0 \quad [\text{K.Suzuki}] \\
 \swarrow \text{index theorem} & & \searrow \text{index theorem} \\
 \text{topological charge} & Q_t & \\
 \text{susceptibility} & \chi_t = \frac{\langle Q_t^2 \rangle}{V} & 
 \end{array}$$

# U(1)<sub>A</sub> and topology

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•  $T < T_c$ :  $\rightarrow SU(N_f)_V \times U(1)_V$

•  $T >$  signal should appear in topology

• able to be investigated through order parameters

• relation through fermionic (near) zero mode:  $\rho(\lambda)$  : density of eigenvalue

SU(N<sub>f</sub>)<sub>A</sub> order parameter

$$-\langle \bar{q}q \rangle = \pi \rho(0) = 0$$

topological charge

susceptibility

U(1)<sub>A</sub> order parameter

$$\Delta_{\pi-\delta} \sim \rho'(0) = 0$$

[K.Suzuki]

index theorem

$$Q_t$$

$$\chi_t = \frac{\langle Q_t^2 \rangle}{V}$$

# U(1)<sub>A</sub> and topology

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•  $T < T_c$ :  $\rightarrow SU(N_f)_V \times U(1)_V$

•  $T >$  signal should appear in topology

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• relative a scenario:  $\chi_t = 0$  for  $0 < m_f < m_c$  eigenvalue

[S.Aoki, Fukaya, Taniguchi (2012)]

$$-\langle \bar{q}q \rangle = \pi \rho(0) = 0 \longleftrightarrow \Delta_{\pi-\delta} \sim \rho'(0) = 0 \quad [\text{K.Suzuki}]$$

index theorem

topological charge

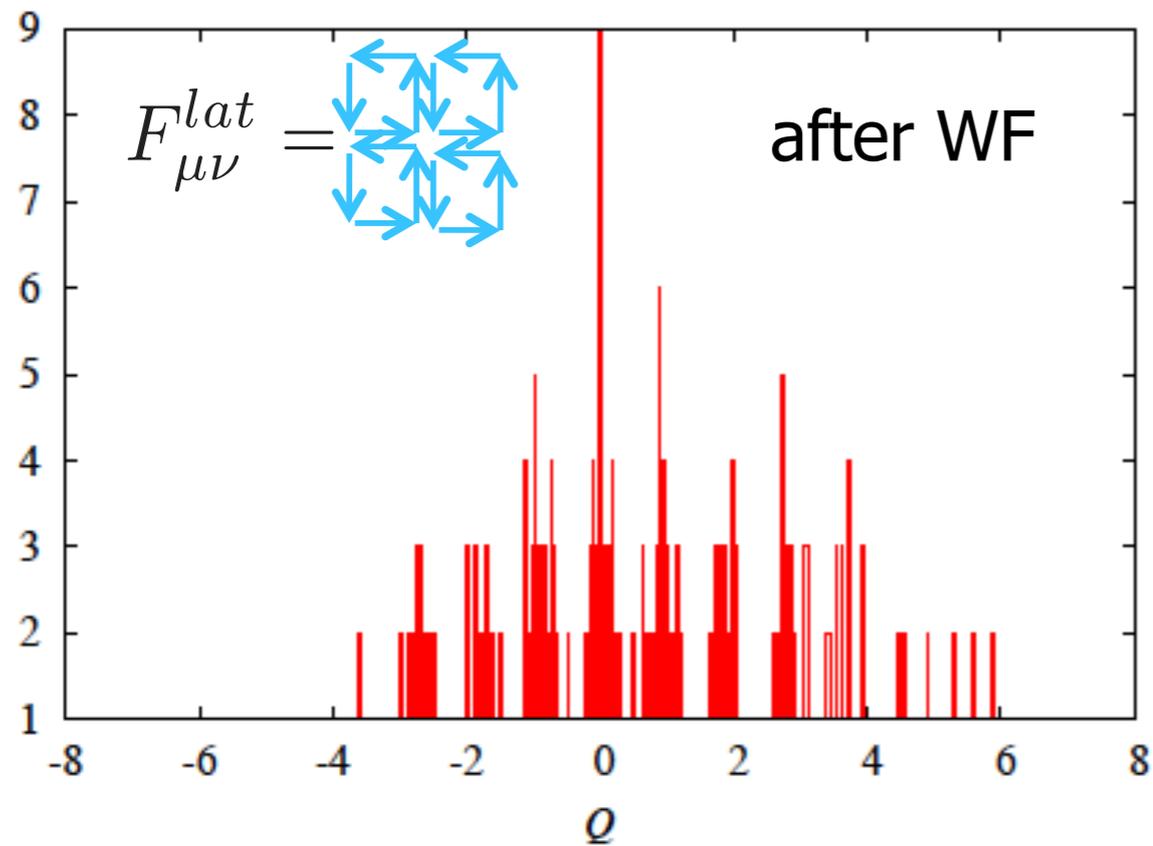
$$Q_t$$

susceptibility

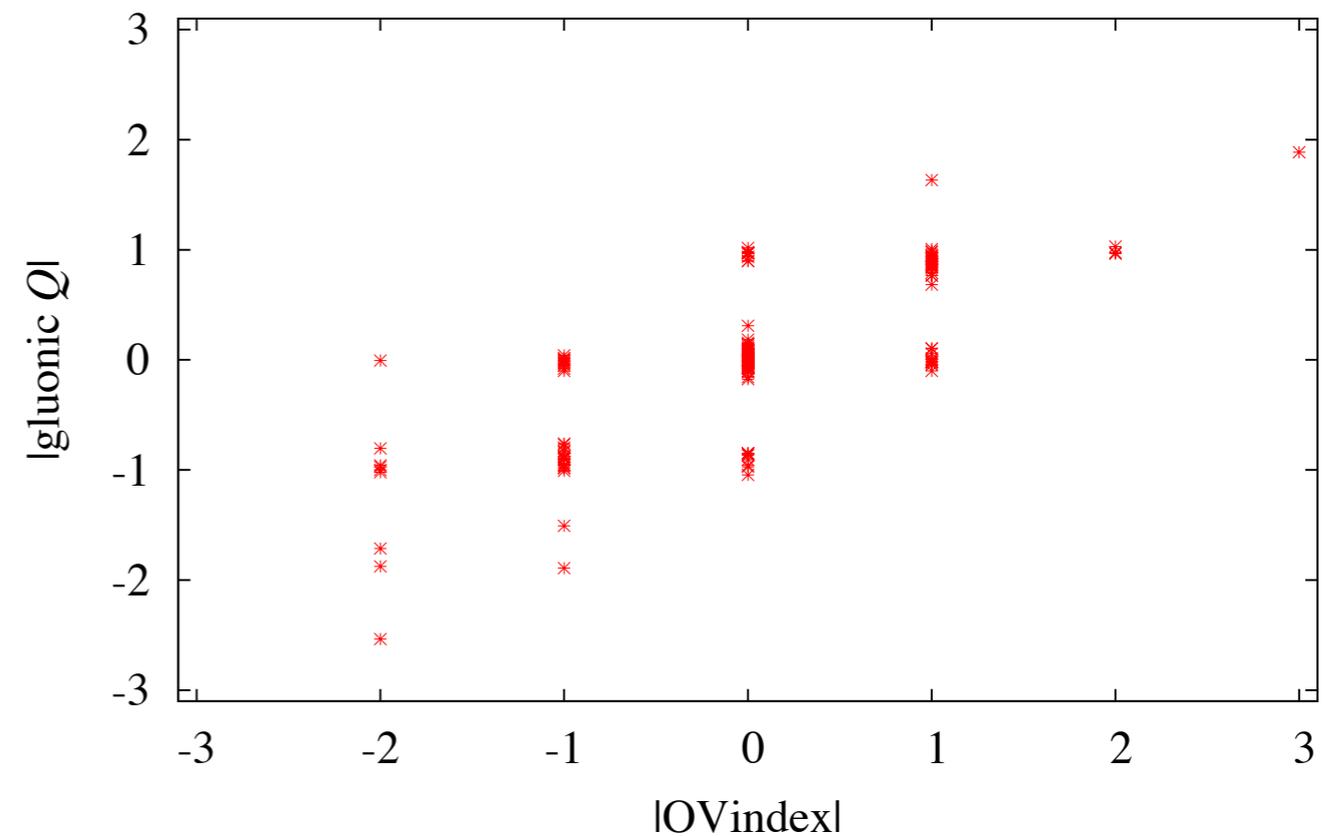
$$\chi_t = \frac{\langle Q_t^2 \rangle}{V}$$

# topological charge $\rightarrow$ susceptibility

$$q(x) = \frac{1}{32\pi^2} \text{Tr} \epsilon^{\mu\nu\rho\sigma} F_{\mu\nu}^{lat} F_{\rho\sigma}^{lat}$$



32x12 beta=4.23 m=0.005 (T=191 MeV)



$$Q = \sum_x q(x) \quad \text{charge}$$

$$\chi_t = \frac{\langle Q^2 \rangle}{V} \quad \text{susceptibility}$$

# $N_f=2$ DWF $\rightarrow$ Overlap

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- DWF ensemble  $\rightarrow$  reweighted to overlap
  - Möbius DWF: almost exact chiral symmetry:  $m_{\text{res}} = 0.05(3)$  MeV ( $\beta=4.3$ ,  $L_s=16$ )
  - Overlap: exact chiral symmetry
- $Q_t$  measurements
  - global sum of the gluonic charge density (clover) after Wilson Flow
  - Overlap Index
- reweighting: before / after and above 2 meas. yield 4  $\chi_t$  values
- current focus:  $1/a = 2.6$  GeV \*\*\* **PRELIMINARY** \*\*\*

# $N_f=2$ DWF $\rightarrow$ Overlap

- DWF ensemble
  - Möbius D
  - Overlap:
- $Q_t$  measurement
  - global sum
  - Overlap Index

GL-DW	gluonic charge on DW ensemble
GL-OV	gluonic charge on OV ensemble
OV-DW	OV index on DW ensemble
OV-OV	OV index on OV ensemble

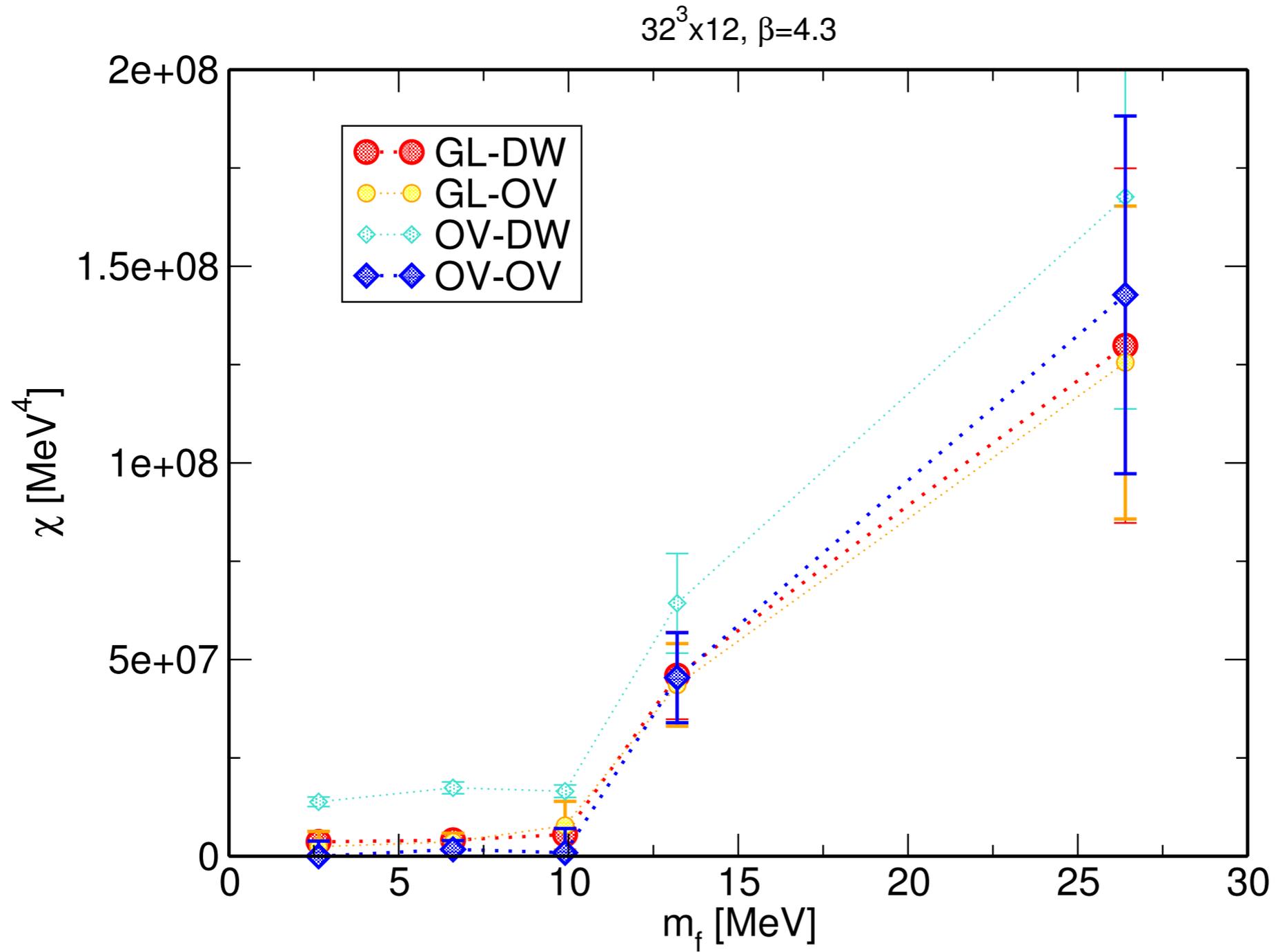
V ( $\beta=4.3, L_s=16$ )

Flow

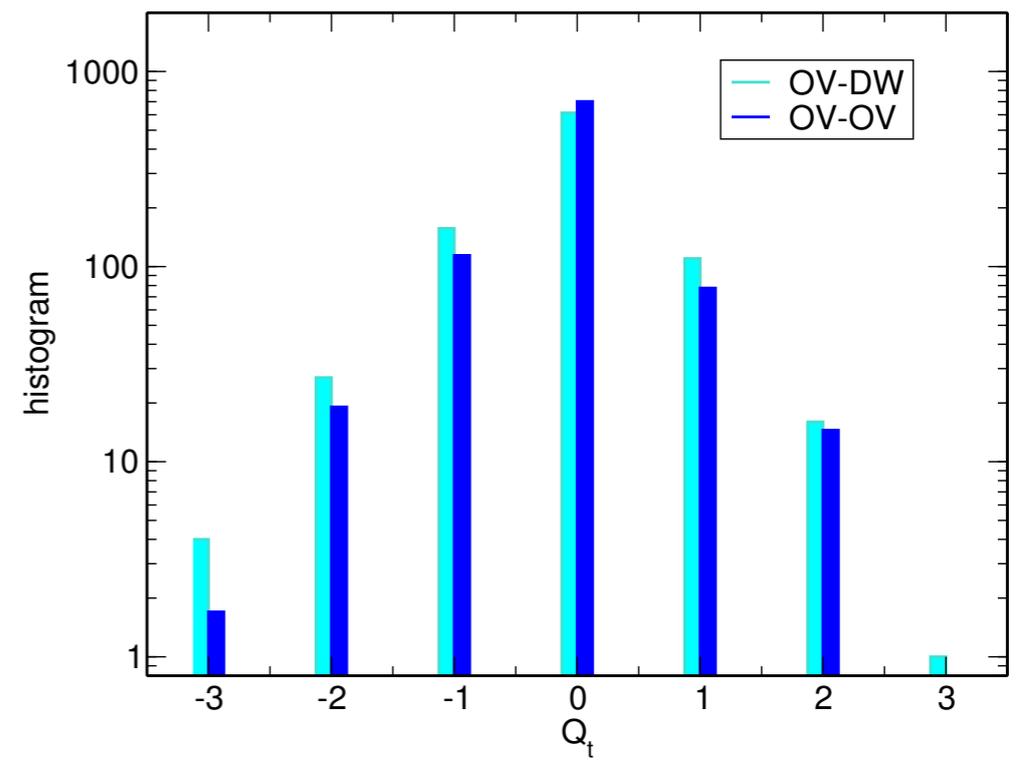
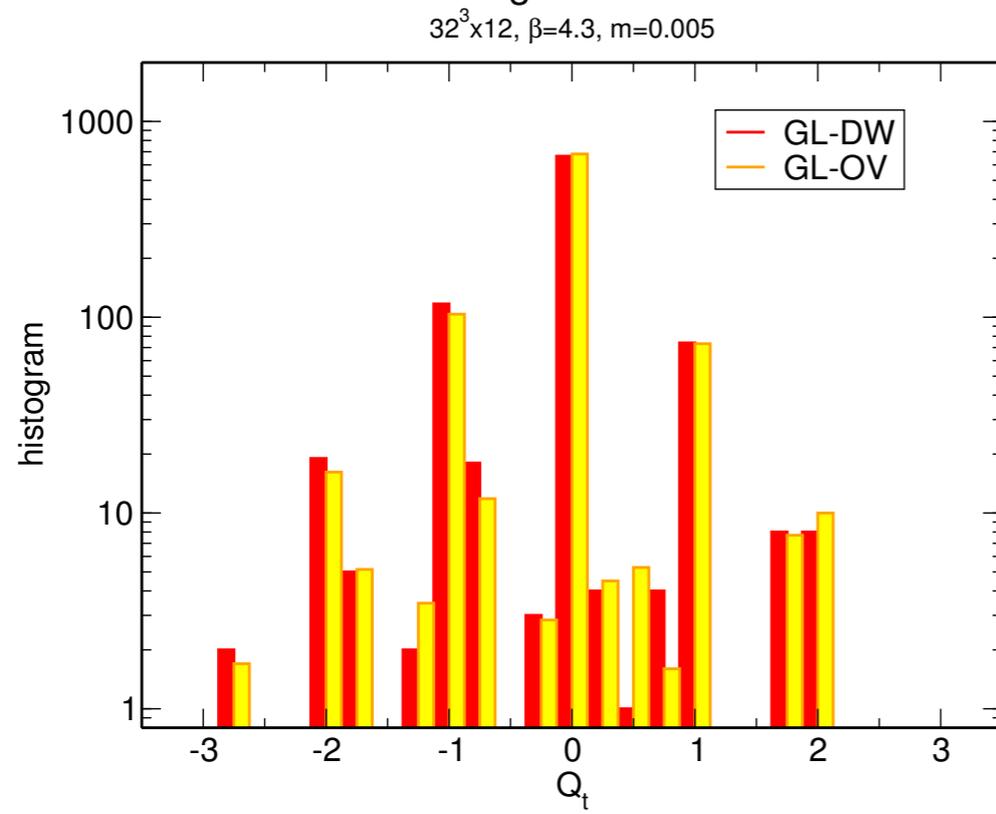
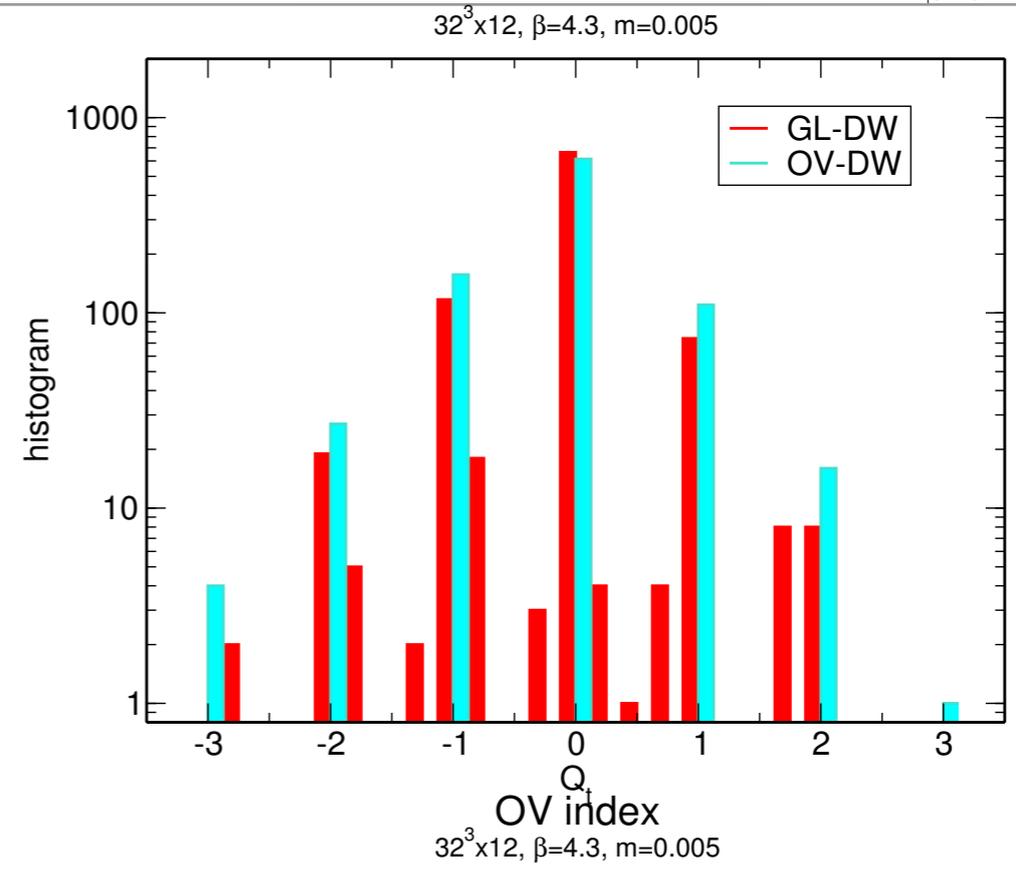
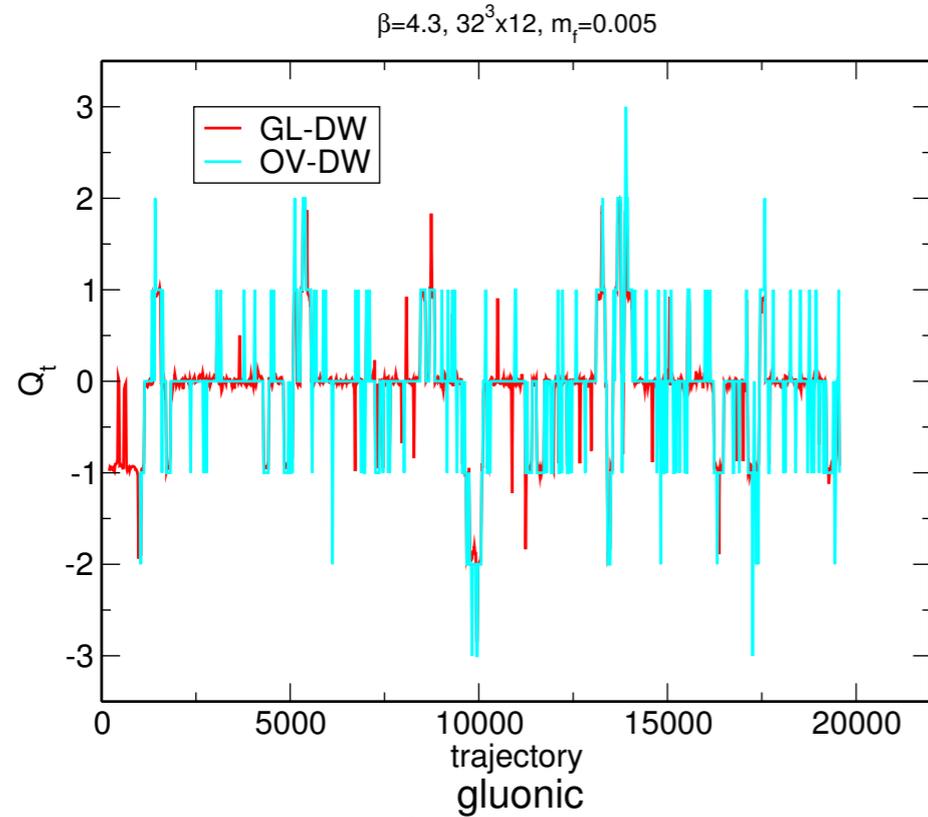
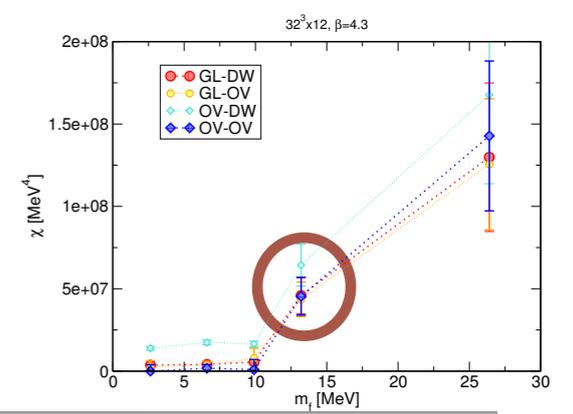
- reweighting: before / after and above 2 meas. yield 4  $\chi_t$  values
- current focus:  $1/a = 2.6$  GeV \*\*\* **PRELIMINARY** \*\*\*

$\chi_t(m_f)$  for  $N_f=2$   $T=220$  MeV

GL-DW	gluonic charge on DW
GL-OV	gluonic charge on OV
OV-DW	OV index on DW ensemble
OV-OV	OV index on OV ensemble



# m=0.005 history and histogram

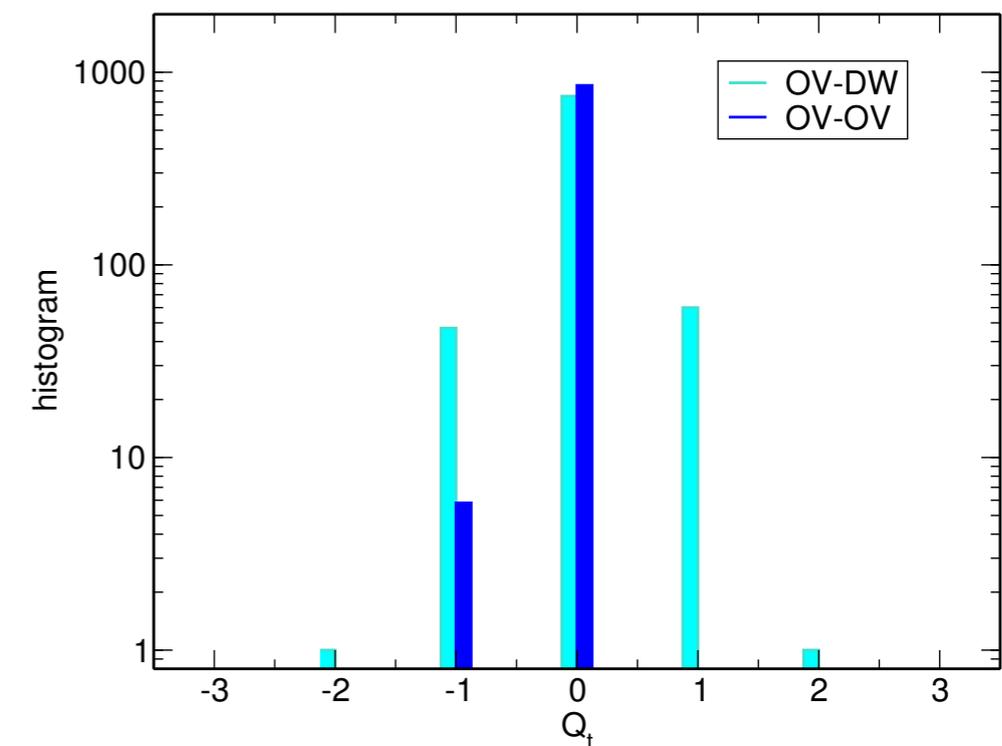
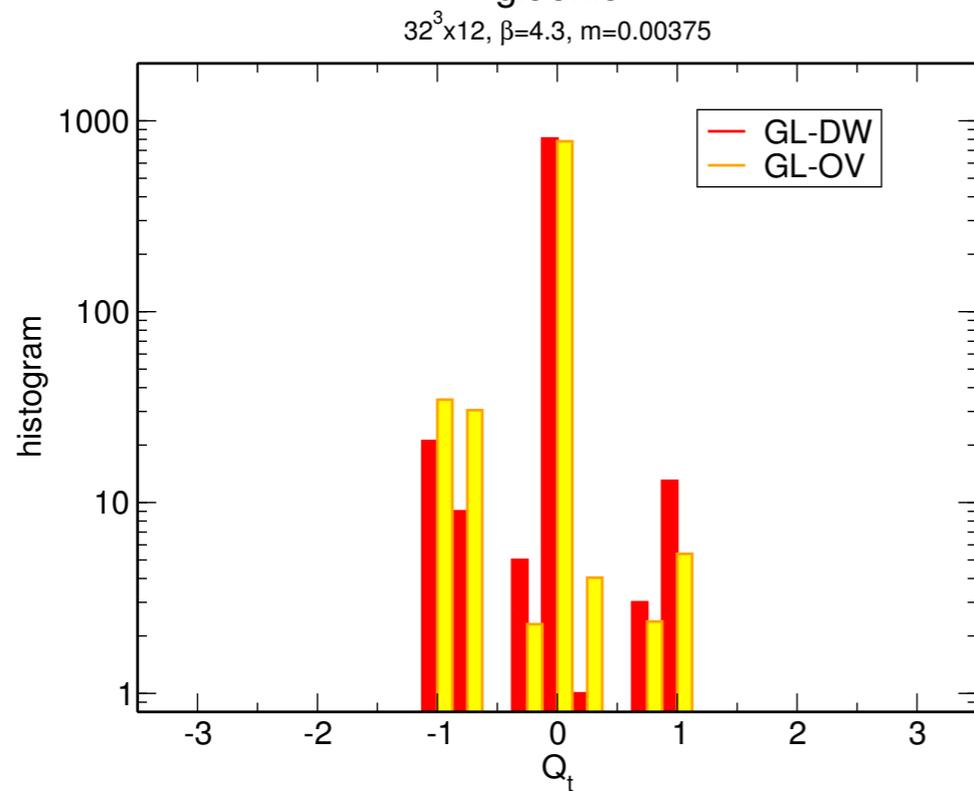
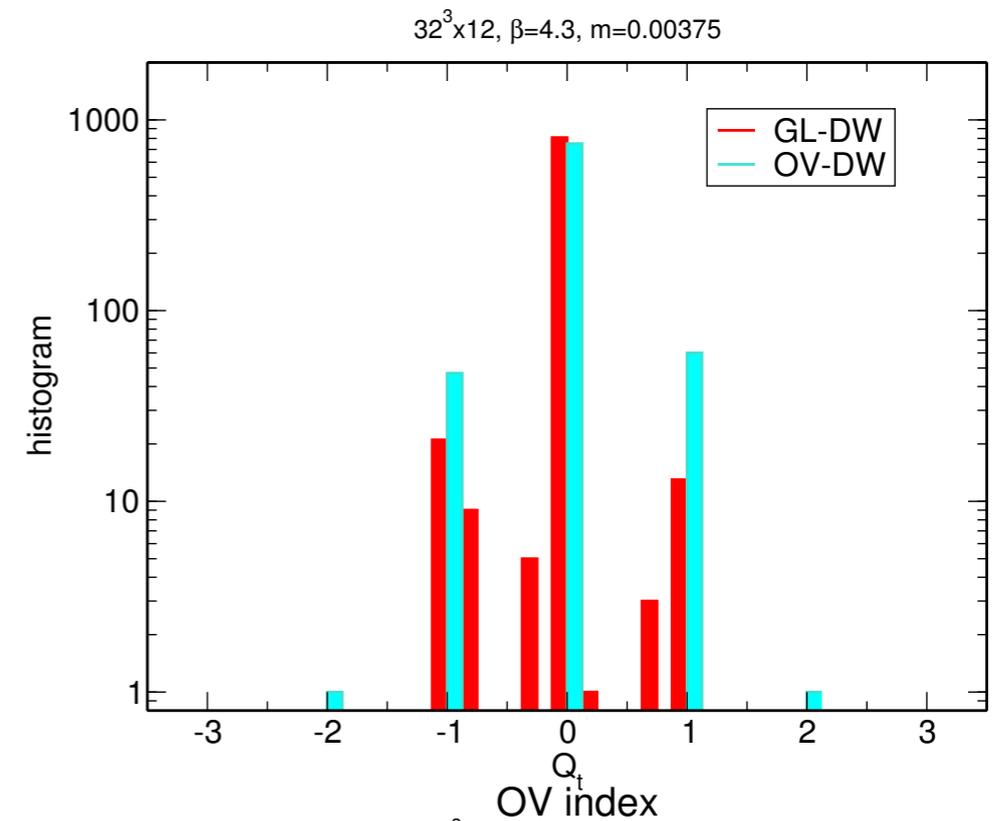
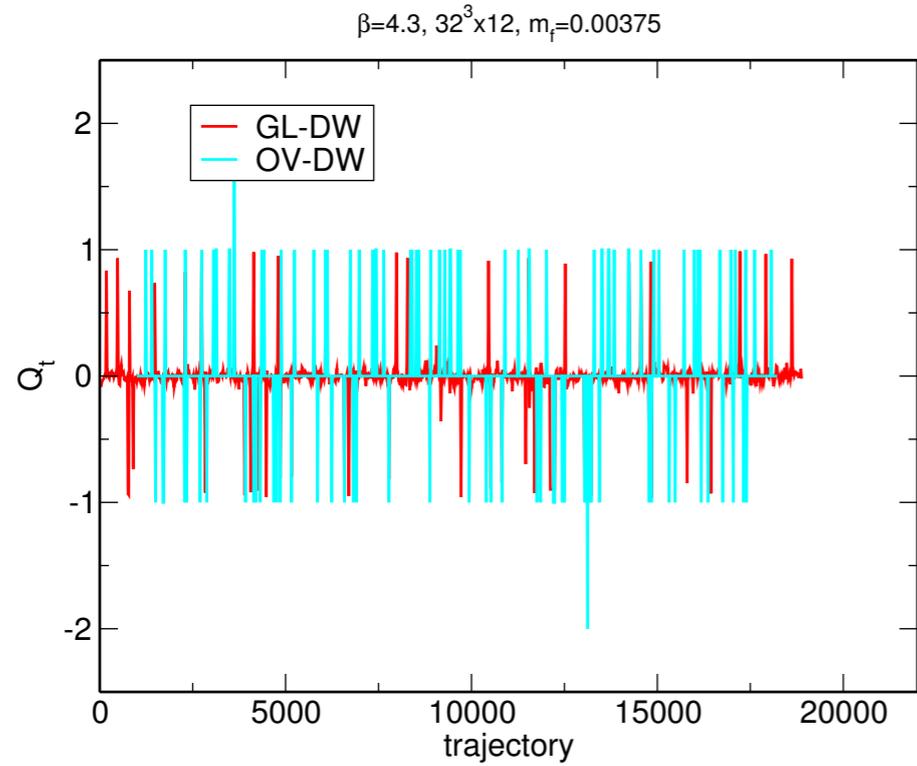
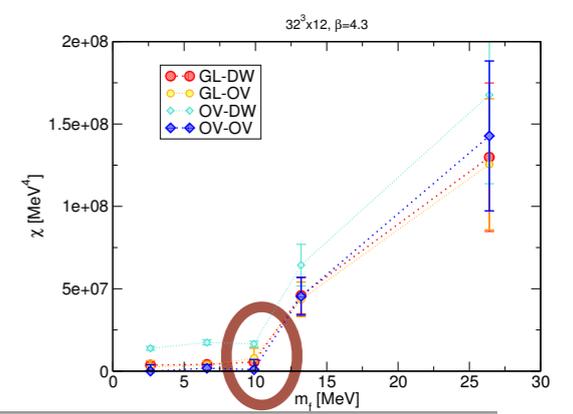


# Out of 4 definitions of susceptibility

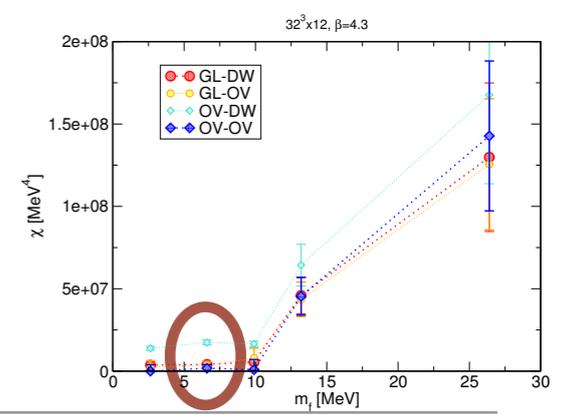
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- **OV** index on **DWF** has pathology: overly enhanced artificial zero modes
  - should be excluded from the discussion
  - reweighting to **OV** ensemble suppresses the weight of fake zero modes
- with **GLuonic Q**: before / after reweighting does not change much
  - Wilson Flow: tends to erase artificial modes
  - OV reweighting: artificial zero modes have suppressed reweighting factor
  - shall take “before” reweighting (**GL-DW**), which is statistically better
- Good ones: **GL-DW**, **OV-OV**
- Discarded: **GL-OV**, **OV-DW**

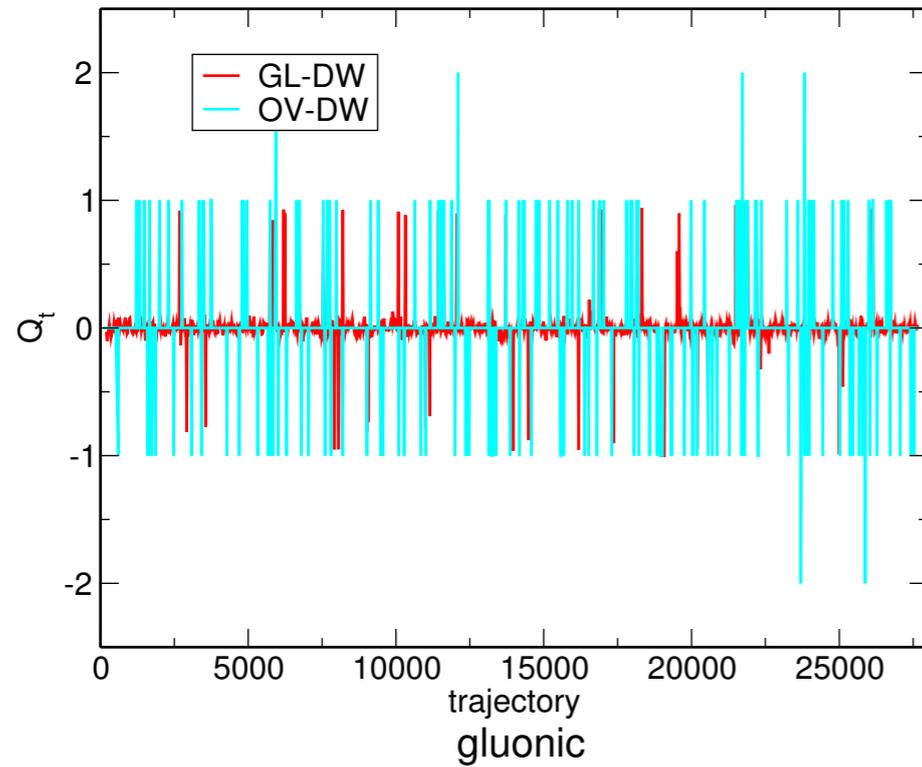
# $m=0.00375$ history and histogram



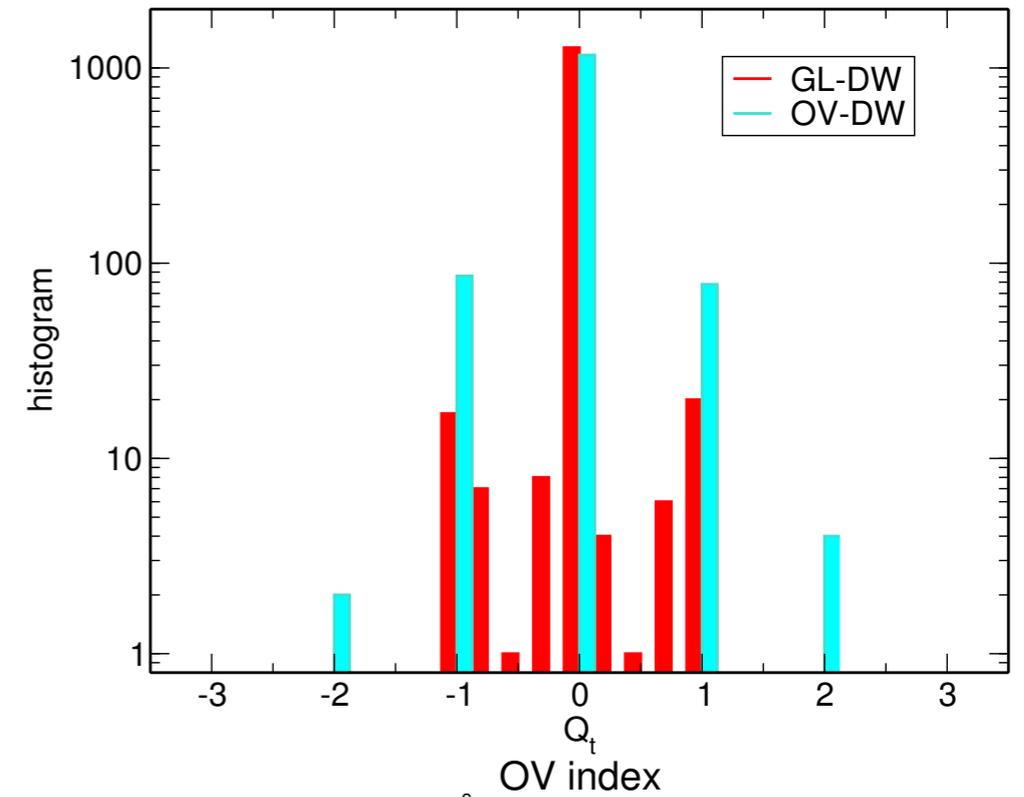
# m=0.0025 history and histogram



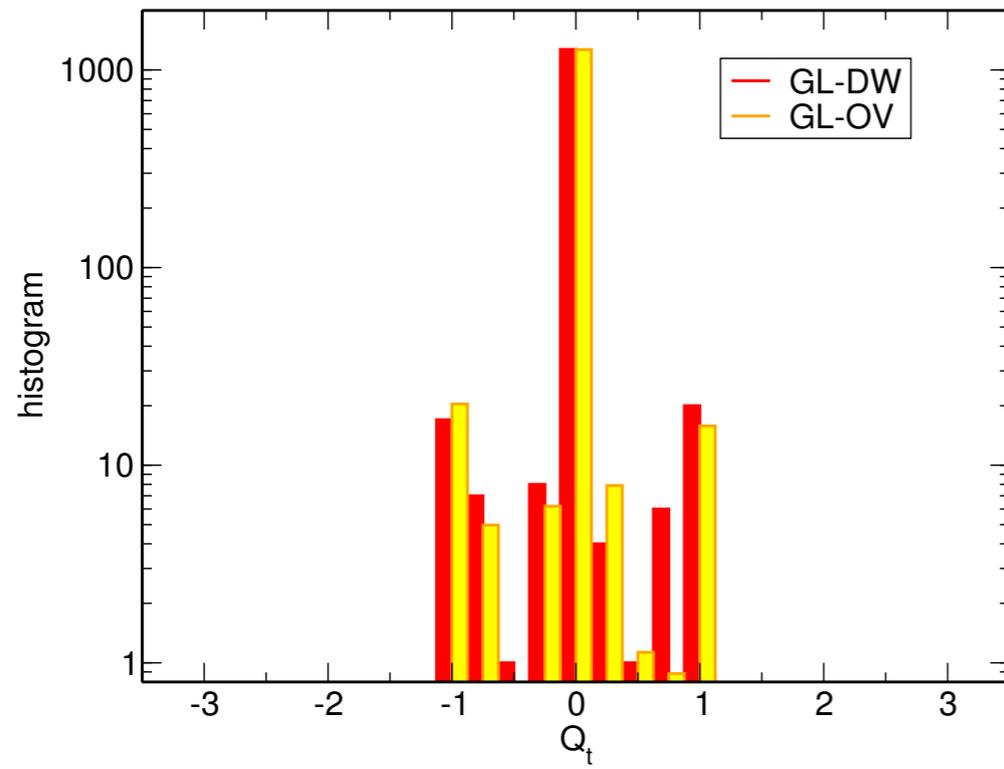
$\beta=4.3, 32^3 \times 12, m_t=0.0025$



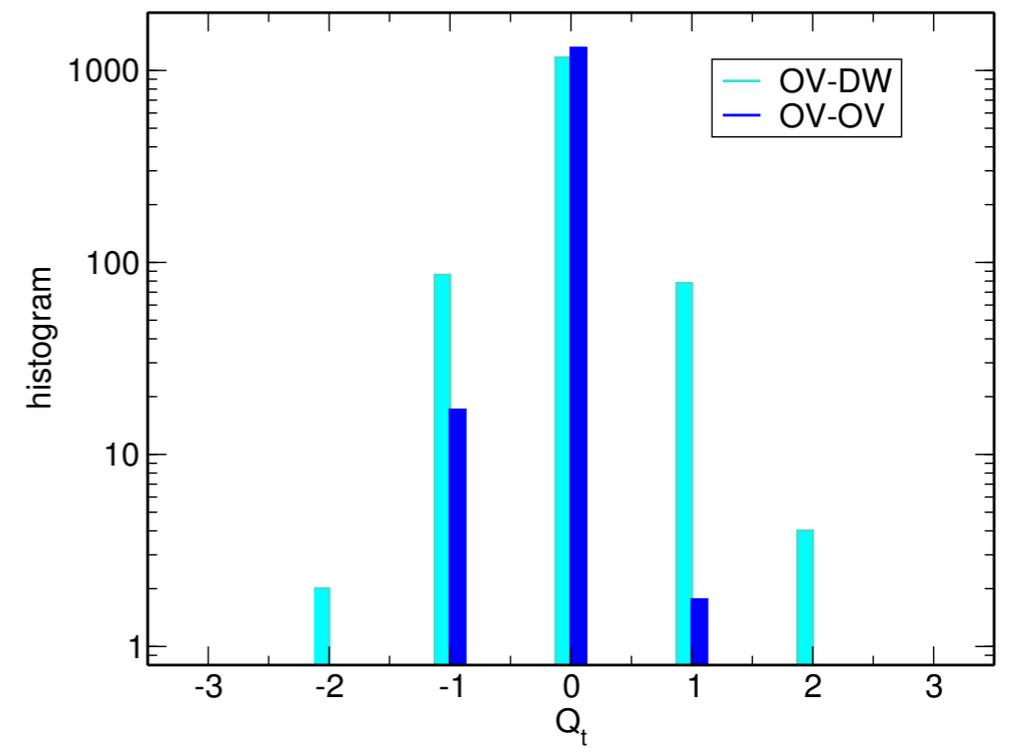
32<sup>3</sup>x12,  $\beta=4.3, m=0.0025$



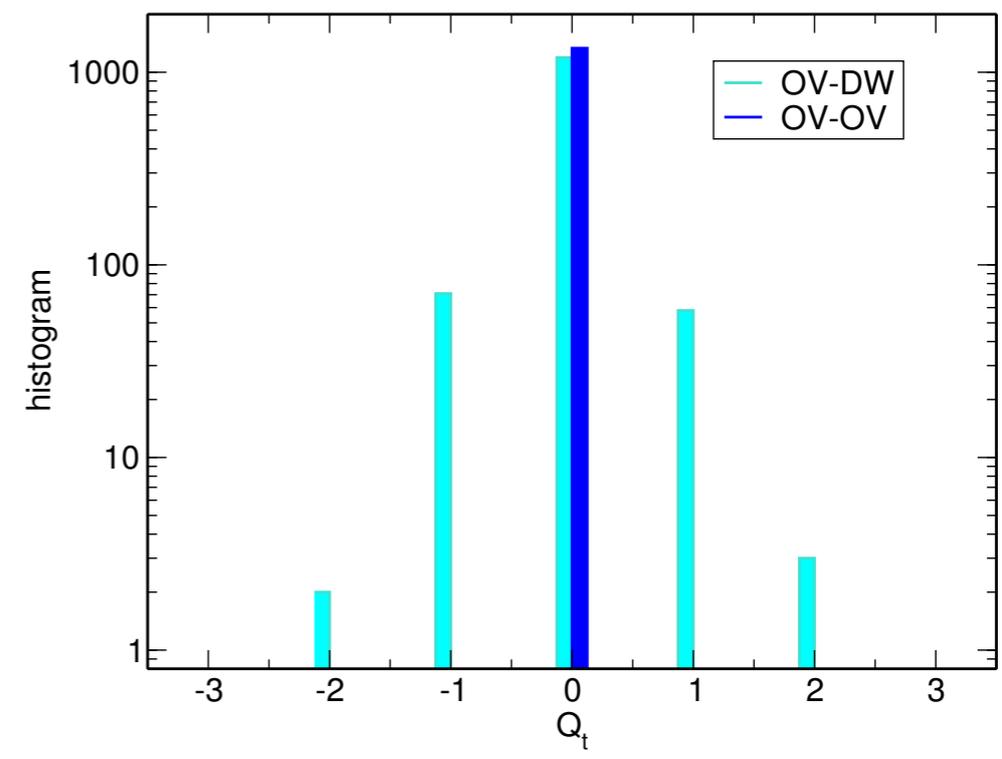
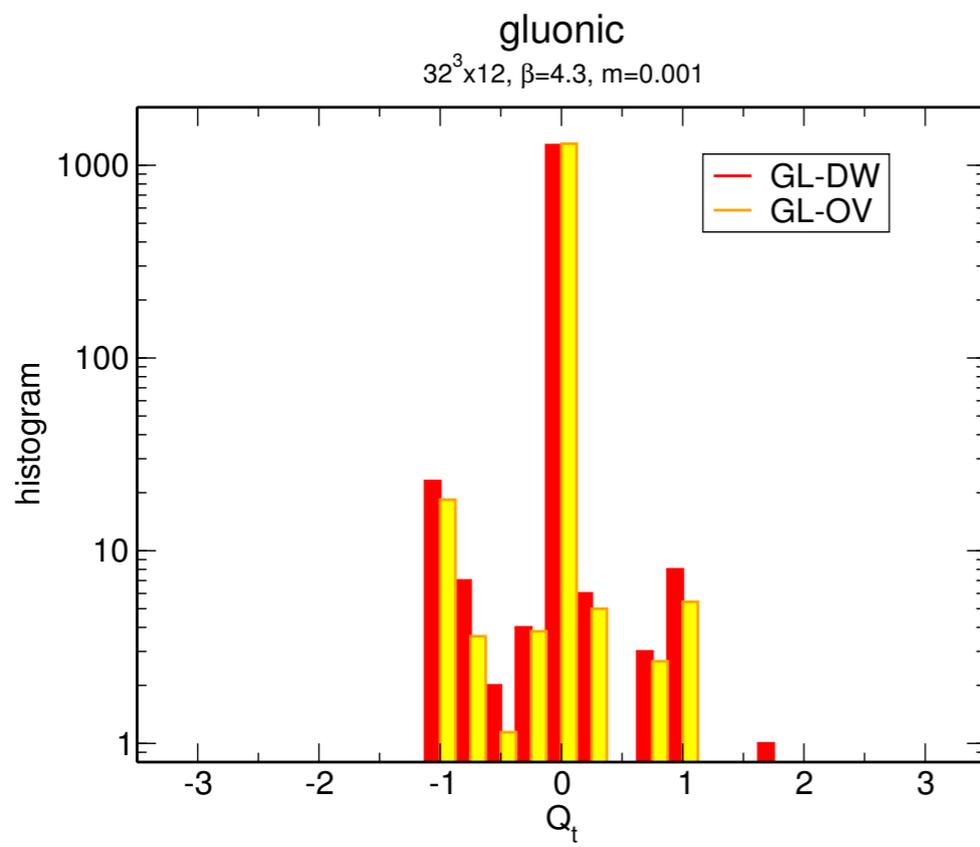
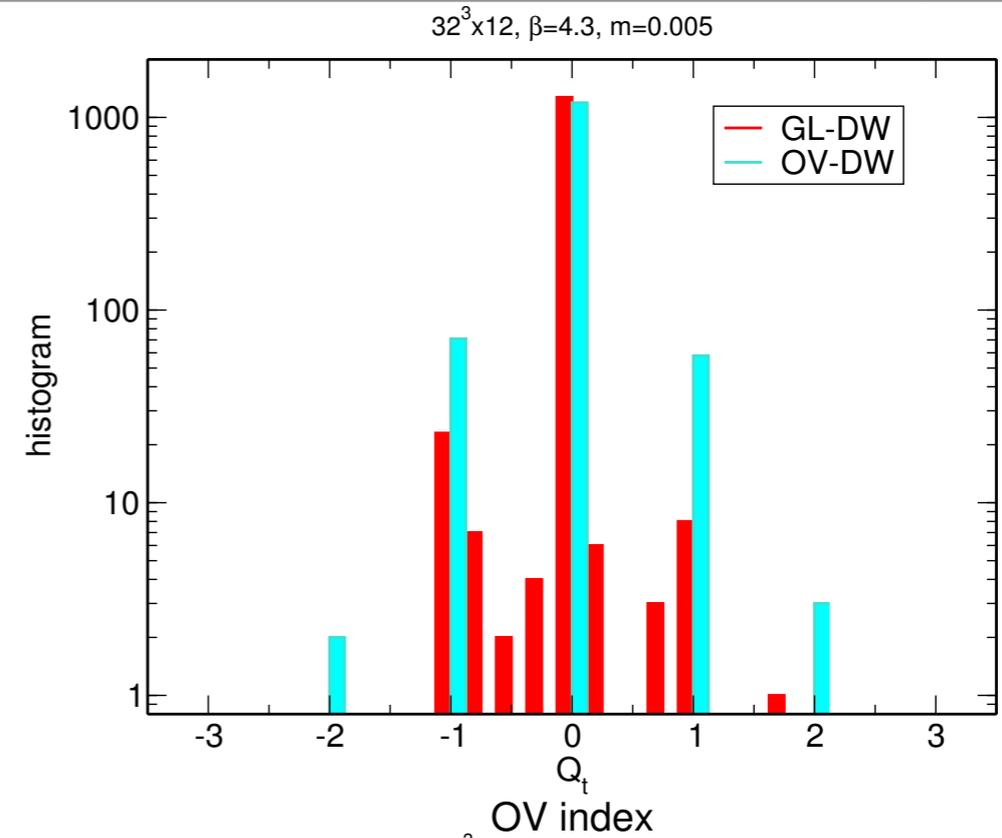
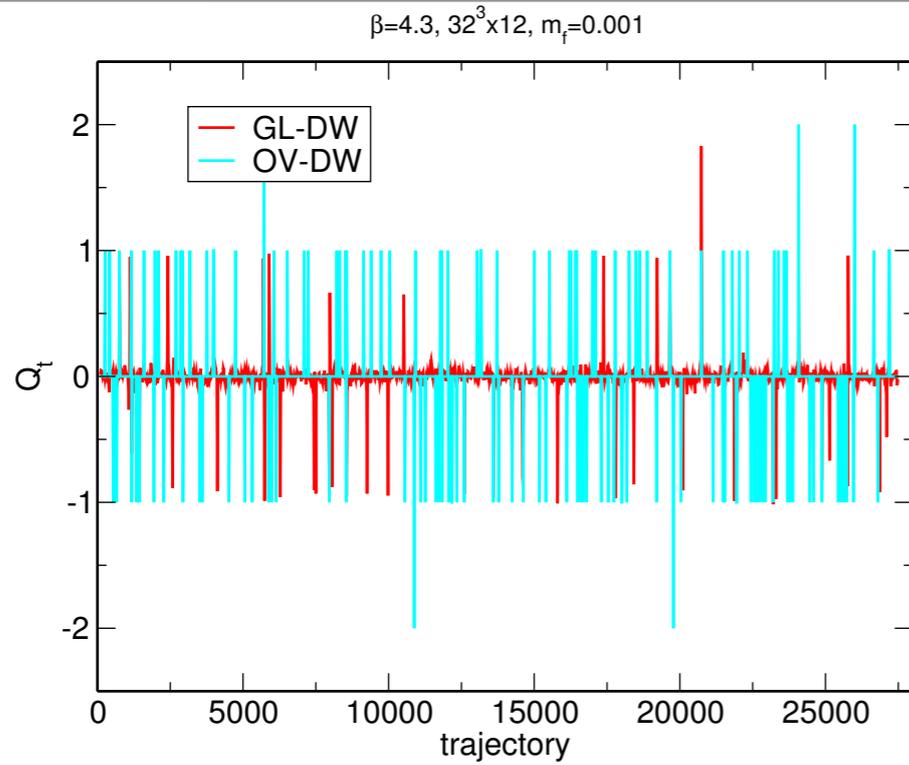
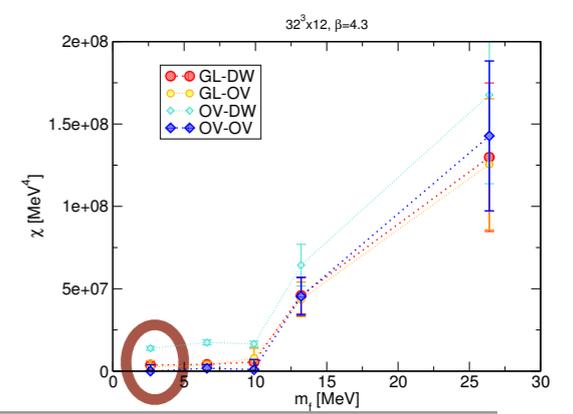
32<sup>3</sup>x12,  $\beta=4.3, m=0.0025$



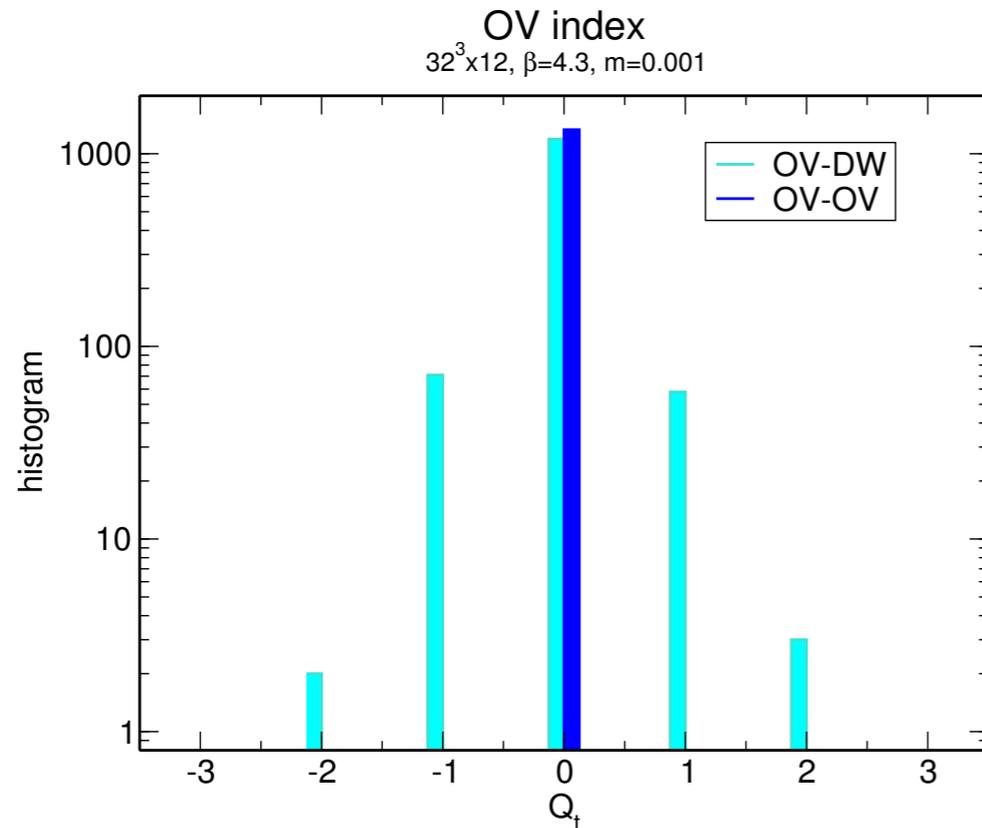
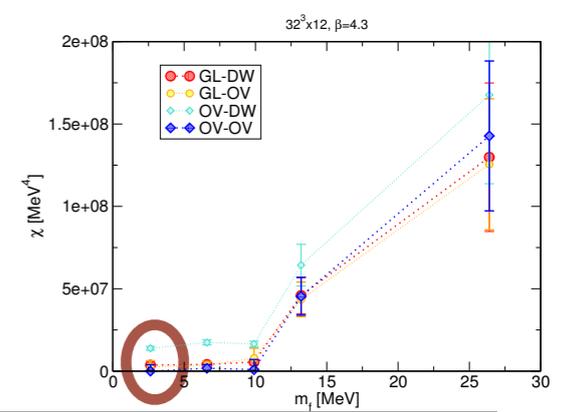
32<sup>3</sup>x12,  $\beta=4.3, m=0.0025$



# m=0.001 history and histogram



# resolution of susceptibility (ex: $m=0.001$ )



## Effective number of statistics

- decreases with reweighting
- $N_{\text{eff}} = N_{\text{conf}} \langle R \rangle / R_{\text{max}}$
- $N_{\text{conf}} = 1326 \rightarrow N_{\text{eff}} = 32$

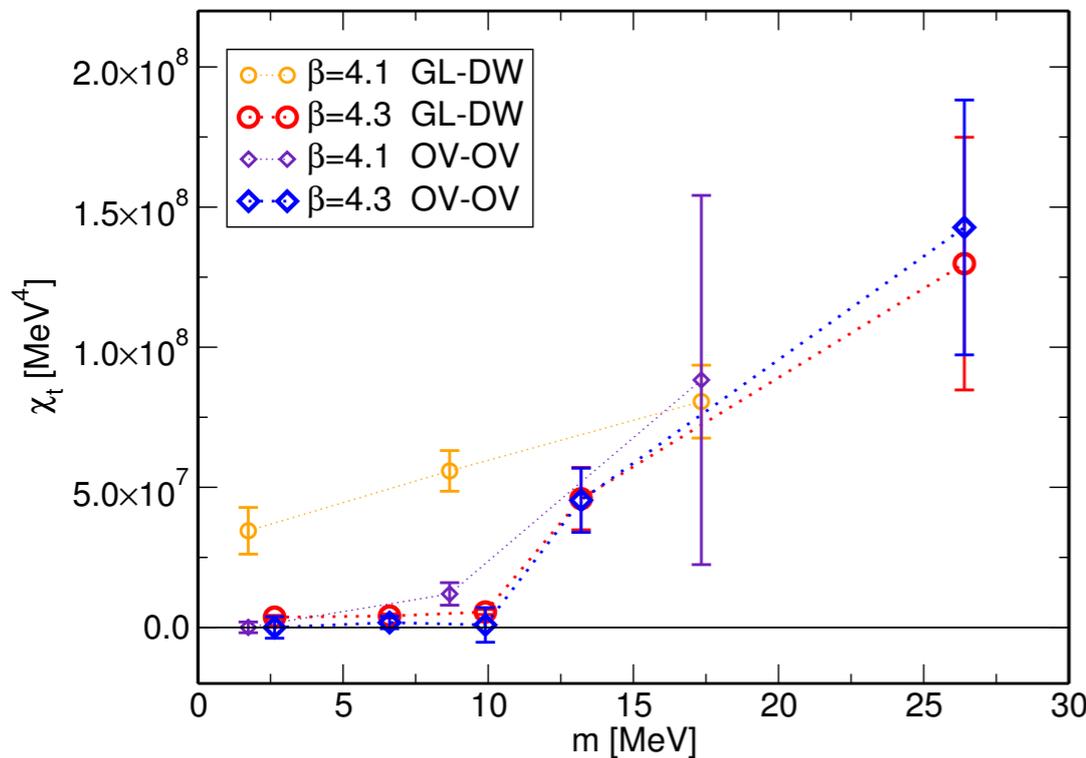
null measurement of topological excitation after reweighting

- does not readily mean  $\chi_t = 0$ : (this case  $\langle Q^2 \rangle = 4(4) \times 10^{-6}$ )
- there must be a resolution of  $\chi_t$  under given statistics
  - [resolution of  $\langle Q^2 \rangle$ ] =  $1/N_{\text{eff}}$
- shall take the “statistical error” of  $\langle Q^2 \rangle = \max(\Delta \langle Q^2 \rangle, 1/N_{\text{eff}})$

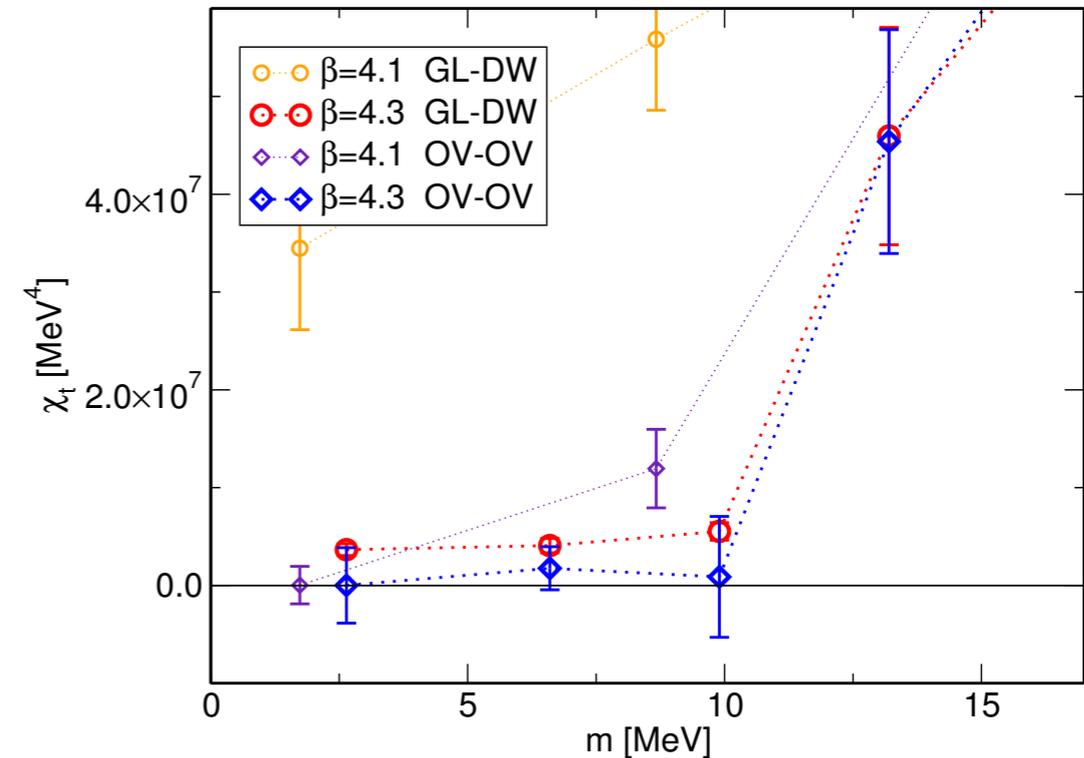
# $\chi_t(m)$ $T \sim 220$ MeV discretization effect

comparing  $1/a=1.7$  GeV and  $1/a=2.6$  GeV (  $(3.6\text{fm})^3$  and  $(2.4\text{fm})^3$  )

compare  $N_t=8(\beta=4.1)$  and  $12(4.3)$  at similar temperature (217 and 220 MeV)



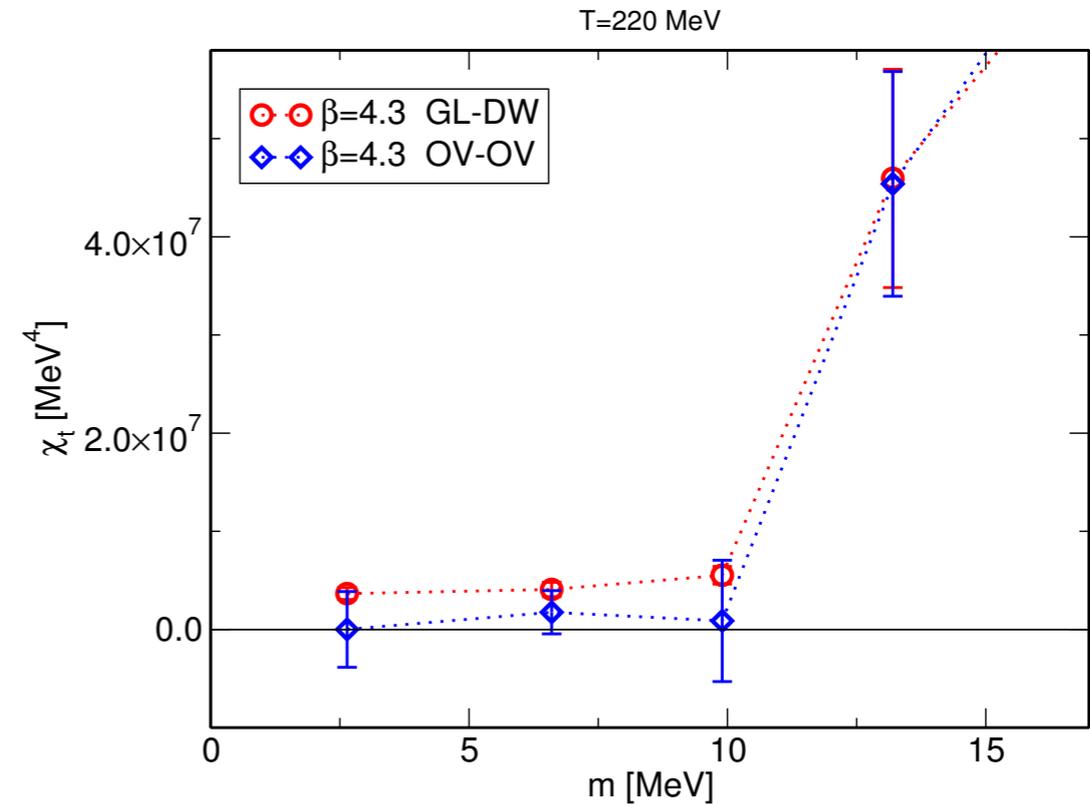
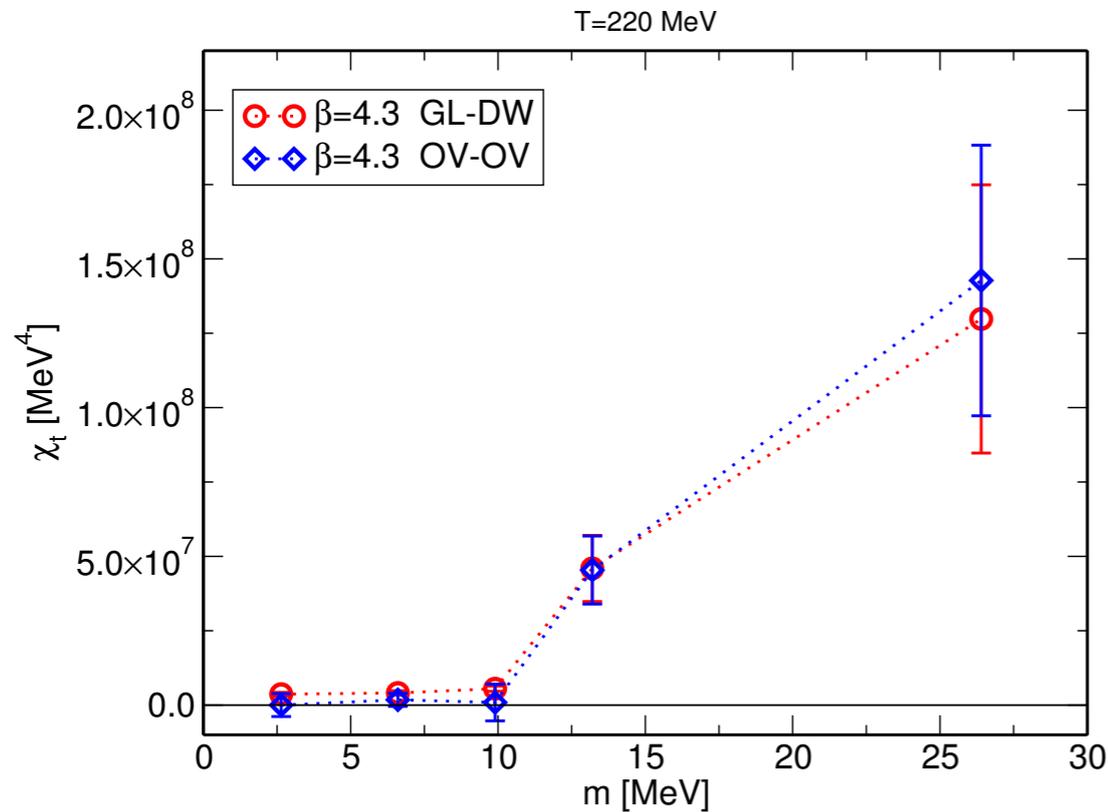
compare  $N_t=8(\beta=4.1)$  and  $12(4.3)$  at similar temperature (217 and 220 MeV)



- **OV-OV**: better scaling
- **GL-DW**: large scaling violation for smaller  $m$
- **OV-OV**:  $\chi_t = 0$  (within error) for  $0 \leq m \lesssim 10$  MeV
- **GL-DW**:  $\chi_t > 0$ , but, may well decrease as  $a$   
 ➔ (consistent with **OV-OV** with large error of **OV-OV**)

$\chi_t(m)$   $T \sim 220$  MeV finer lattice

$1/a = 2.6$  GeV



suggesting 2 regions

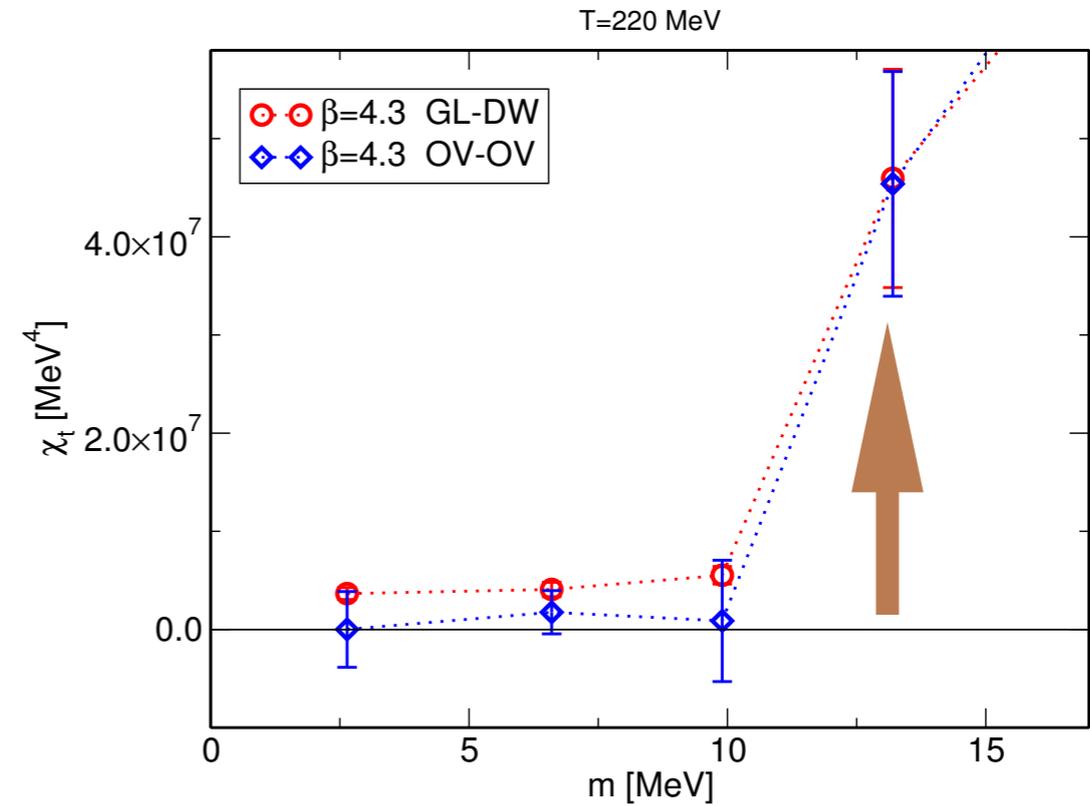
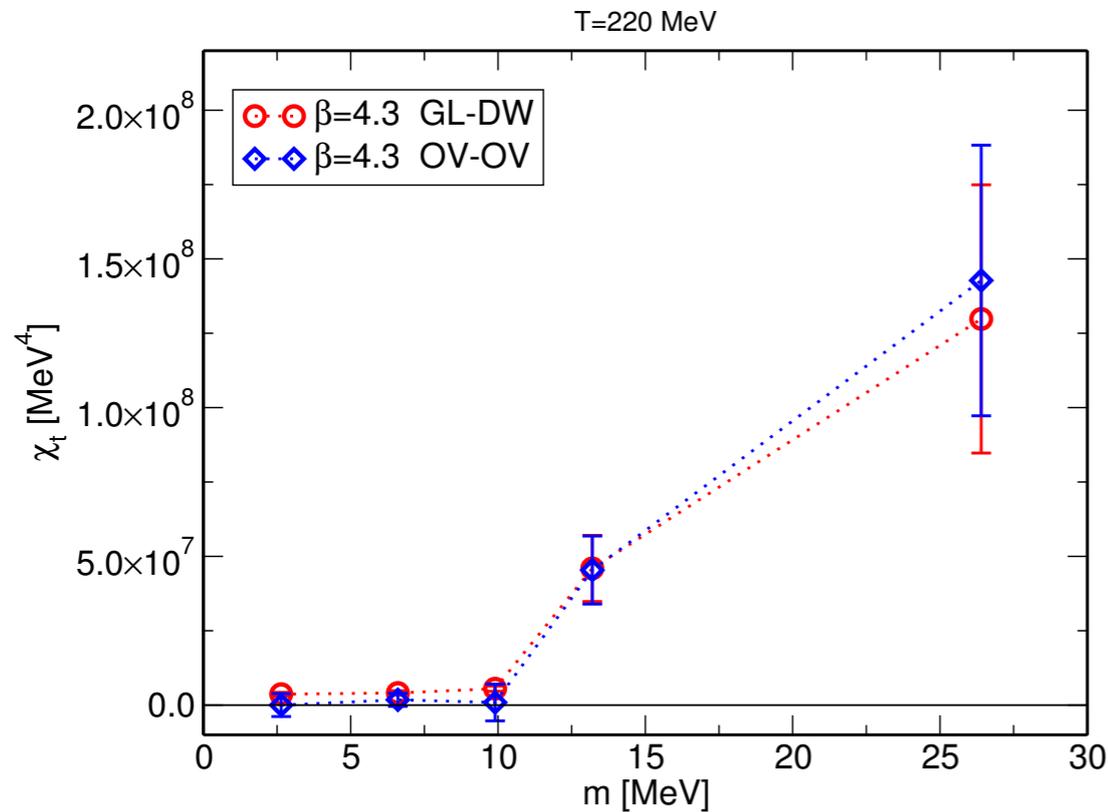
1:  $\chi_t$  is very small (may vanish in  $a \rightarrow 0$ ):  $0 \leq m \lesssim 10$  MeV

2: sudden **growth** of  $\chi_t$  :  $10$  MeV  $\lesssim m$

- physical ud mass point:  $m \approx 4$  MeV

$\chi_t(m)$   $T \sim 220$  MeV finer lattice

$1/a = 2.6$  GeV



suggesting 2 regions

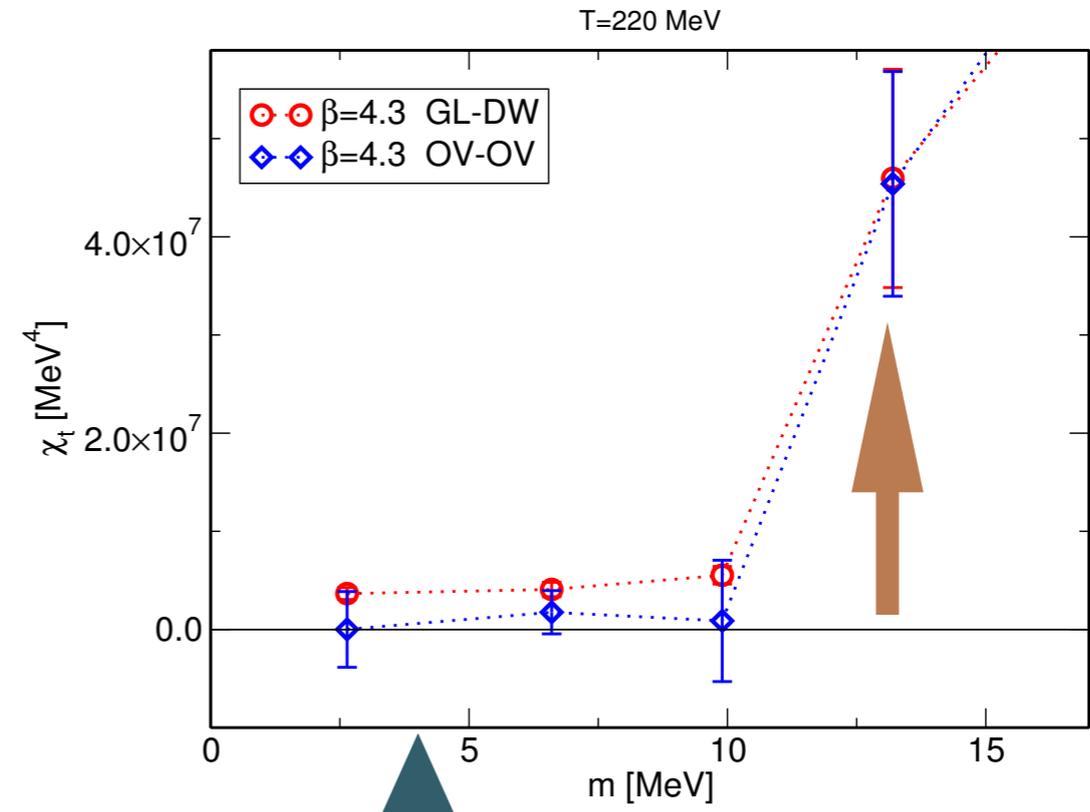
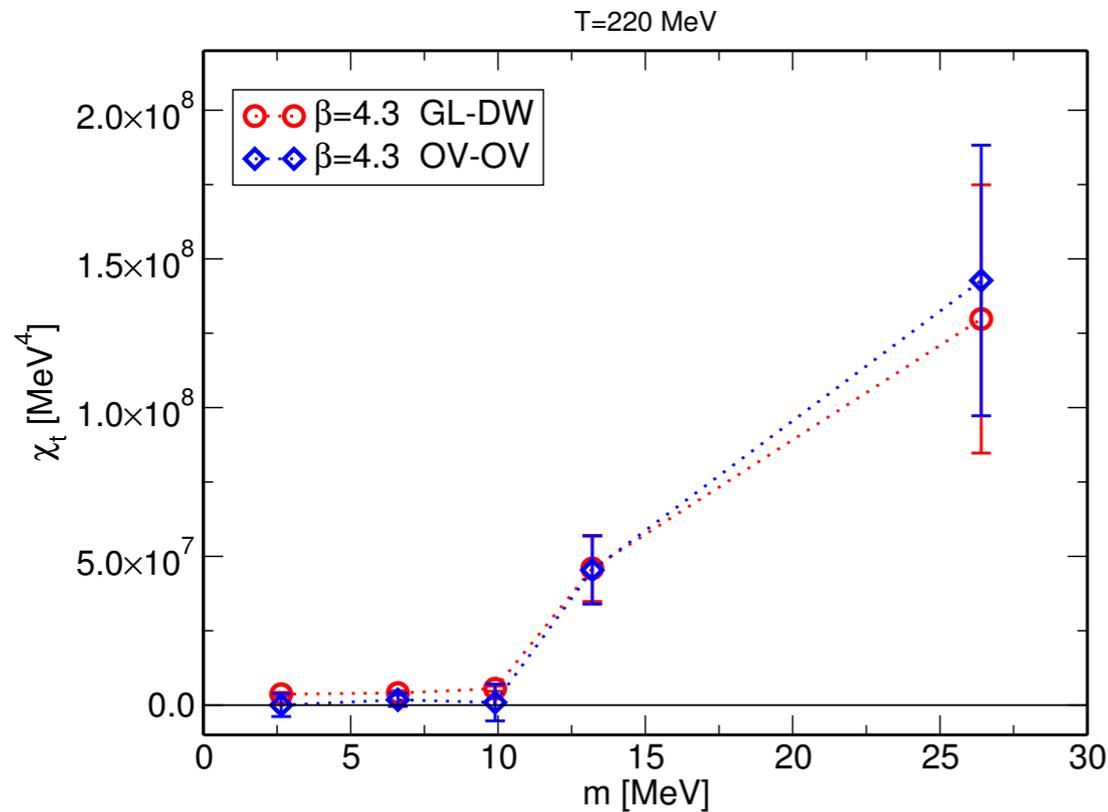
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$\chi_t(m)$   $T \sim 220$  MeV finer lattice

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suggesting 2 regions

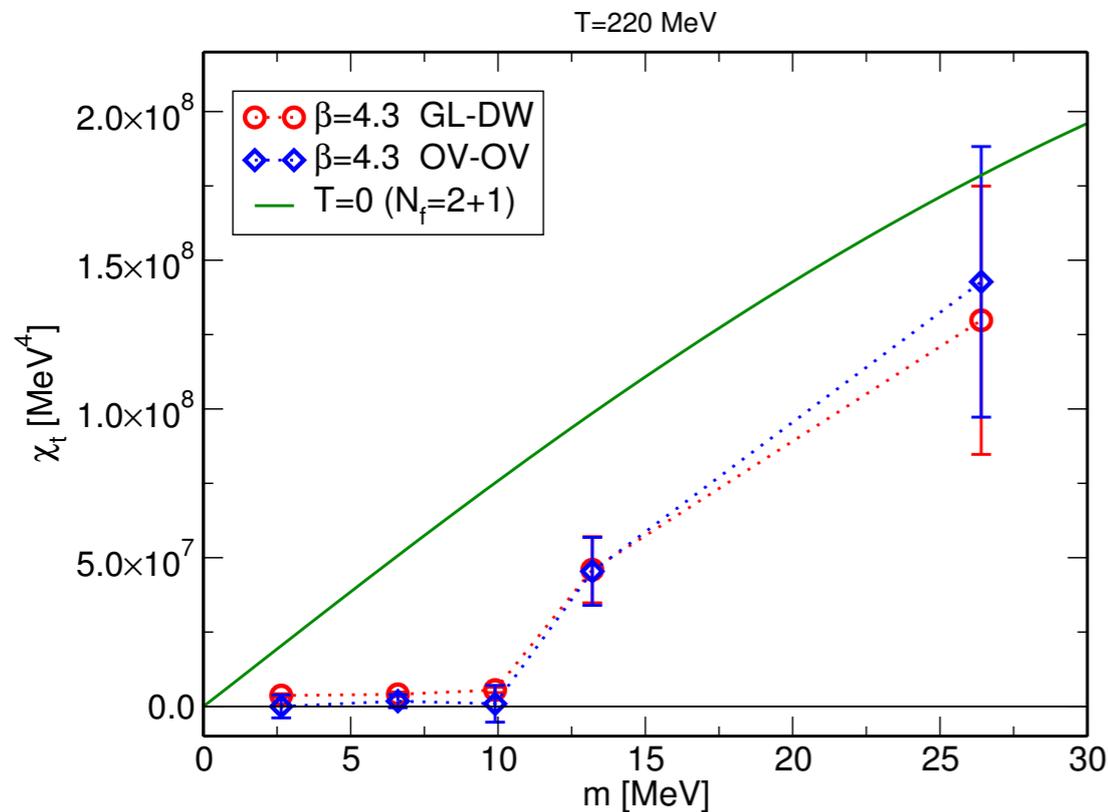
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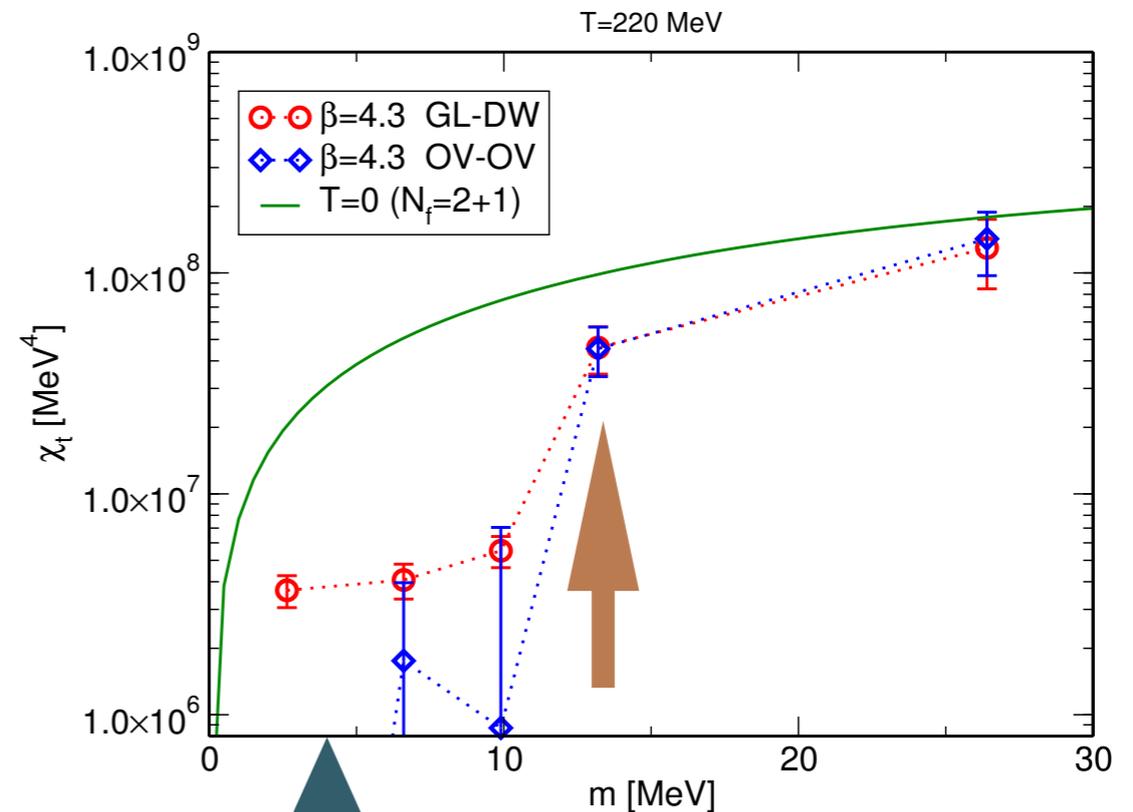
- physical ud mass point:  $m \approx 4$  MeV

# $\chi_t(m)$ $T \sim 220$ MeV finer lattice against $T=0$

$1/a = 2.6$  GeV



[ $T=0$ : JLQCD: S.Aoki et al 2017,  $N_f=2+1$ ]



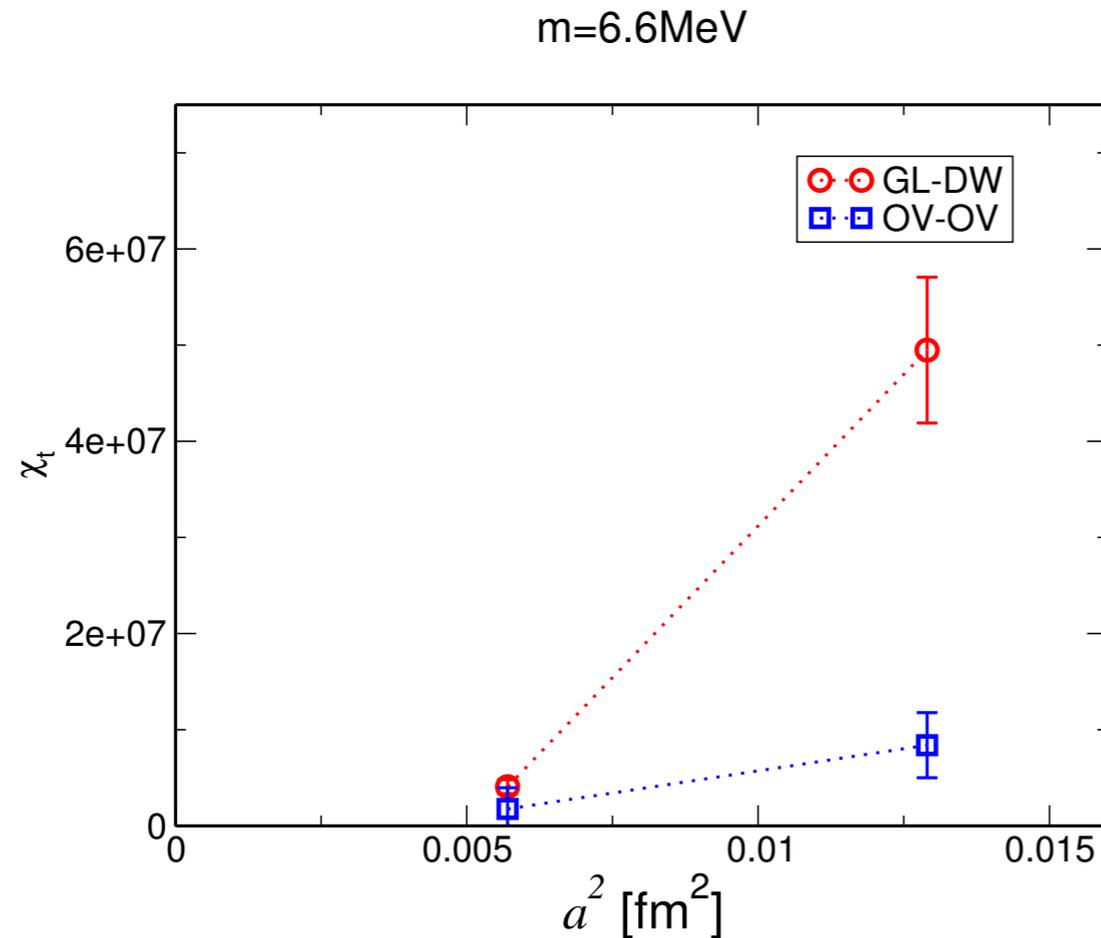
suggesting 2 regions

1:  $\chi_t$  is very small (may vanish in  $a \rightarrow 0$ ):  $0 \leq m \lesssim 10$  MeV

2: sudden **growth** of  $\chi_t$  :  $10$  MeV  $\lesssim m$

→ surprisingly close to  $T=0$  value

$\chi_t(m)$   $T=220$  MeV  $a^2$  scaling:  $m=6.6$  MeV

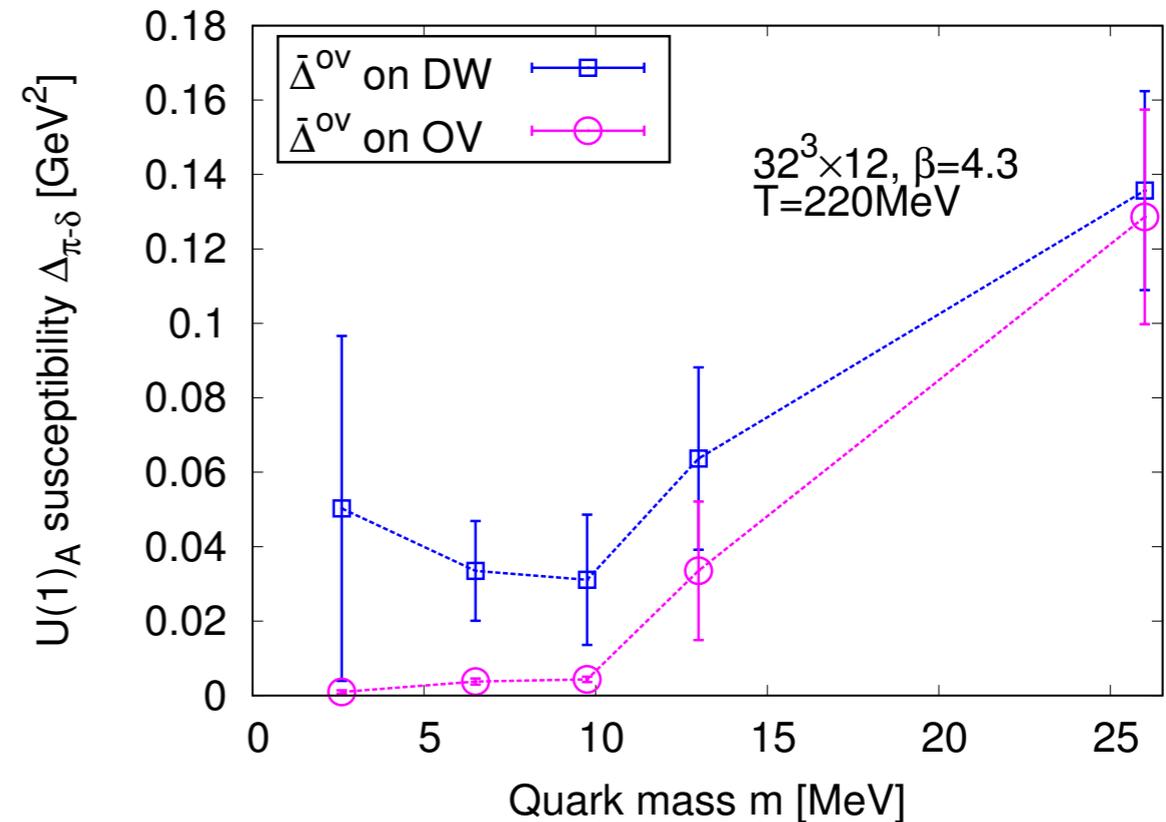
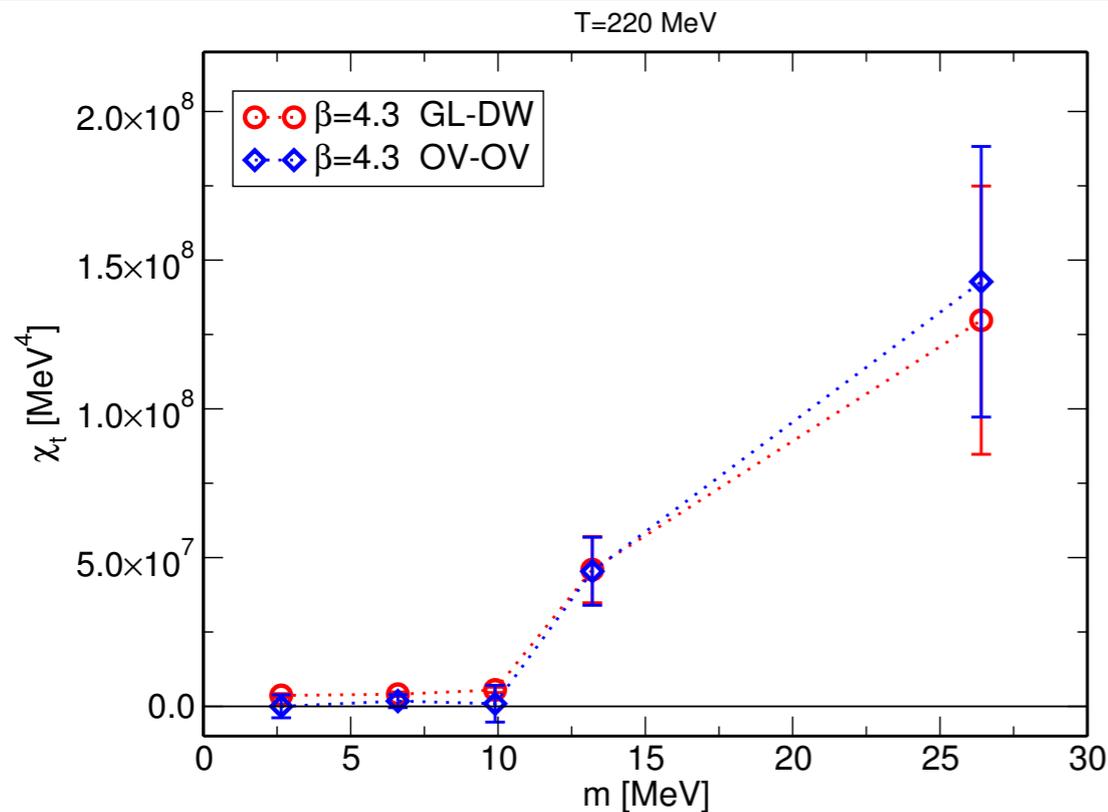


(  $V=(3.6\text{fm})^3$  and  $(2.4\text{fm})^3$  )

continuum scaling in 1st region

- $m=6.6$  MeV
- vanishing towards continuum limit
- caveat: physical volume is different  $\rightarrow$  needs further invest.

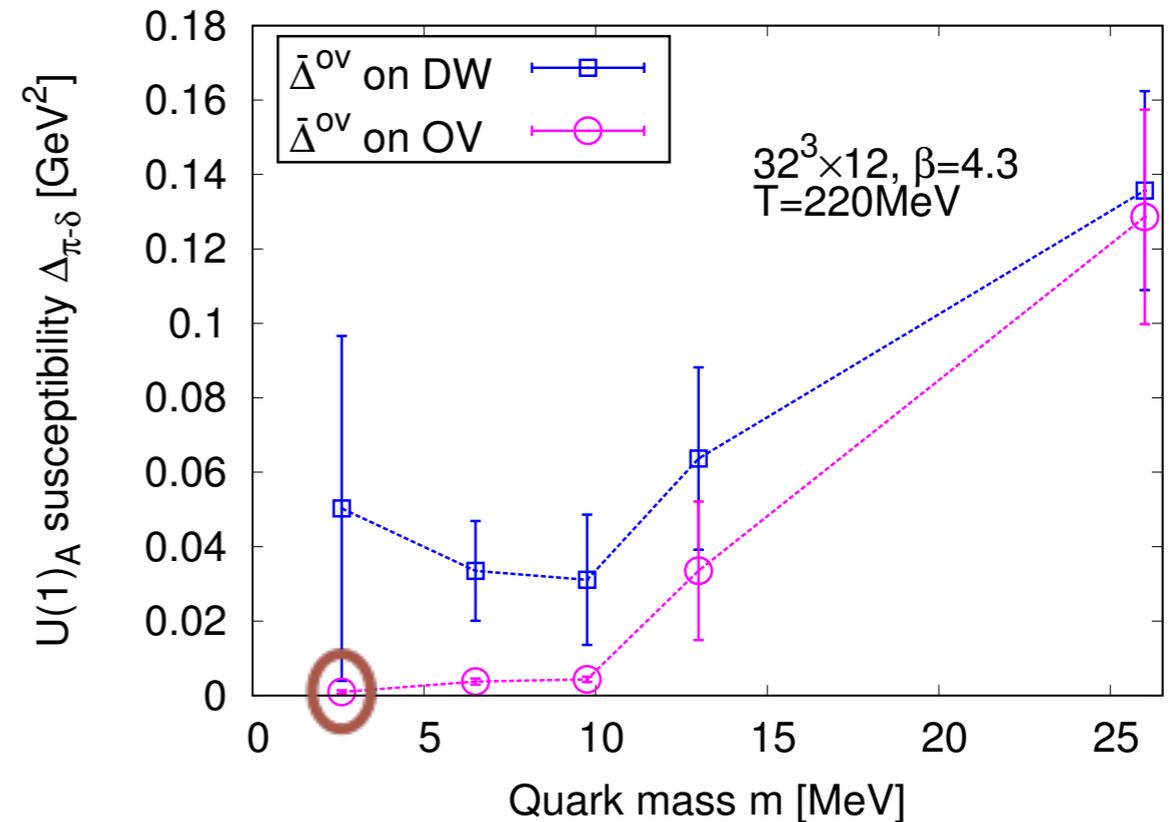
# $\chi_t$ and $\Delta_{\pi-\delta}$ ( $U_A(1)$ order parameter) @ $T \sim 220$ MeV



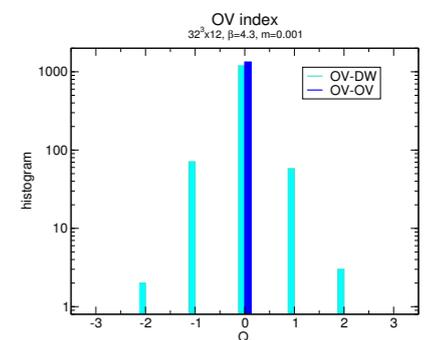
- $\Delta_{\pi-\delta}$  has a similar trend as function of  $m$
- AFT scenario: [Aoki, Fukaya, Taniguchi]
  - $\chi = 0$  for  $0 < m < m_c \rightarrow$  consistent with left figure
  - $\Delta_{\pi-\delta} = 0$  for  $m=0$ , but does not say anything for  $m>0$ : OK
- KY scenario...

competing scenarios for

$\chi_t$  and  $\Delta_{\pi-\delta}$  ( $U_A(1)$  order parameter) @  $T \sim 220$  MeV

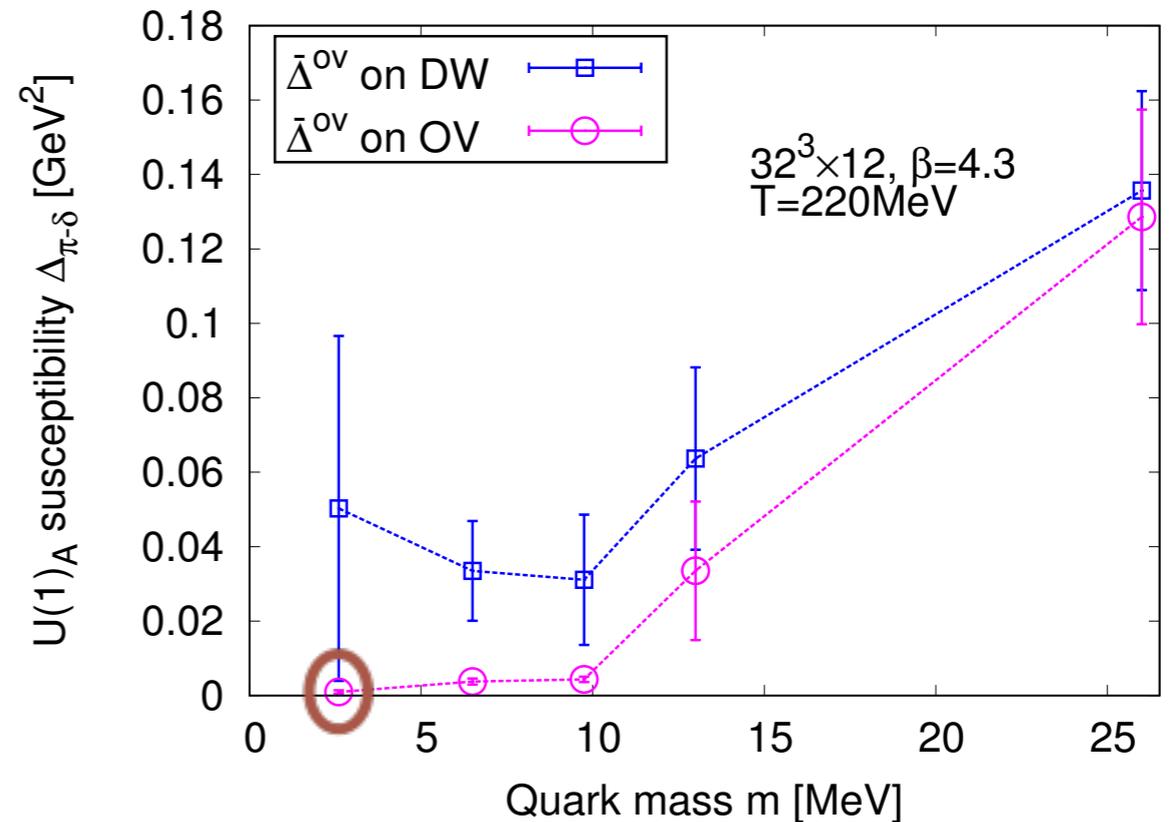
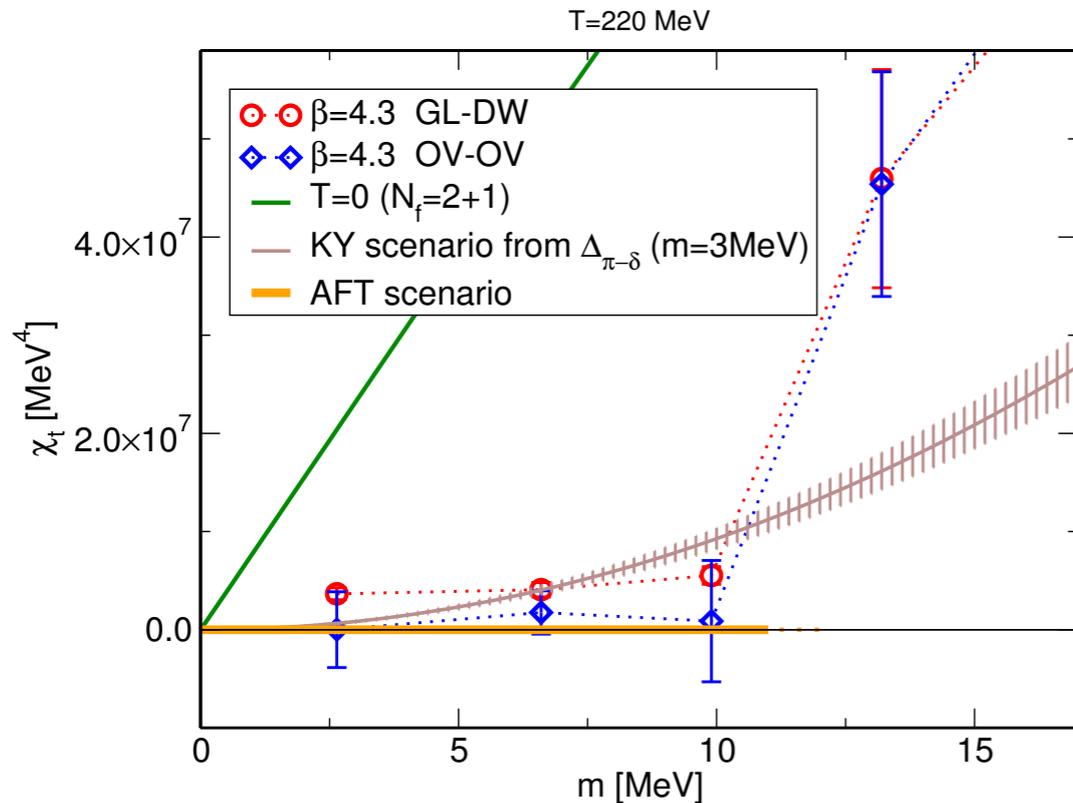


- KY scenario [Kanazawa, Yamamoto 2016]
  - $\Delta_{\pi-\delta}$ : including zero mode cont. is proper
  - $\Delta_{\pi-\delta} = \text{const} > 0$
  - $\Delta_{\pi-\delta} \approx 8 V f_A^2 m^2$  for **Q=0 sector** (for  $2V f_A m^2 < 1$ )
- $\Delta_{\pi-\delta}$  @ lightest point only from Q=0
- $\chi_t = 2 f_A m^2$
- inconsistent with  $m > 10$  MeV  $\chi_t$  growth

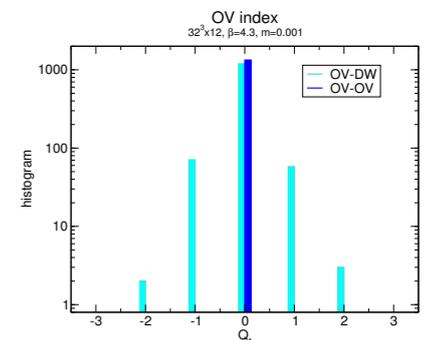


# competing scenarios for

$\chi_t$  and  $\Delta_{\pi-\delta}$  ( $U_A(1)$  order parameter) @  $T \sim 220$  MeV

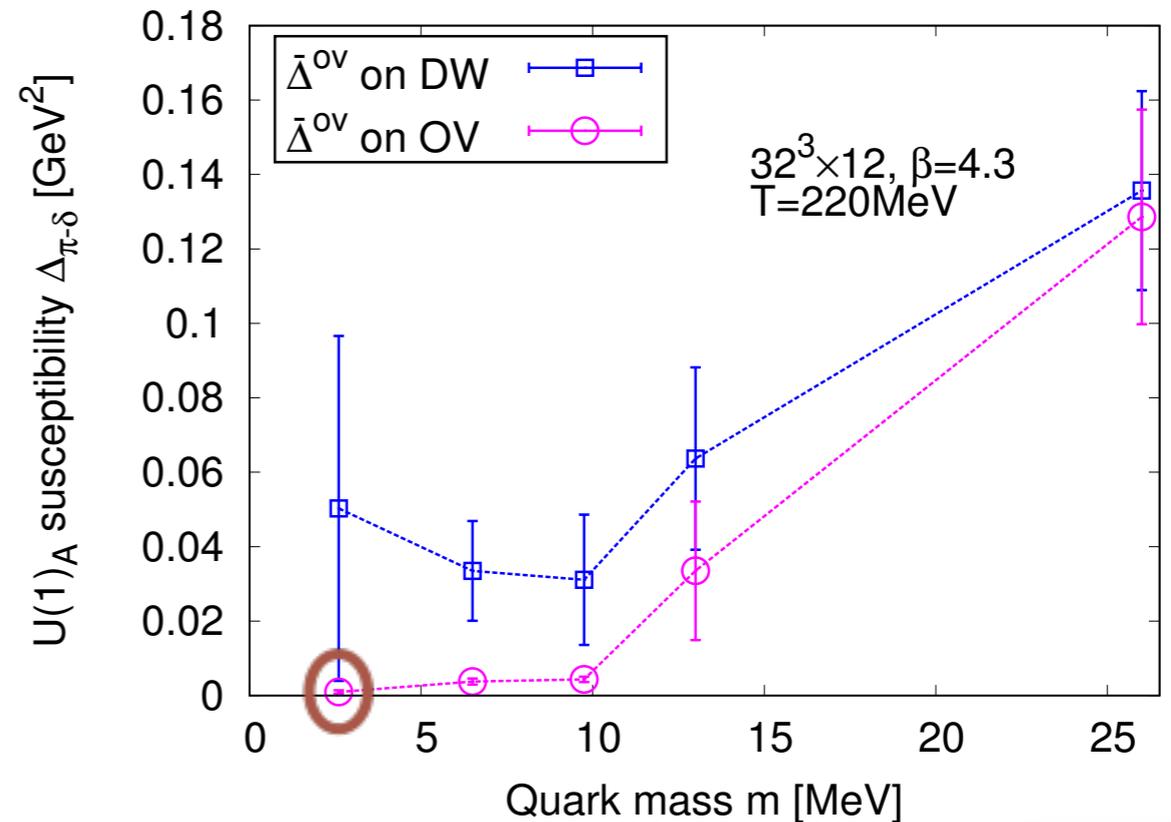
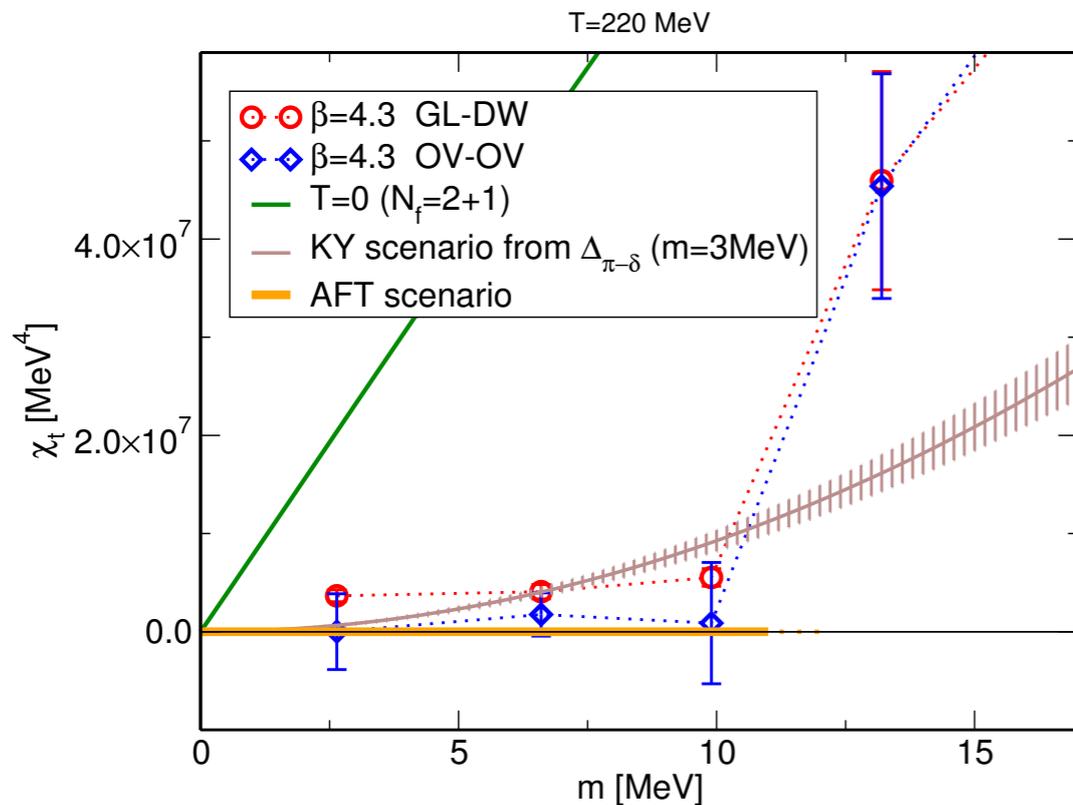


- KY scenario [Kanazawa, Yamamoto 2016]
  - $\Delta_{\pi-\delta}$ : including zero mode cont. is proper
  - $\Delta_{\pi-\delta} = \text{const} > 0$
  - $\Delta_{\pi-\delta} \approx 8 V f_A^2 m^2$  for **Q=0 sector** (for  $2V f_A m^2 < 1$ )
- $\Delta_{\pi-\delta}$  @ lightest point only from Q=0
- $\chi_t = 2 f_A m^2$
- inconsistent with  $m > 10$  MeV  $\chi_t$  growth



# competing scenarios for

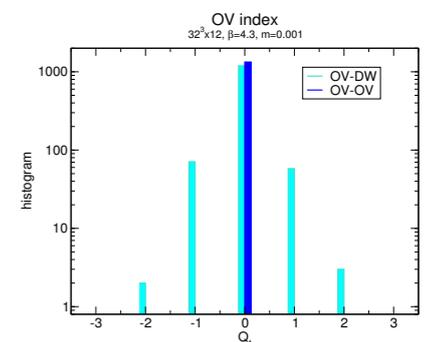
$\chi_t$  and  $\Delta_{\pi-\delta}$  ( $U_A(1)$  order parameter) @  $T \sim 220$  MeV



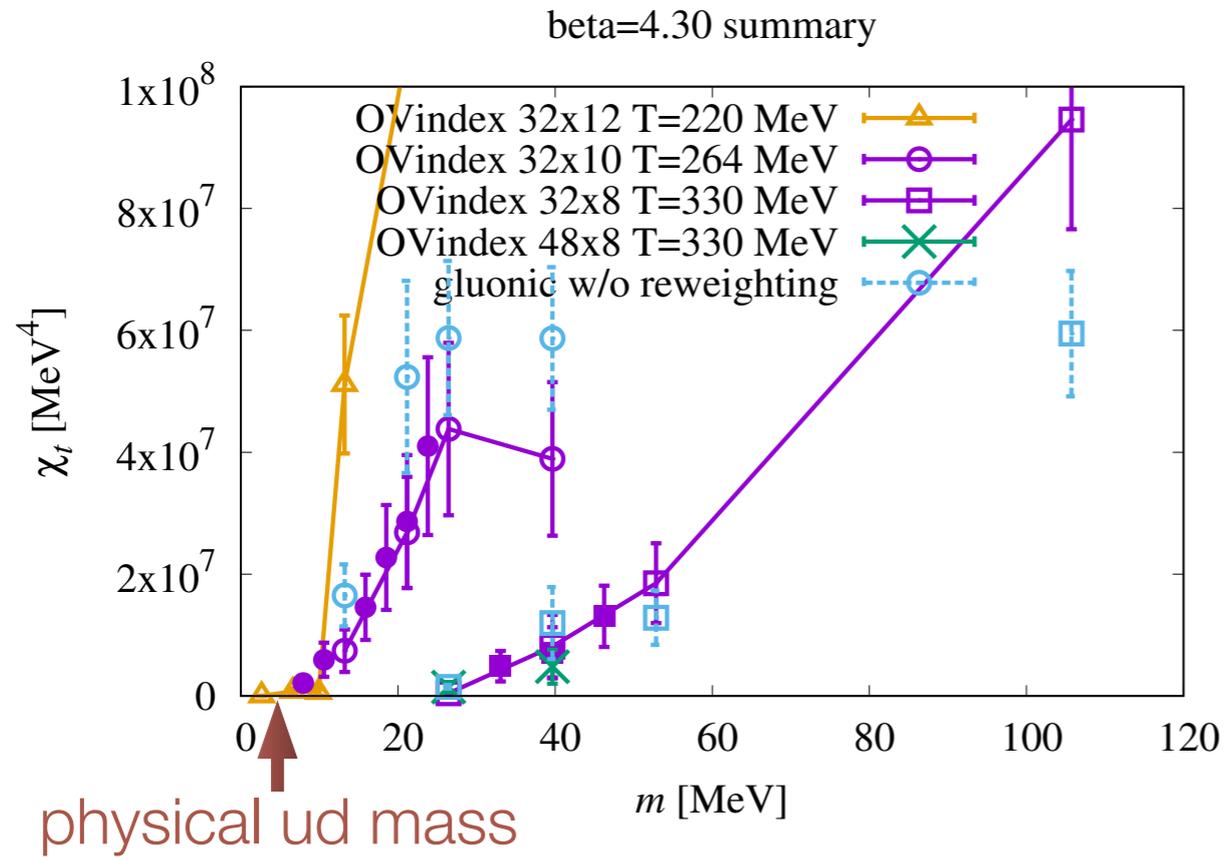
- KY scenario [Kobayashi, Yamamoto, 2016]

Volume study would be useful to check this

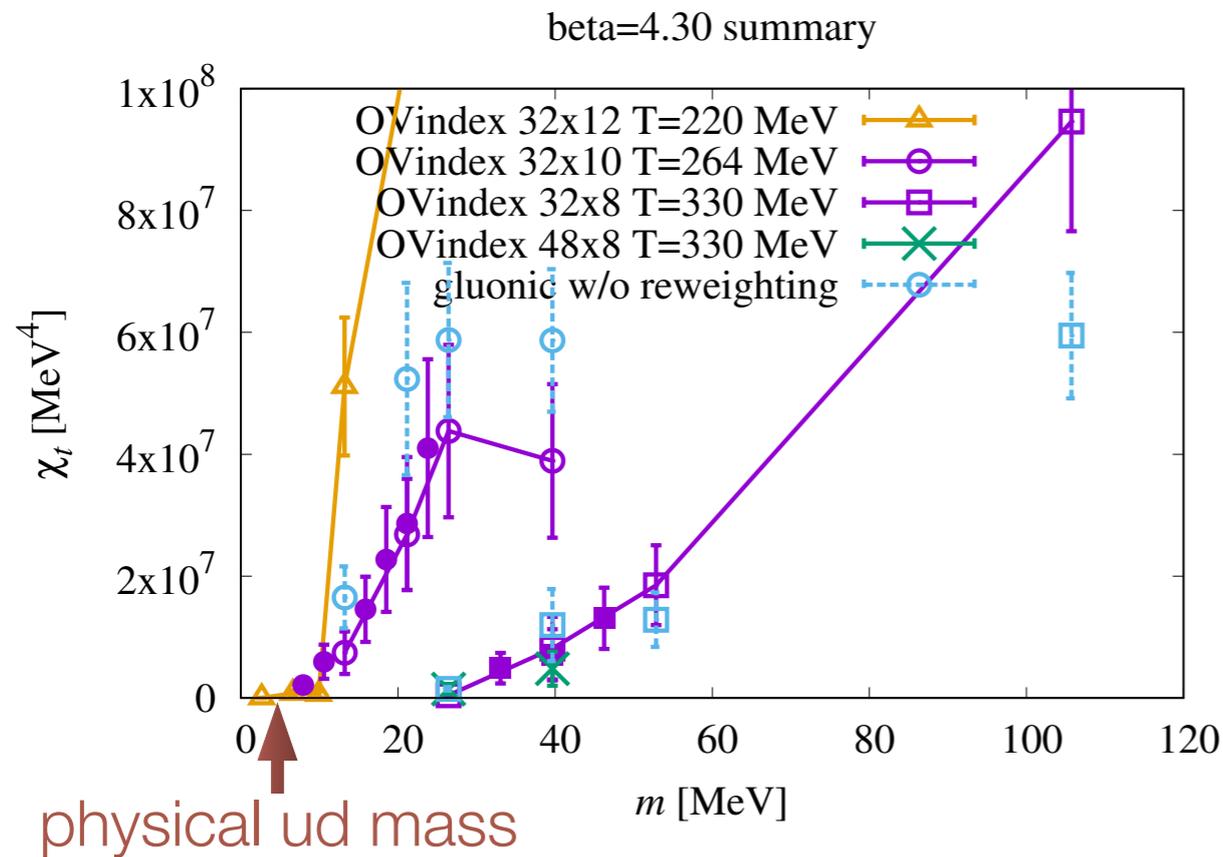
- $\Delta_{\pi-\delta}$ : incl
- $\Delta_{\pi-\delta} = \text{const} > 0$
- $\Delta_{\pi-\delta} \approx 8 V f_A^2 m^2$  for  $Q=0$  sector (for  $2V f_A m^2 < 1$ )
- $\Delta_{\pi-\delta}$  @ lightest point only from  $Q=0$
- $\chi_t = 2 f_A m^2$
- inconsistent with  $m > 10$  MeV  $\chi_t$  growth



T=220, 264 & 330 MeV ; 1/a=2.6 GeV fixed

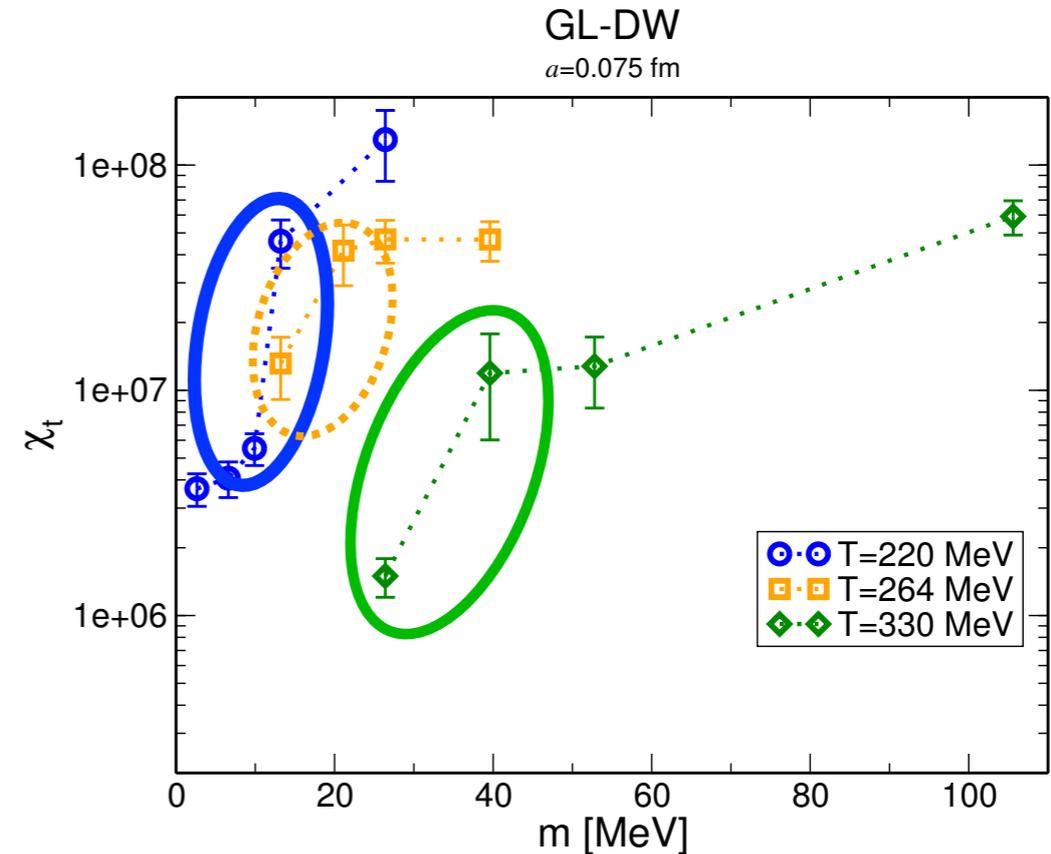
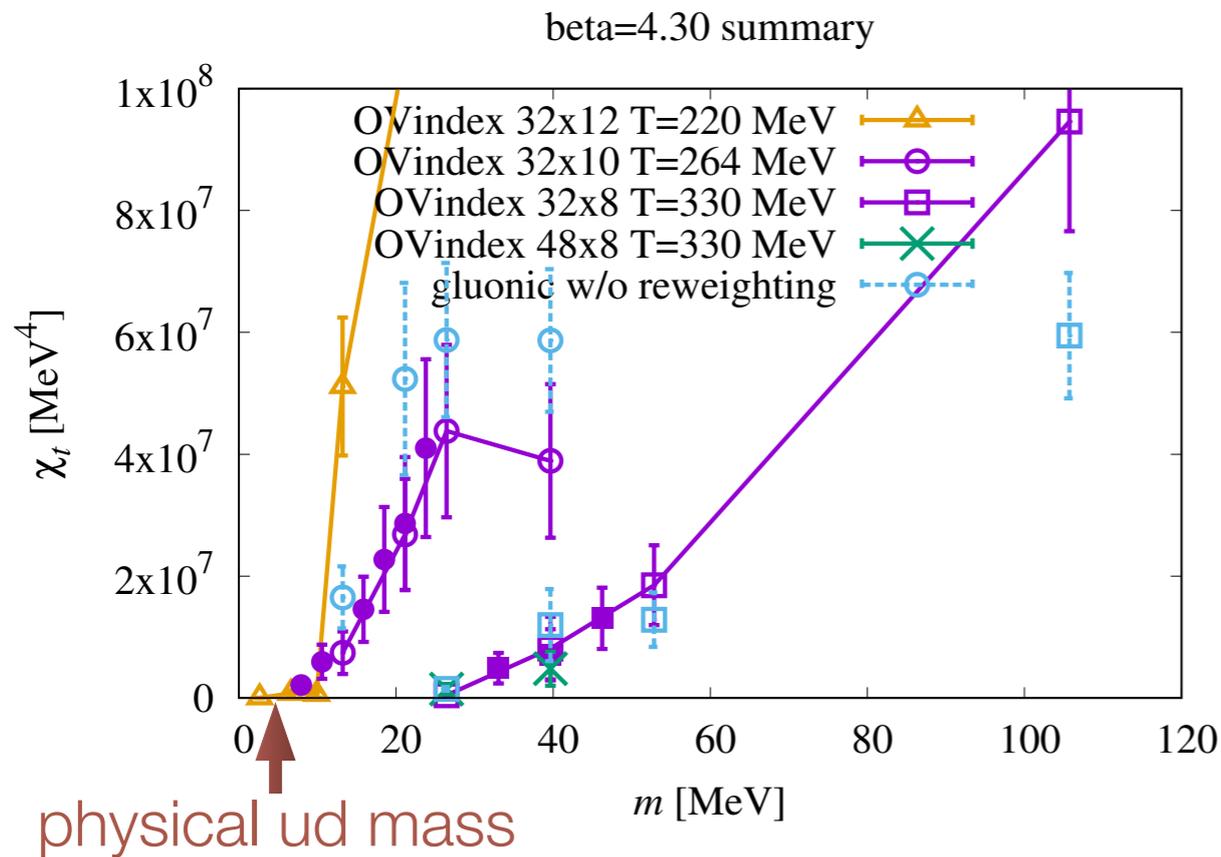


T=220, 264 & 330 MeV ; 1/a=2.6 GeV fixed



- filled symbol: by mass reweighting
- green: L x 1.5 (T=330 MeV) finite size effect small
- getting tiny value for larger mass with higher temperature
  - then lose the resolution
- similar drop as T=220 might have stated to be seen...

$T=220, 264 \text{ \& } 330 \text{ MeV} ; 1/a=2.6 \text{ GeV fixed}$



- filled symbol: by mass reweighting
- green:  $L \times 1.5$  (T=330 MeV) finite size effect small
- getting tiny value for larger mass with higher temperature
  - then lose the resolution
- similar drop as T=220 might have stated to be seen...

# Summary and Discussion

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- $N_f=2$  QCD under investigation for topological susceptibility  $\chi_t$  for high T
- fully utilizing the exact chiral symmetry Mobius DW (fine latt)  $\rightarrow$  OV reweighting
- Extensive study for  $T=220$  MeV with  $a=0.07$  fm (& 0.11 fm)
- $\chi_t$ : sudden grow for  $m>10$  MeV observed
- naive continuum scaling from two different definition of chit suggest it vanishes in the continuum limit for  $m<\sim 10$  MeV (eg: physical  $m_{ud}=\sim 4$  MeV)
- ➔ consistent with the prediction of AFT  $\chi_t=0$  for  $0\leq m<m_c$  (small on-zero  $m_c$ )
- caveat of this scaling analysis is the difference in physical volume
  - ➔ needs further investigation with different volume
- The numerical data shown here does not rule out KY scenario  $\rightarrow$  vol. study
- any case, sudden change of the  $\chi_t$  may suggest an existence of phase transition  $\rightarrow$  volume study is useful to check this. too.