

# Real-time CSFT simulations of chiral plasma with overlap fermions

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# Outline

- 1 Introduction
  - Motivation
  - Chiral Plasma Instability
- 2 Program
- 3 Numerical results
  - Helical modes
  - Comparison Wilson-Dirac and Overlap results
  - Axial charge

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# Why studying chiral plasma ?

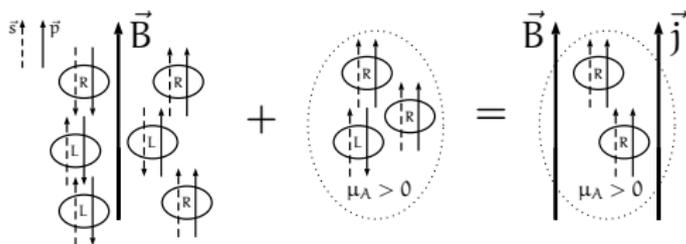
- Important in lot of different fields:
  - High energy physics:
    - Quark-Gluon Plasma
    - Hadronic matter
    - Cosmology of the early universe
  - Condensed matter phenomena
    - Weyl semi-metals
    - Topological insulators
    - Liquid Helium
- We understand very little about it beyond hydrodynamics and kinetic approaches

# Anomalous transport phenomena

- Chiral imbalance (axial charge  $\neq 0$ )  
 $\Rightarrow$  Anomalous transport phenomena!
- Chiral Separation Effect, Chiral Vortex effect, ...
- Chiral Magnetic Effect (CME):

$$\vec{j}_V = \sigma_{CME} \vec{B}$$

with  $\sigma_{CME} = qN_c\mu_A/2\pi^2$ : the CM conductivity



See: [axXiv:1511.04050](https://arxiv.org/abs/1511.04050)

# Anomalous Maxwell Equation

## Maxwell eq. + ohmic conductivity + CME

- $\partial_t \vec{B} = -\vec{\nabla} \times \vec{E}$
- $\partial_t \vec{E} = -\vec{\nabla} \times \vec{B} - \sigma \vec{E} - \sigma_{CME} \vec{B}$

Applied to plane waves:

## Maxwell eq. + ohmic conductivity + CME

- $i\omega \vec{B} = -i\vec{k} \times \vec{E}$
- $i\omega \vec{E} = -\vec{k} \times \vec{B} - \sigma \vec{E} - \sigma_{CME} \vec{B}$

Dispersion relation:

$$\omega = \frac{i\sigma}{2} \pm \sqrt{k^2 \pm \sigma_{CME} k - \frac{\sigma^2}{4}}$$

## Chiral Plasma instability

$$w = \frac{i\sigma}{2} \pm \sqrt{k^2 \pm \sigma_{CME} k - \frac{\sigma^2}{4}}$$

## Instable Solutions

- $E_x = fe^{\gamma t} \cos(kz)$        $E_y = -fe^{\gamma t} \sin(kz)$        $E_z = 0$
- $B_x = -f \frac{k}{\kappa} e^{\gamma t} \cos(kz)$        $B_y = f \frac{k}{\kappa} e^{\gamma t} \sin(kz)$        $B_z = 0$
- $\gamma = -\frac{\sigma}{2} + \sqrt{\frac{\sigma^2}{4} + \sigma_{CME} k - k^2}$

- For  $\sigma_{CME} > |k|$ :  $\gamma > 0$  and the amplitude of solution grows exponentially.
- What can stop the instability ? The decay of the axial charge

$$\partial_t Q_A = \frac{g^2}{2\pi^2} \int d^3x \vec{E} \cdot \vec{B} < 0 \Rightarrow \partial_t \sigma_{CME} < 0$$

# Different Approach

- Previous approach is far too simple as the  $k$  and  $w$  dependences of  $\sigma$  and  $\sigma_{CME}$  have been neglected
- Other approaches used so far:
  - Chiral kinetic theory (linear response, relaxation time, long - wavelength...)
  - Chiral Hydrodynamics (long - wavelength)
  - ...
- What has not been taken into account:
  - Non-trivial dispersion of conductivities
  - Developing (axial) charge inhomogeneities
  - Non-linear responses

⇒ We can use simulations

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## CSFT simulations

- CSFT: Classical Statistical Field Theory: Out of equilibrium simulation.
- Idea: Semi-classical approximation along Keldysh contour

more details in [Phys. Rev. D 94, 025009]

- In practice: solving the following set of equations:

$$\partial_t^2 A_{x,i}(t) = -J_{ext,x,i}(t) - \langle j_{x,i}(t) \rangle - \partial_j F_{x,ij}(t)$$

$$\langle j_{x,i}(t) \rangle = \sum_n n_f(\epsilon_n) \psi_{n,x}^\dagger(t) j_{x,i}(t) \psi_{n,x}(t)$$

$$\partial_t \psi_{n,x}(t) = -i h_{x,y} \psi_{n,y}(t)$$

$$j_{x,i} = \partial h_{x,y} / \partial A_{x,i}$$

with  $h_{x,y}$  the one particle Hamiltonian.

- Solving equation using Leapfrog algorithm (symplectic integrator: energy conservation under control)

# Wilson Dirac and Overlap Hamiltonian

Simulations with Wilson-Dirac and Overlap Hamiltonian.

- Wilson-Dirac Hamiltonian:

$$h_{x,y}^{WD}(m) = (3v_f\beta + m)\delta_{x,y} + \frac{iv_f}{2} \sum_{i=1}^3 (i\beta + \alpha_i) e^{igA_{x,i}} \delta_{y,x+\mathbf{e}_i} + \frac{iv_f}{2} \sum_{i=1}^3 (i\beta - \alpha_i) e^{-igA_{x-\mathbf{e}_i,i}} \delta_{y,x-\mathbf{e}_i}$$

- Overlap Hamiltonian:

$$h_{x,y}^{OV}(m) = \left(1 + \frac{m}{2\rho}\right) \beta + \left(1 - \frac{m}{2\rho}\right) \text{sign} \left( h_{x,y}^{WD}(\rho) \right)$$

- Important as chiral related problems.
- Approximation for the sign function: Polynomial and Zolotarev
- Comparison of computation time:

Exact  $\gg$  Polynomial  $>$  Zolotarev

# Introducing chiral imbalance and excitations

## Insertion of the chiral imbalance:

- left-handed eigenstates  $>$  right-handed eigenstates

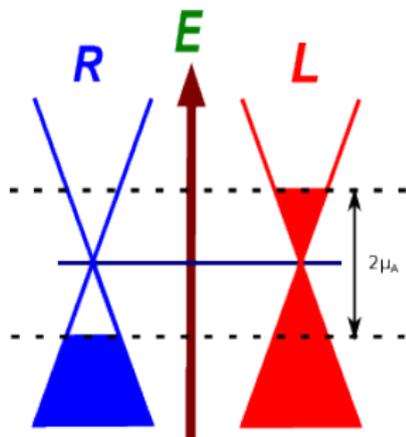
## Insertions of seeds of instabilities:

- $n$  modes

$$A_{x,i}(t=0) = \sum_{m=1}^n f \frac{\cos(k_m x_3 + \phi_m)}{\sqrt{4 \sin^2(k_m/2)}} \vec{n}_m$$

- Random polarization ( $\vec{n}_m \in \{\vec{x}, \vec{y}\}$ ), same energy (same  $f$ )
- To have instable mode:

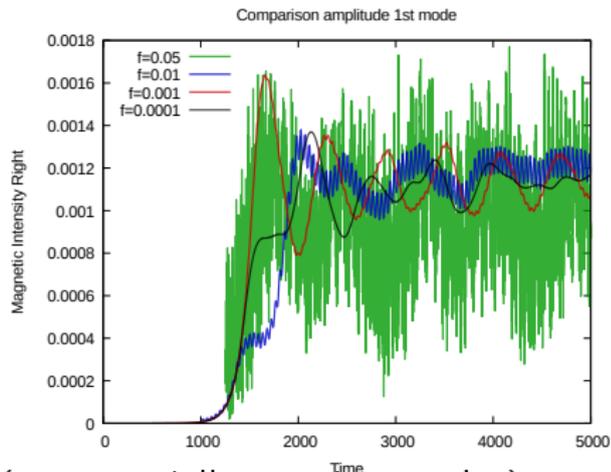
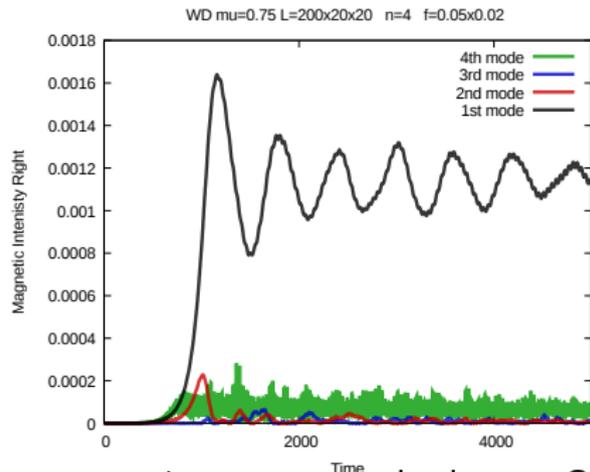
$$\sigma_{CME} > k \Rightarrow L > (4\pi^3) / \mu_A \quad (L > 125 \text{ for } \mu_A = 1)$$



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# Helical modes



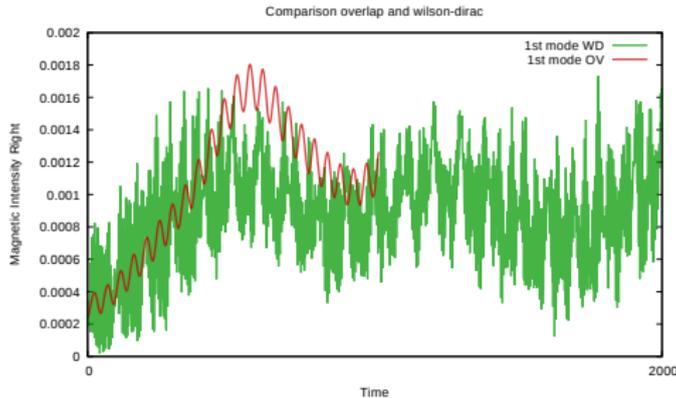
- Inverse cascade due to CPI (exponentially growing modes) with saturation!
- Saturation does not depend on amplitude of perturbations
- Reason for saturation?
  - WD artifacts?

$Q_A$  decays?

Non-linearity?

# Comparison Wilson-Dirac/Overlap

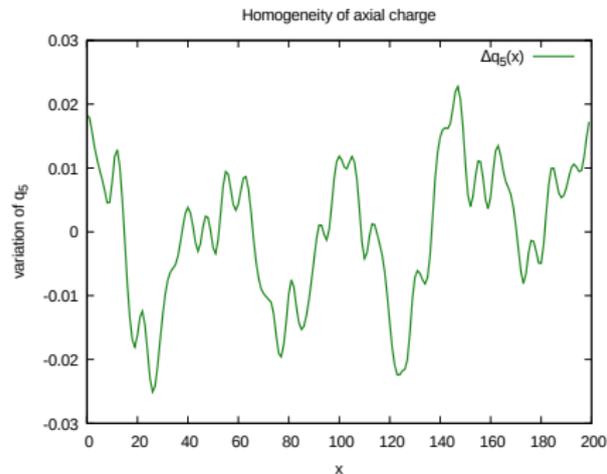
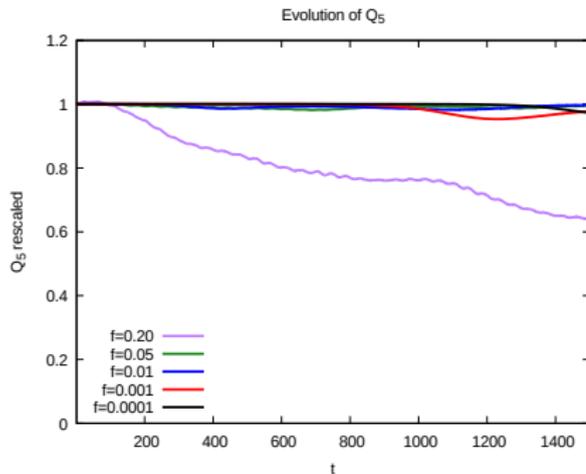
- Both behaved similarly (not a WD artifact)
  - Same inverse cascade (CPI)
  - Same saturation



- Possibility to study with WD (cheaper) with a check of the results with OV.

# Axial charge

- Decay only appears for large energy
  - ⇒ Cannot explain the Saturation !
  - ⇒ Strong non-linear effect !



- Homogeneity: Variation  $< 5\%$

# Conclusion

- Observation of saturation of the instabilities!
  - No Decay of the axial charge
- ⇒ Important non-linear phenomena: Simulation necessary!

## Outlook:

- Explanation of the saturation.
- Computing occupation number (production of particles).
- Understanding the behavior with analytical model.