

# QCD at finite isospin chemical potential: Reweighting

**Sebastian Schmalzbauer**

with Bastian Brandt & Gergely Endrődi



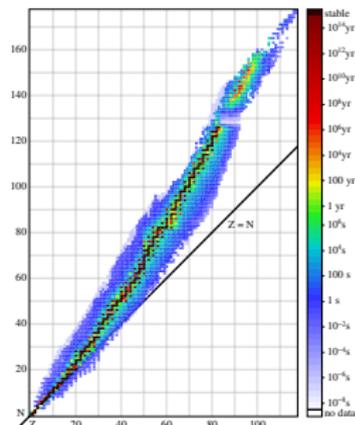
- QCD with isospin: motivation
- introduction of a symmetry breaking parameter  $\lambda$
- $\lambda \rightarrow 0$  limit
- explore phase diagram for  $\mu_B > 0$  via reweighting
- summary & outlook

# Motivation

- two quark flavors  $u, d$
- isospin density  $n_I = n_u - n_d$
- $n_I \neq 0$ :
  - systems with charged pions
  - neutron stars ( $udd$ )
  - heavy-ion collisions ( $N > Z$ )



(taken from [thescienceexplorer](#))



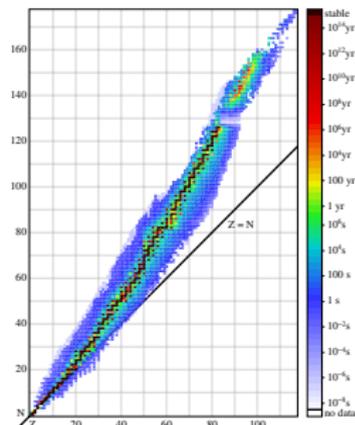
(taken from [wikipedia](#))

# Motivation

- two quark flavors  $u, d$
- isospin density  $n_I = n_u - n_d$
- $n_I \neq 0$ :
  - systems with charged pions
  - neutron stars ( $udd$ )
  - heavy-ion collisions ( $N > Z$ )
- no sign problem  $\Rightarrow$  lattice simulations
- analogies to baryon density
  - Silver Blaze
  - hadron condensation (small eigenvalues)
  - saturation



(taken from [thescienceexplorer](#))



(taken from [wikipedia](#))

# Symmetry breaking

- QCD with light quarks

$$\mathcal{M} = \not{D} + m_{ud}\mathbb{1}$$

- corresponding symmetry (without considering trivial  $U(1)_V$ )

$$SU(2)_V$$

# Symmetry breaking

- QCD with light quarks

$$\mathcal{M} = \not{D} + m_{ud}\mathbb{1} + \mu_1\gamma_0\tau_3$$

- corresponding symmetry (without considering trivial  $U(1)_V$ )

$$SU(2)_V \rightarrow U(1)_{\tau_3}$$

- **problem:** cannot directly observe spontaneous symmetry breaking
  - pion condensate  $\langle \bar{\psi}\gamma_5\tau_{1,2}\psi \rangle = 0$  (Goldstone mode)
  - accumulation of zero eigenvalues

# Symmetry breaking

- QCD with light quarks

$$\mathcal{M} = \not{D} + m_{ud}\mathbb{1} + \mu_1\gamma_0\tau_3 + i\lambda\gamma_5\tau_2$$

- corresponding symmetry (without considering trivial  $U(1)_V$ )

$$SU(2)_V \rightarrow U(1)_{\tau_3} \rightarrow \emptyset$$

- **problem**: cannot directly observe spontaneous symmetry breaking
  - pion condensate  $\langle \bar{\psi}\gamma_5\tau_{1,2}\psi \rangle = 0$  (Goldstone mode)
  - accumulation of zero eigenvalues
- **solution**: add explicit unphysical breaking
  - can directly observe spontaneous symmetry breaking
  - no zero eigenvalues

# Symmetry breaking

- QCD with light quarks

$$\mathcal{M} = \not{D} + m_{ud}\mathbb{1} + \mu\nu\gamma_0\tau_3 + i\lambda\gamma_5\tau_2$$

- corresponding symmetry (without considering trivial  $U(1)_V$ )

$$SU(2)_V \rightarrow U(1)_{\tau_3} \rightarrow \emptyset$$

- **problem**: cannot directly observe spontaneous symmetry breaking
  - pion condensate  $\langle \bar{\psi}\gamma_5\tau_{1,2}\psi \rangle = 0$  (Goldstone mode)
  - accumulation of zero eigenvalues
- **solution**: add explicit unphysical breaking
  - can directly observe spontaneous symmetry breaking
  - no zero eigenvalues
- extrapolate results  $\lambda \rightarrow 0$

# Simulation Details

- QCD partition function

$$\mathcal{Z} = \int \mathcal{D}[U] (\det \mathcal{M})^{\frac{1}{4}} e^{-S_G}$$

- $N_f = 2$  rooted staggered fermions,  $\eta_5 = (-1)^{n_t+n_x+n_y+n_z}$

$$\mathcal{M} = \begin{pmatrix} \not{D}_{\mu_l} + m & \lambda \eta_5 \\ -\lambda \eta_5 & \not{D}_{-\mu_l} + m \end{pmatrix}$$

- plaquette gauge action  $S_G$
- positivity of determinant:

$$\det \mathcal{M} = \det \left( M^\dagger M + \lambda^2 \right) \in \mathbb{R}_{>0} \quad M = \not{D}_{\mu_l} + m$$

ensured by  $\eta_5 \tau_1 \mathcal{M} \tau_1 \eta_5 = \mathcal{M}^\dagger$

- first studies [[Kogut, Sinclair '02](#), [de Forcrand, Stephanov, Wenger '07](#)]
- $8^4$  lattice,  $T = 0$

- condensates, isospin density

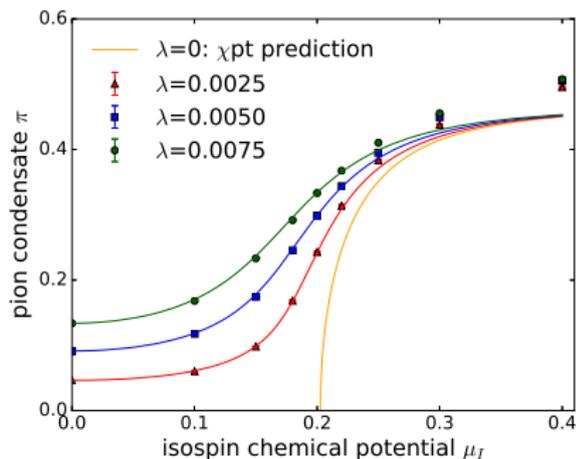
$$\langle \bar{\psi}\psi \rangle = \frac{T}{V} \frac{\partial \ln \mathcal{Z}}{\partial m} \quad \langle \pi \rangle = \frac{T}{V} \frac{\partial \ln \mathcal{Z}}{\partial \lambda} \quad \langle n_I \rangle = \frac{T}{V} \frac{\partial \ln \mathcal{Z}}{\partial \mu_I}$$

- condensates, isospin density

$$\langle \bar{\psi}\psi \rangle = \frac{T}{V} \frac{\partial \ln \mathcal{Z}}{\partial m} \quad \langle \pi \rangle = \frac{T}{V} \frac{\partial \ln \mathcal{Z}}{\partial \lambda} \quad \langle n_I \rangle = \frac{T}{V} \frac{\partial \ln \mathcal{Z}}{\partial \mu_I}$$

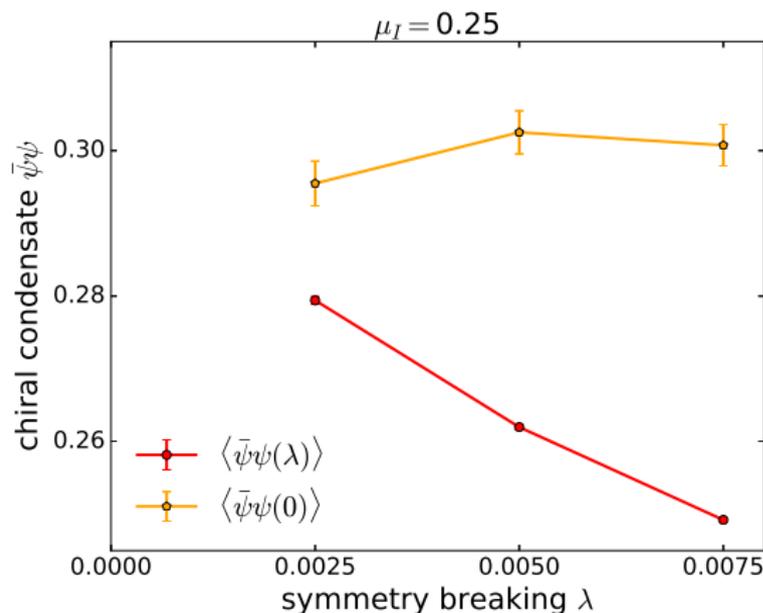
- $O(\lambda)$  strongly depend on  $\lambda$
- for  $\lambda \rightarrow 0$ :

$$\langle \pi \rangle \begin{cases} = 0 & : \mu_I \leq m_\pi/2 \\ > 0 & : \mu_I > m_\pi/2 \end{cases}$$



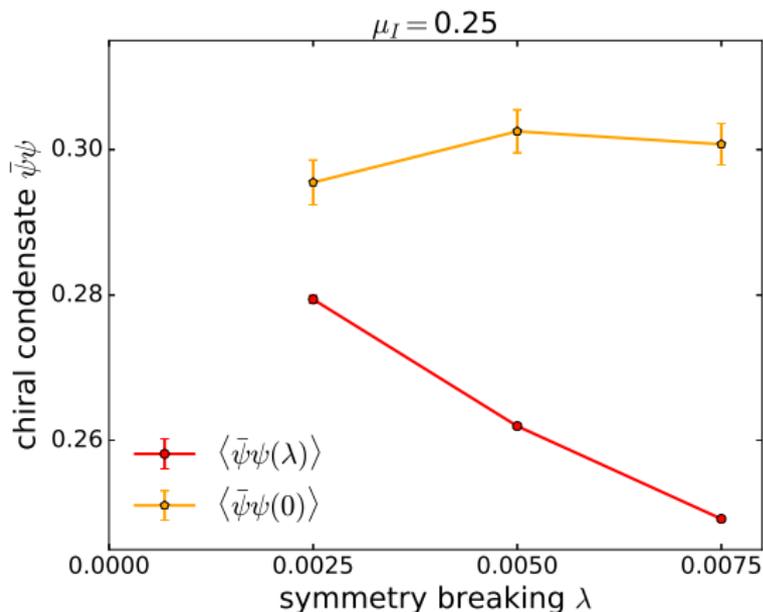
# $\lambda$ -extrapolation

- naive limit (steep):  $\langle O \rangle = \lim_{\lambda \rightarrow 0} \int \mathcal{D}[U] e^{-S_G} \det [\mathcal{M}(\lambda)]^{\frac{1}{4}} O(\lambda)$
- better (almost flat):  $\langle O \rangle = \lim_{\lambda \rightarrow 0} \int \mathcal{D}[U] e^{-S_G} \det [\mathcal{M}(\lambda)]^{\frac{1}{4}} O(0)$



# $\lambda$ -extrapolation

- naive limit (steep):  $\langle O \rangle = \lim_{\lambda \rightarrow 0} \int \mathcal{D}[U] e^{-S_G} \det [\mathcal{M}(\lambda)]^{\frac{1}{4}} O(\lambda)$
- better (almost flat):  $\langle O \rangle = \lim_{\lambda \rightarrow 0} \int \mathcal{D}[U] e^{-S_G} \det [\mathcal{M}(\lambda)]^{\frac{1}{4}} O(0)$



- instead of taking lim:  
reweight  $\lambda \rightarrow 0$

# Reweighting in $\lambda$

- reweighting factor

$$R_\lambda = \left[ \frac{\det(M^\dagger M)}{\det(M^\dagger M + \lambda^2)} \right]^{\frac{1}{4}} \in \mathbb{R}$$

- leading order expansion in  $\lambda$

$$R_\lambda \approx R_{LO} = \exp\left(-\frac{\lambda T}{2V}\pi\right)$$

# Reweighting in $\lambda$

- reweighting factor

$$R_\lambda = \left[ \frac{\det(M^\dagger M)}{\det(M^\dagger M + \lambda^2)} \right]^{\frac{1}{4}} \in \mathbb{R}$$

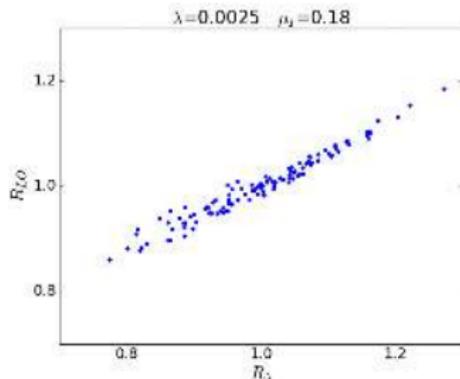
- leading order expansion in  $\lambda$

$$R_\lambda \approx R_{LO} = \exp\left(-\frac{\lambda T}{2V}\pi\right)$$

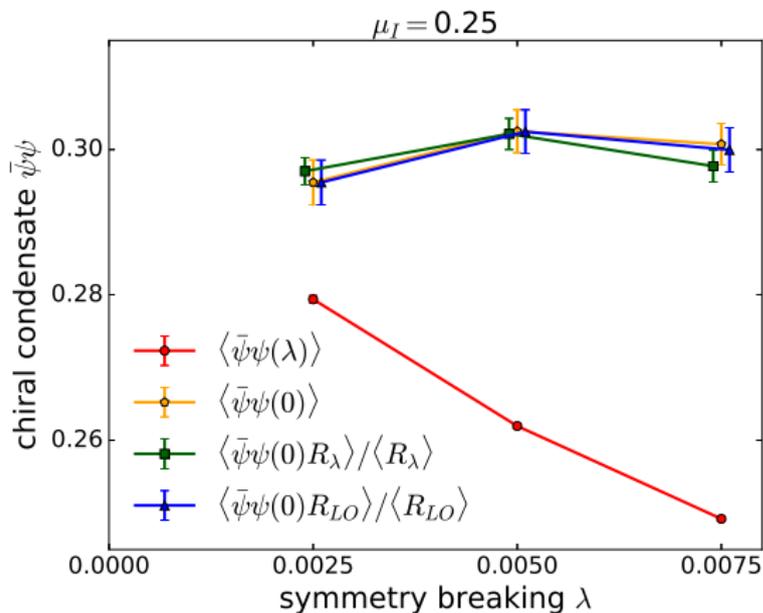
- any observable measured at  $\lambda = 0$  can then be expressed as:

$$\langle O \rangle = \frac{\langle OR_\lambda \rangle}{\langle R_\lambda \rangle} \approx \frac{\langle OR_{LO} \rangle}{\langle R_{LO} \rangle}$$

- $R_\lambda$  expensive
- $R_{LO}$  cheap
- for  $\lambda \rightarrow 0$ :  $R_\lambda = R_{LO}$



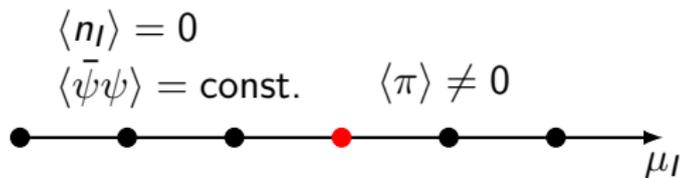
# Reweighting in $\lambda$ : comparison



- $O(\lambda)$ : steep
- $O(0)$ : flat
- LO  $\hat{=}$  full reweighting

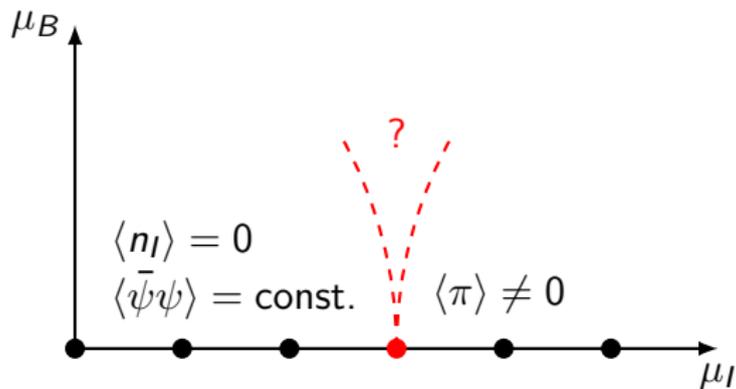
## Reweighting in $\mu$

# Reweighting in $\mu$ : motivation



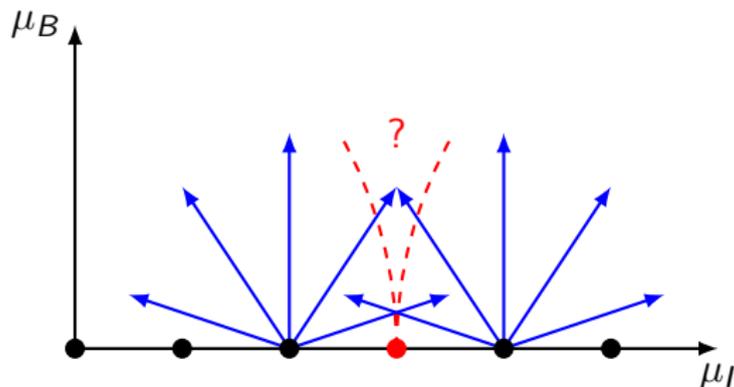
- until now:  $\mu_B = 0$

# Reweighting in $\mu$ : motivation



- until now:  $\mu_B = 0$
- want to explore the  $\mu_B > 0$  direction: **phase boundary**

# Reweighting in $\mu$ : motivation



- until now:  $\mu_B = 0$
- want to explore the  $\mu_B > 0$  direction: **phase boundary**
- **reweight** from pure  $\mu_I$  into  $\mu_I - \mu_B$  plane
- look for:
  - lines of constant observables
  - small error on reweighted observables (big overlap, small sign problem)

## Reweighting in $\mu$ : how-to

ensemble:  $(\mu_I, \lambda, 0) \rightarrow (\mu_I, 0, 0) \rightarrow (\mu'_I, 0, \mu_B)$

## Reweighting in $\mu$ : how-to

ensemble:  $(\mu_I, \lambda, 0) \rightarrow (\mu_I, 0, 0) \rightarrow (\mu'_I, 0, \mu_B)$

$$\mathcal{M} = \begin{pmatrix} M_{\mu_I} & \eta_5 \lambda \\ -\eta_5 \lambda & M_{-\mu_I} \end{pmatrix} \rightarrow \begin{pmatrix} M_{\mu_I} & 0 \\ 0 & M_{-\mu_I} \end{pmatrix} \rightarrow \begin{pmatrix} M_{\mu_u} & 0 \\ 0 & M_{\mu_d} \end{pmatrix}$$

with  $\mu_u = \mu_B + \mu'_I$      $\mu_d = \mu_B - \mu'_I$

## Reweighting in $\mu$ : how-to

ensemble:  $(\mu_I, \lambda, 0) \rightarrow (\mu_I, 0, 0) \rightarrow (\mu'_I, 0, \mu_B)$

$$\mathcal{M} = \begin{pmatrix} M_{\mu_I} & \eta_5 \lambda \\ -\eta_5 \lambda & M_{-\mu_I} \end{pmatrix} \rightarrow \begin{pmatrix} M_{\mu_I} & 0 \\ 0 & M_{-\mu_I} \end{pmatrix} \rightarrow \begin{pmatrix} M_{\mu_u} & 0 \\ 0 & M_{\mu_d} \end{pmatrix}$$

with  $\mu_u = \mu_B + \mu'_I$      $\mu_d = \mu_B - \mu'_I$

- can now factorize

$$\det \mathcal{M} = \det M(\mu_u) \det M(\mu_d)$$

- reweighting factor for  $(\mu_I, 0, 0) \rightarrow (\mu'_I, 0, \mu_B)$ :

$$R_\mu = \left[ \frac{\det M(\mu_u) \det M(\mu_d)}{\det M(\mu_I) \det M(-\mu_I)} \right]^{\frac{1}{4}} \in \mathbb{C}$$

# Reweighting in $\mu$ : determinant reduction

- need to compute  $\det M(\mu)$  for many different  $\mu$
- determinant reduction [Toussaint '90 & Fodor, Katz '02]

$$\begin{aligned}\det M(\mu) &= e^{-3V_s L_t \mu} \det \left( P - e^{L_t \mu} \right) \\ &= e^{-3V_s L_t \mu} \prod_{i=1}^{6V_s} \left( p_i - e^{L_t \mu} \right)\end{aligned}$$

- analytic  $\mu$ -dependence: calculate eigenvalues of  $P$  **just once**
- construction of  $P$  is well defined

## Reweighting in $\mu$ : observables

- combined reweighting step

$$\langle O \rangle = \frac{\langle OR_\lambda R_\mu \rangle}{\langle R_\lambda R_\mu \rangle}$$

# Reweighting in $\mu$ : observables

- combined reweighting step

$$\langle O \rangle = \frac{\langle OR_\lambda R_\mu \rangle}{\langle R_\lambda R_\mu \rangle}$$

- have to measure observables in target ensemble ( $\mu_B \neq 0$ ):

$$\bar{\psi}\psi = \frac{T}{4V} \frac{\partial}{\partial m} [\ln \det M(\mu_u) + \ln \det M(\mu_d)]$$

via numerical derivative

## Reweighting in $\mu$ : observables

- combined reweighting step

$$\langle O \rangle = \frac{\langle OR_\lambda R_\mu \rangle}{\langle R_\lambda R_\mu \rangle}$$

- have to measure observables in target ensemble ( $\mu_B \neq 0$ ):

$$\bar{\psi}\psi = \frac{T}{4V} \frac{\partial}{\partial m} [\ln \det M(\mu_u) + \ln \det M(\mu_d)]$$

via numerical derivative,

$$n_l = \frac{T}{4V} \frac{\partial}{\partial \mu_l} [\ln \det M(\mu_u) + \ln \det M(\mu_d)]$$

via closed formula (follows from determinant reduction):

$$\frac{\partial \ln \det M(\mu)}{\partial \mu} = -3L_t V_s - L_t \sum_{i=1}^{6V_s} \frac{e^{L_t \mu}}{p_i - e^{L_t \mu}}$$

## Reweighting in $\mu$ : observables

- combined reweighting step

$$\langle O \rangle = \frac{\langle OR_\lambda R_\mu \rangle}{\langle R_\lambda R_\mu \rangle}$$

- have to measure observables in target ensemble ( $\mu_B \neq 0$ ):

$$\bar{\psi}\psi = \frac{T}{4V} \frac{\partial}{\partial m} [\ln \det M(\mu_u) + \ln \det M(\mu_d)]$$

via numerical derivative,

$$n_l = \frac{T}{4V} \frac{\partial}{\partial \mu_l} [\ln \det M(\mu_u) + \ln \det M(\mu_d)]$$

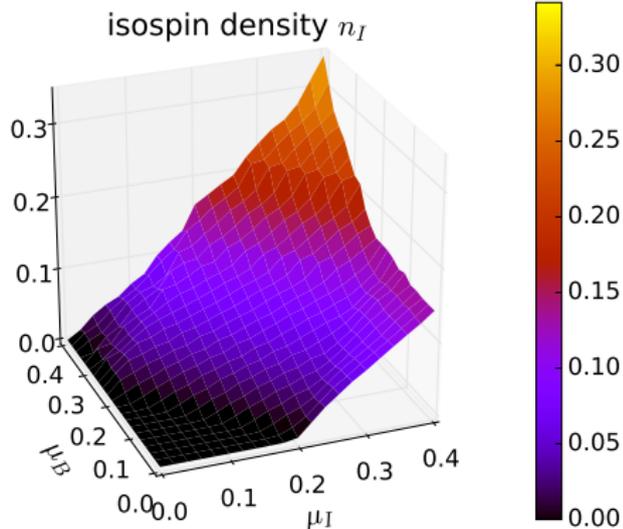
via closed formula (follows from determinant reduction):

$$\frac{\partial \ln \det M(\mu)}{\partial \mu} = -3L_t V_s - L_t \sum_{i=1}^{6V_s} \frac{e^{L_t \mu}}{p_i - e^{L_t \mu}}$$

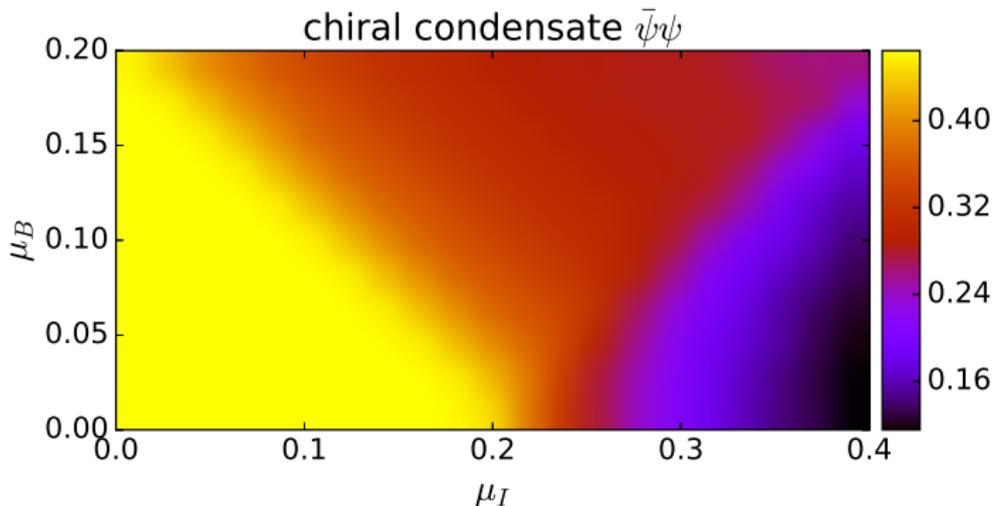
- measure  $\pi(\mu_B \neq 0)$ ?

# Reweighting in $\mu$ : results ( $T = 0$ )

- take auxiliary ensemble with biggest overlap
- unphysical surfaces
  - severe sign problem
  - underestimated error
- similar for chiral condensate



# Reweighting in $\mu$ : results ( $T = 0$ )



- phase boundary bends to the right for increasing  $\mu_B$

# Summary, outlook & transition

- lattice QCD with isospin chemical potential
- observe pion condensation with  $\lambda \neq 0$
- reweight  $\lambda \rightarrow 0$
  
- determinant reduction + reweighting to explore the  $\mu_I - \mu_B$  phase-space
  
- finite T: next talk by B. Brandt

