

Complex Langevin simulation of
QCD at finite density and low temperature
using the deformation technique

Shinji Shimasaki (Keio U.)

Collaborators

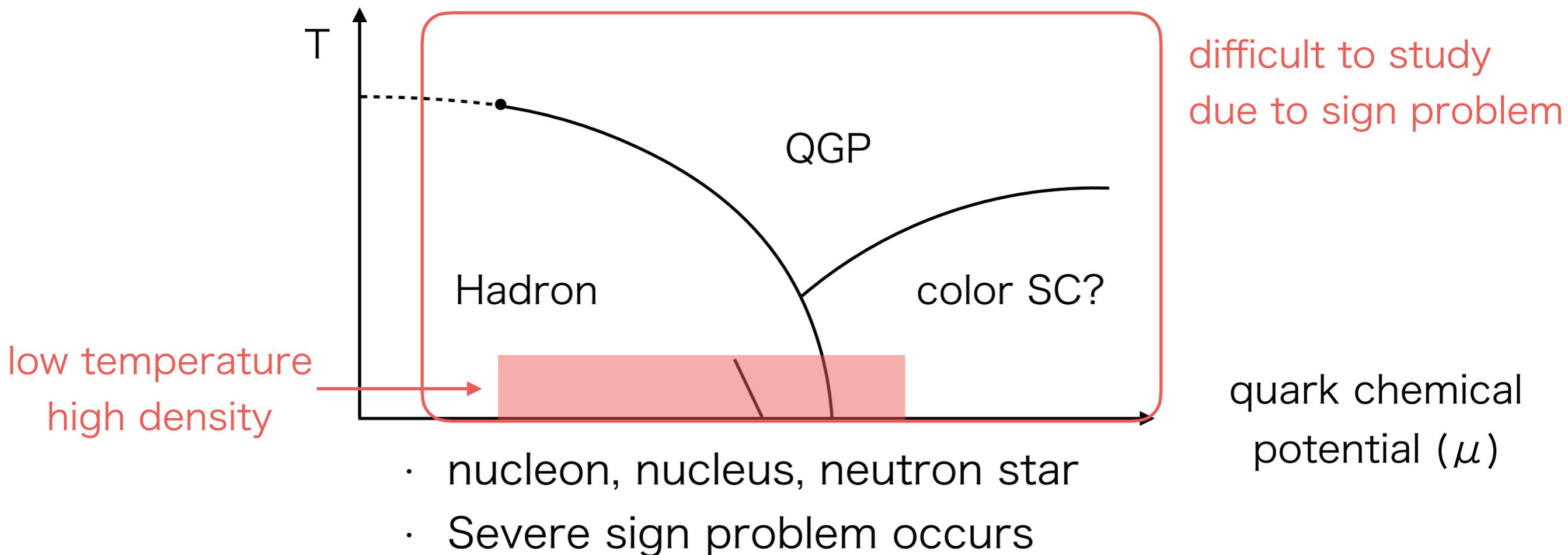
Keitaro Nagata (Kochi U.)

Jun Nishimura (KEK, SOKENDAI)

Sign problem in finite density QCD

- QCD partition function $Z = \int dU \det M[U, \mu] e^{-S_g[U]}$
- QCD phase diagram

complex ($\mu \neq 0$)



Our work

- We study QCD at finite μ and low T ($N_f=4$ staggered fermion) by using **the complex Langevin method (CLM)**
[Parisi 83][Klauder 84]
[Aarts, Seiler, Stamatescu 09][Aarts, James, Seiler, Stamatescu 11]...
- We successfully observe “the Silver Blaze phenomenon”
- Crucial techniques in our work :
 - ✓ **Gauge cooling** [Seiler, Sexty, Stamatescu 13]
 - ✓ **Deformation of the Dirac operator** [Ito, Nishimura 16]
 - ✓ **Reliability check of the obtained results by the probability distribution of the drift term** [Nagata, Nishimura, SS 16]

CLM for lattice QCD

- $Z = \int dU \det M[U, \mu] e^{-S_g[U]}$

- complexification $U_{x\mu} \in SU(3) \rightarrow \mathcal{U}_{x\mu} \in SL(3, C)$

- action $S[U] \rightarrow S[\mathcal{U}]$ holomorphic function

- gauge transformation $\mathcal{U}_{x\mu} \rightarrow g_x \mathcal{U}_{x\mu} g_{x+\hat{\mu}}^{-1}$ $g_x \in SL(3, C)$

- complex Langevin equation

$$\mathcal{U}_{x\mu}(t + \epsilon) = \exp \left\{ i \left(\epsilon v_{x\mu}(\mathcal{U}) + \sqrt{\epsilon} \eta_{x\mu} \right) \right\} \mathcal{U}_{x\mu}(t)$$

$v_{x\mu}(\mathcal{U}) = -\mathcal{D}_{x\mu} S[\mathcal{U}]$

Condition for the correct convergence

[Nagata, Nishimura, SS 16]

- Probability distribution of the magnitude of the drift term

$$p(u) = \frac{1}{4N_V} \left\langle \sum_{x\nu} \delta(u - u_{x\nu}) \right\rangle \quad u_{x\nu} = \sqrt{\frac{1}{N_c^2 - 1} \text{tr}(v_{x\nu} v_{x\nu}^\dagger)}$$

- Asymptotic behavior of $p(u)$ at large u

exponential fall-off or faster	→	correct result
power-law fall-off	→	wrong result

- The main causes of the power-law fall

- large deviations of $\mathcal{U}_{x\mu}$ from SU(3)
(excursion problem)

← Gauge cooling

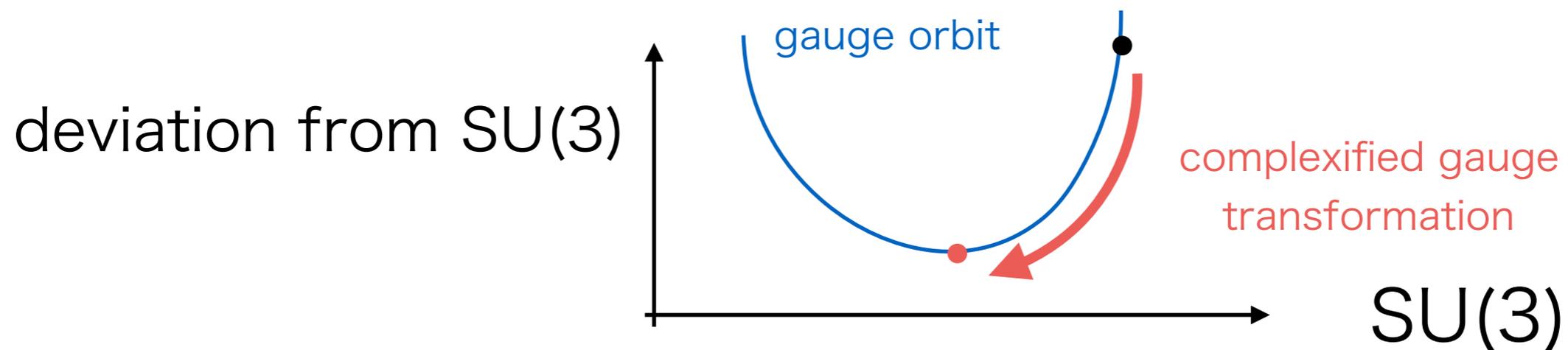
- small eigenvalues of $M[\mathcal{U}, \mu]$
(singular drift problem)

← Deformation of
the Dirac operator

Gauge cooling

[Seiler, Sexty, Stamatescu 13]

- Perform the complexified gauge transformation after each Langevin step in such a way that the deviation of $\mathcal{U}_{x\mu}$ from SU(3) is minimized



⌘ theoretical justification [Nagata, Nishimura, SS 16]

Deformation of the Dirac operator

[Ito, Nishimura 16]

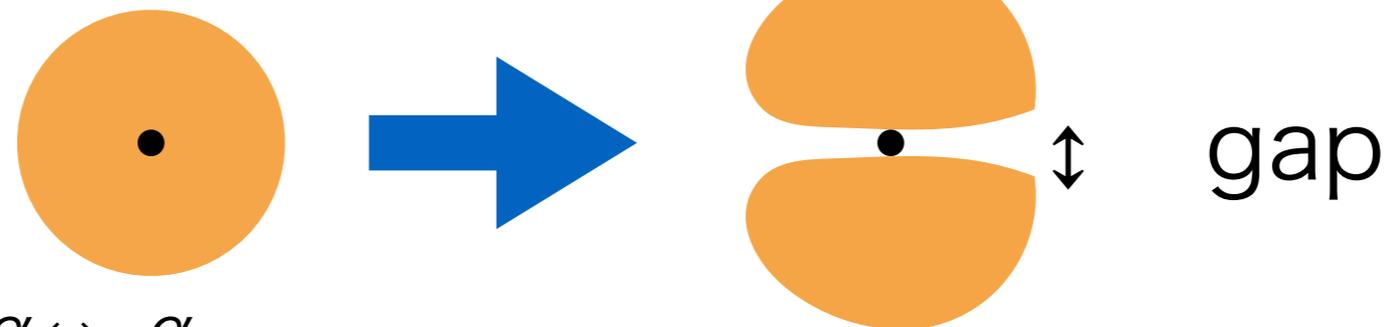
Next talk by Y. Ito

- $\det[M(\mathcal{U}, \mu)] \rightarrow \det[M(\mathcal{U}, \mu) + i\alpha\gamma_4]$

- Formally, equivalent to adding the imaginary chemical potential in the continuum theory

→ Chiral sym and rotation sym kept intact

- $\alpha \neq 0 \rightarrow$ a gap in the imaginary direction



- symmetric under $\alpha \leftrightarrow -\alpha$

⇒ We fit the expectation values of observables to $c_0 + c_2 \alpha^2 + \dots$ by using only data points, which pass the reliability test, and extrapolate $\alpha \rightarrow 0$.

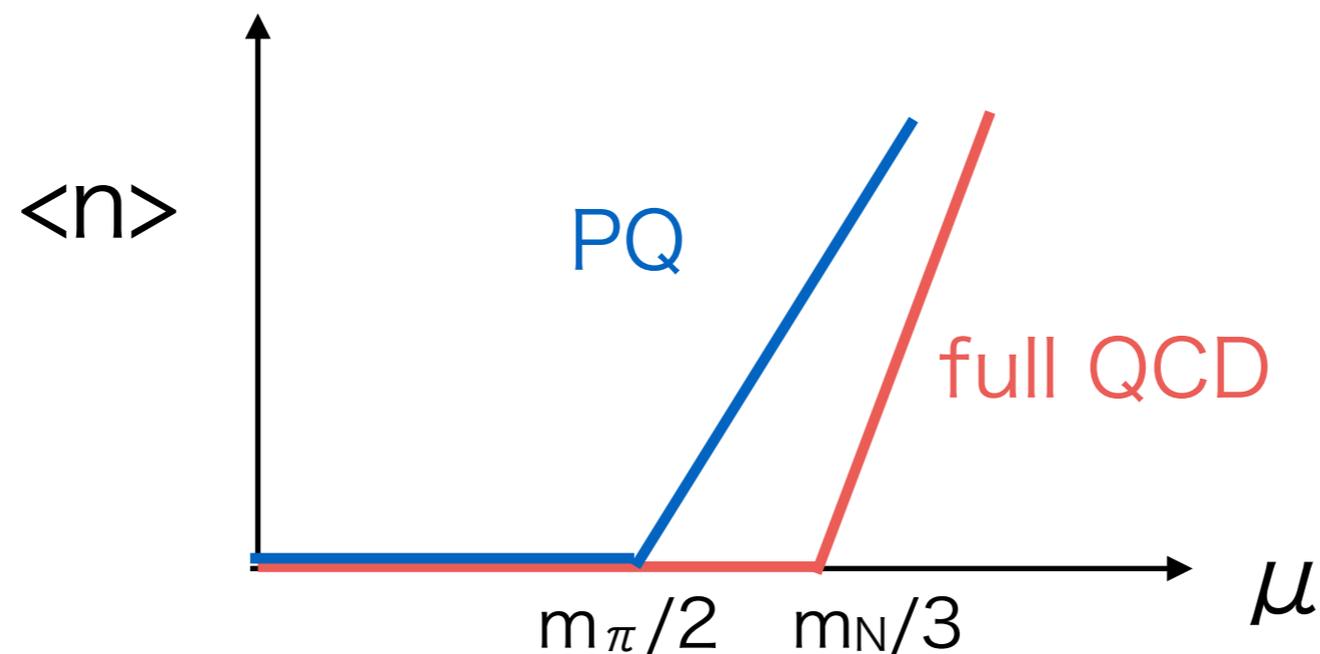
Simulation setup

cf) related talk by D. K. Sinclair

- Lattice setup :
 - * Staggered fermion ($N_f=4$)
 - * lattice size : $N_x=N_y=N_z=4$, $N_t=8$ [Nagata, Nishimura, SS, to appear]
 - * $\beta=5.7$
 - * quark mass : $ma=0.05$
 - * quark chemical potential : $\mu a=0.4, 0.5, 0.6, 0.7$
- CLM setup :
 - * $\varepsilon=10^{-4}$,
 - * total Langevin steps = $5 \times 10^5 \sim 15 \times 10^5$
 - * gauge cooling (10~50 times)
- We also report some preliminary results for $N_x=N_y=N_z=8$, $N_t=16$
[Ito, Matsufuru, Moritake, Nishimura, SS, Tsuchiya, Tsutsui]

Observables

- Baryon number density $\langle n \rangle = \frac{1}{N_V N_c} \frac{\partial}{\partial(\mu a)} \log Z$
- Chiral condensate $\langle \Sigma \rangle = \frac{1}{N_V} \frac{\partial}{\partial(ma)} \log Z$
- We compare the results with the RHMC results of the phase quenched model to confirm “the Silver Blaze phenomenon”.



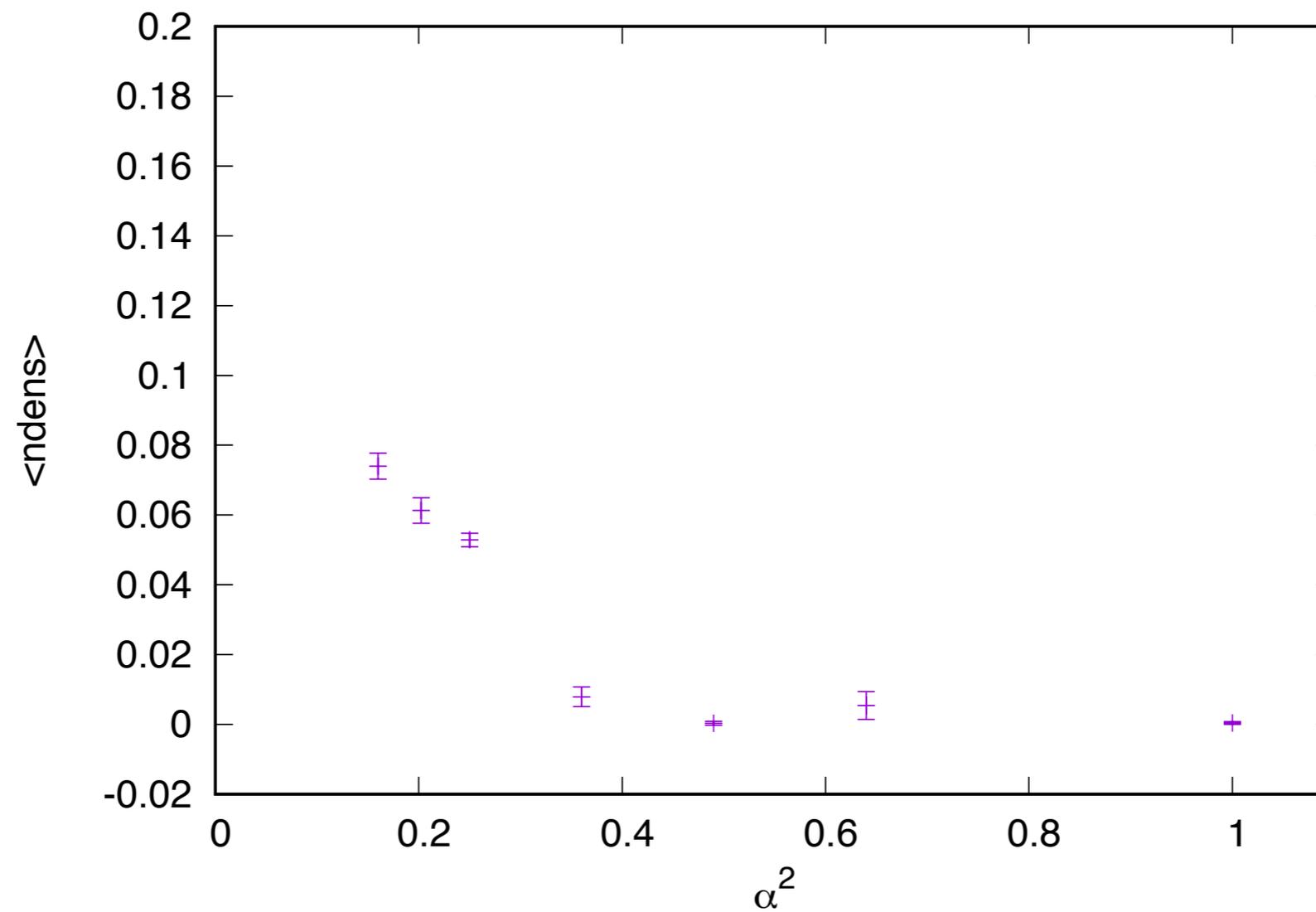
Results for

$$N_x=N_y=N_z=4, N_t=8$$

[Nagata, Nishimura, SS, to appear]

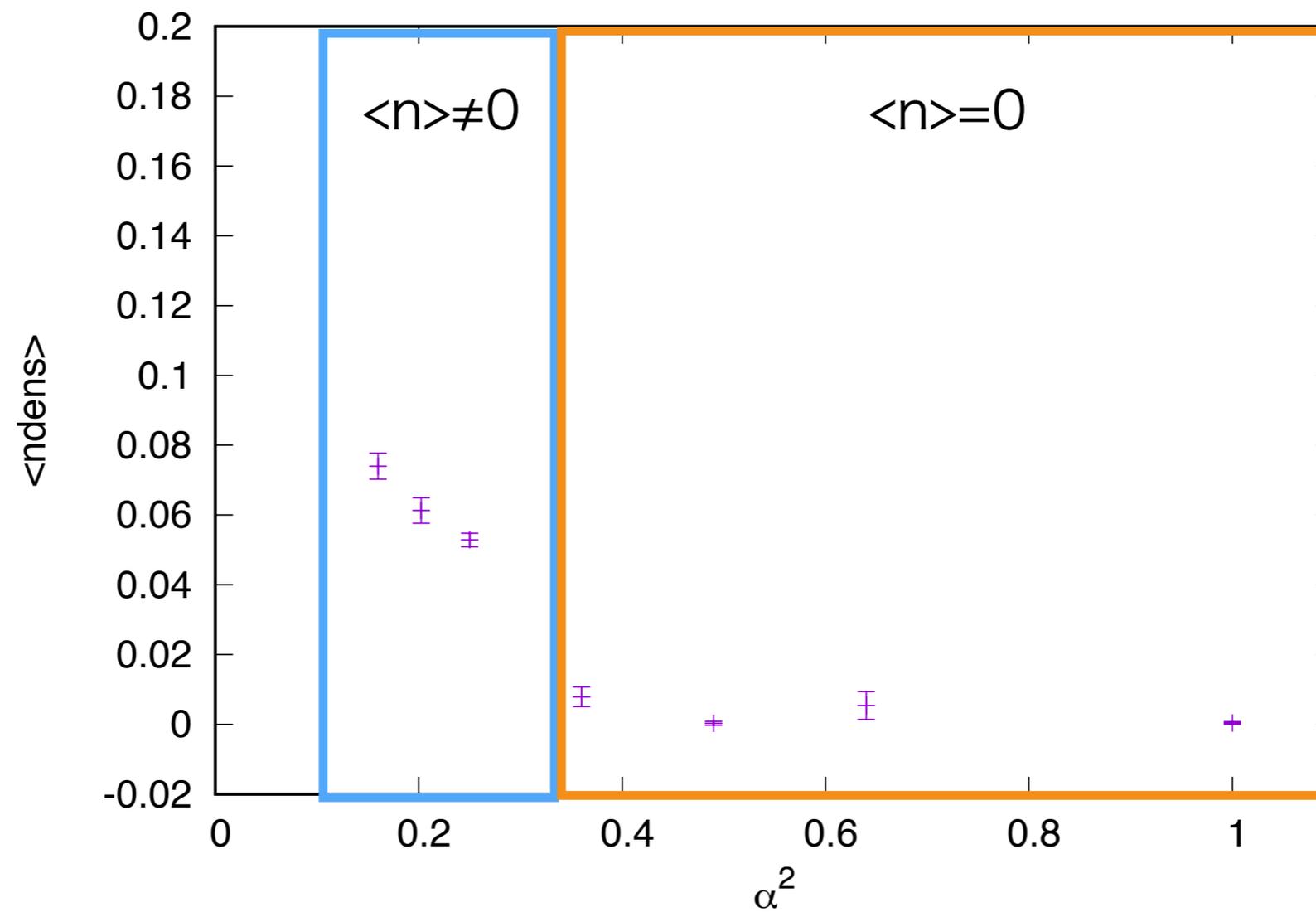
“ $\langle n \rangle$ vs α^2 ” at $\mu = 0.7$

($4^3 \times 8$ lattice)



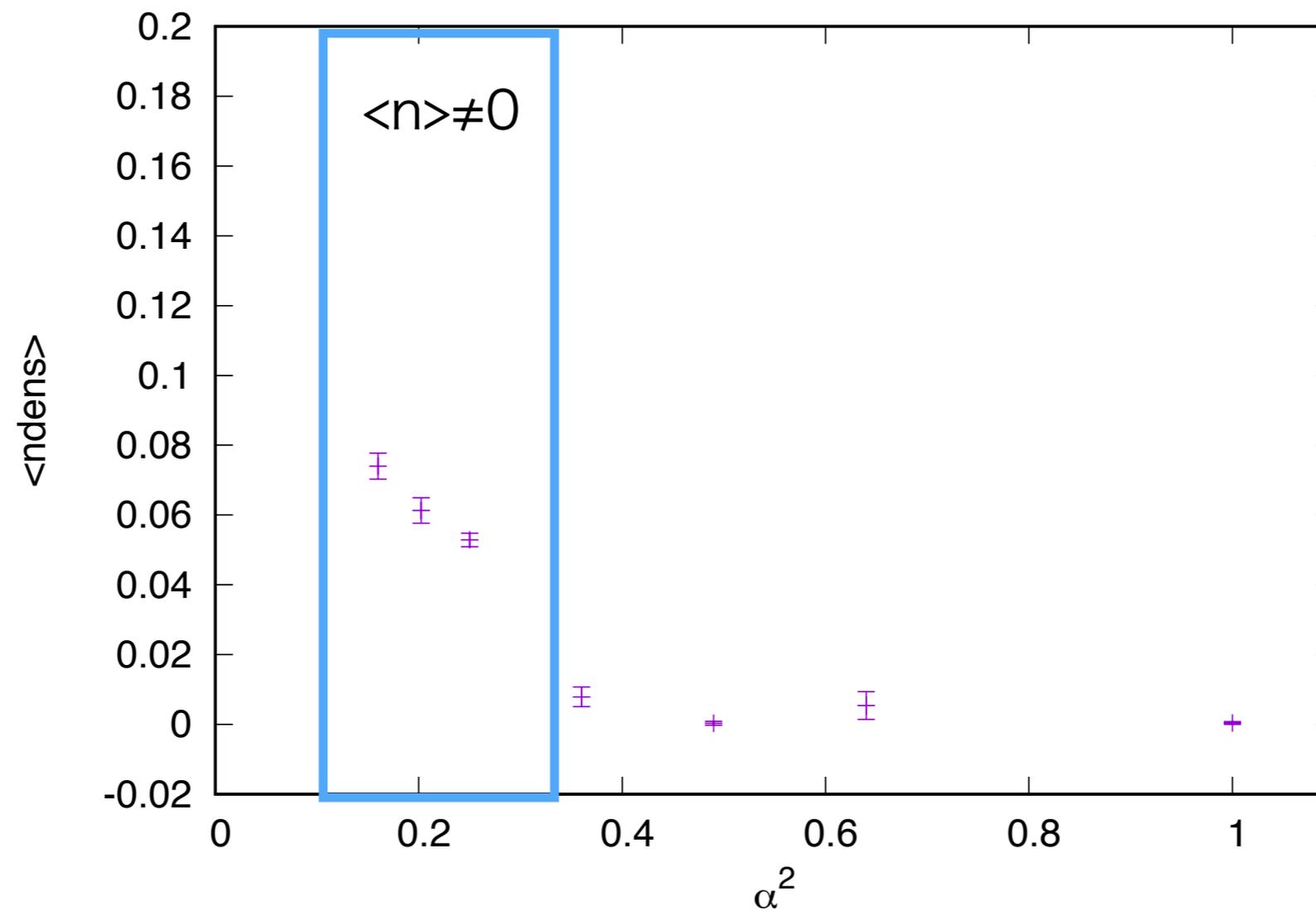
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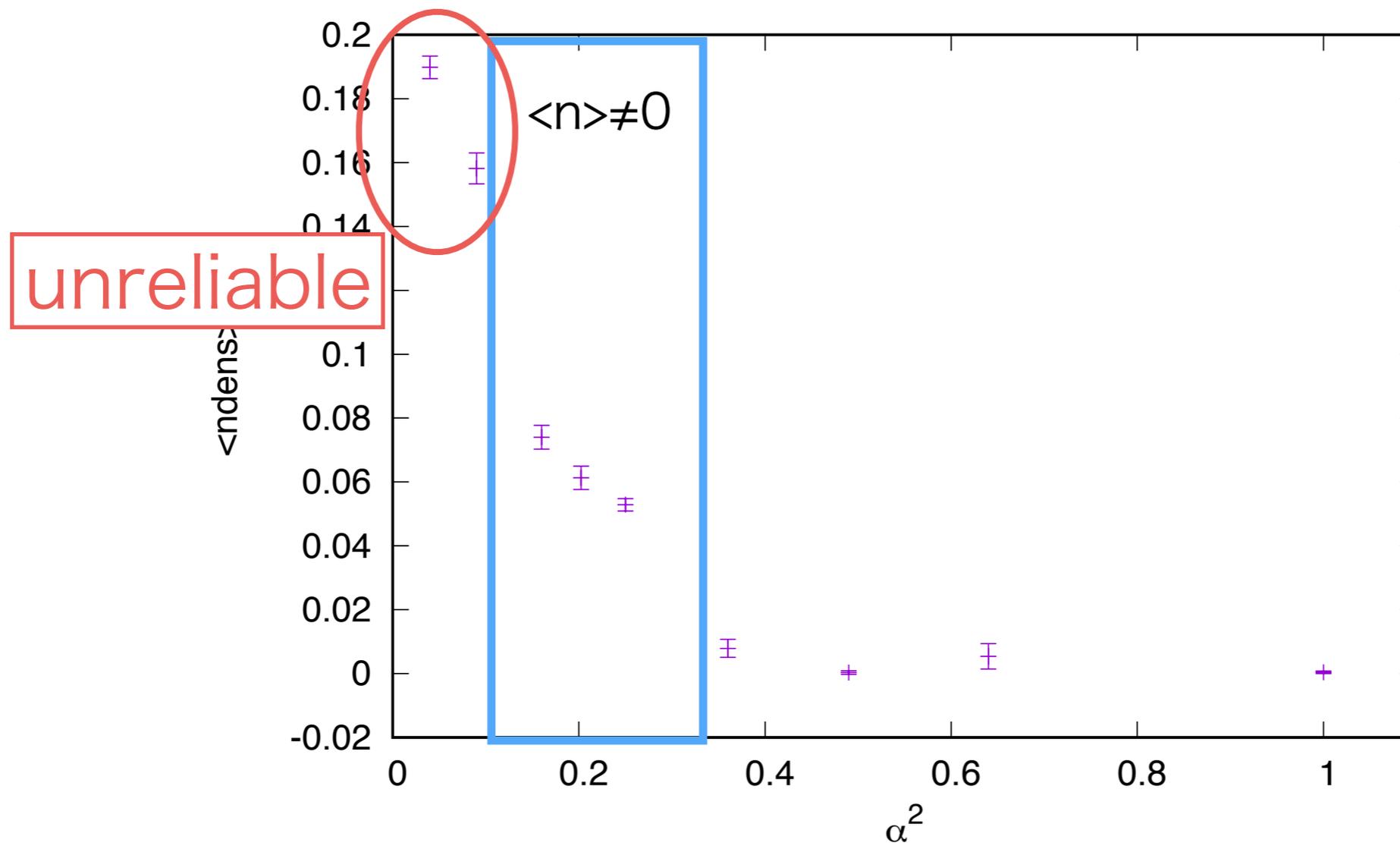
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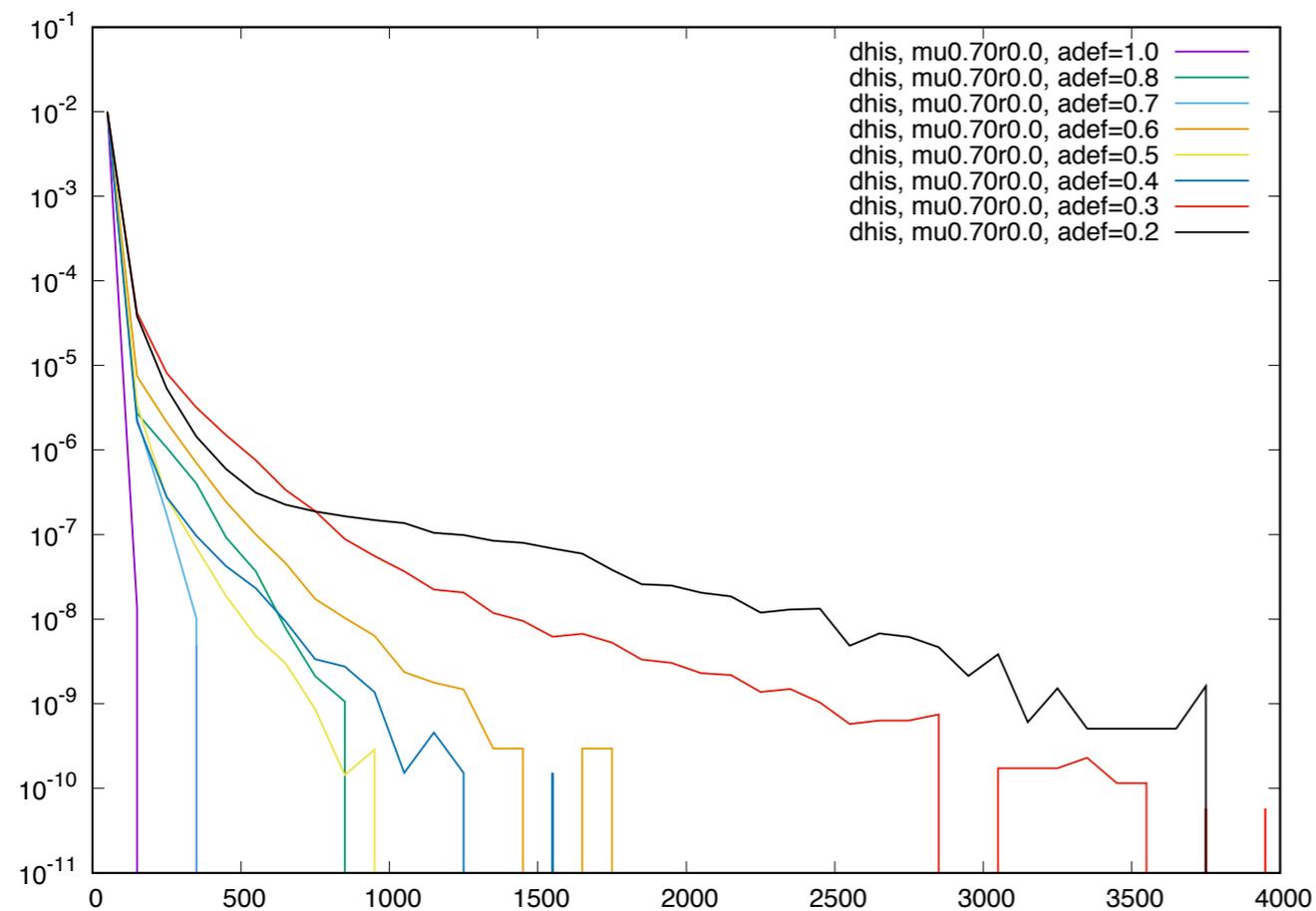


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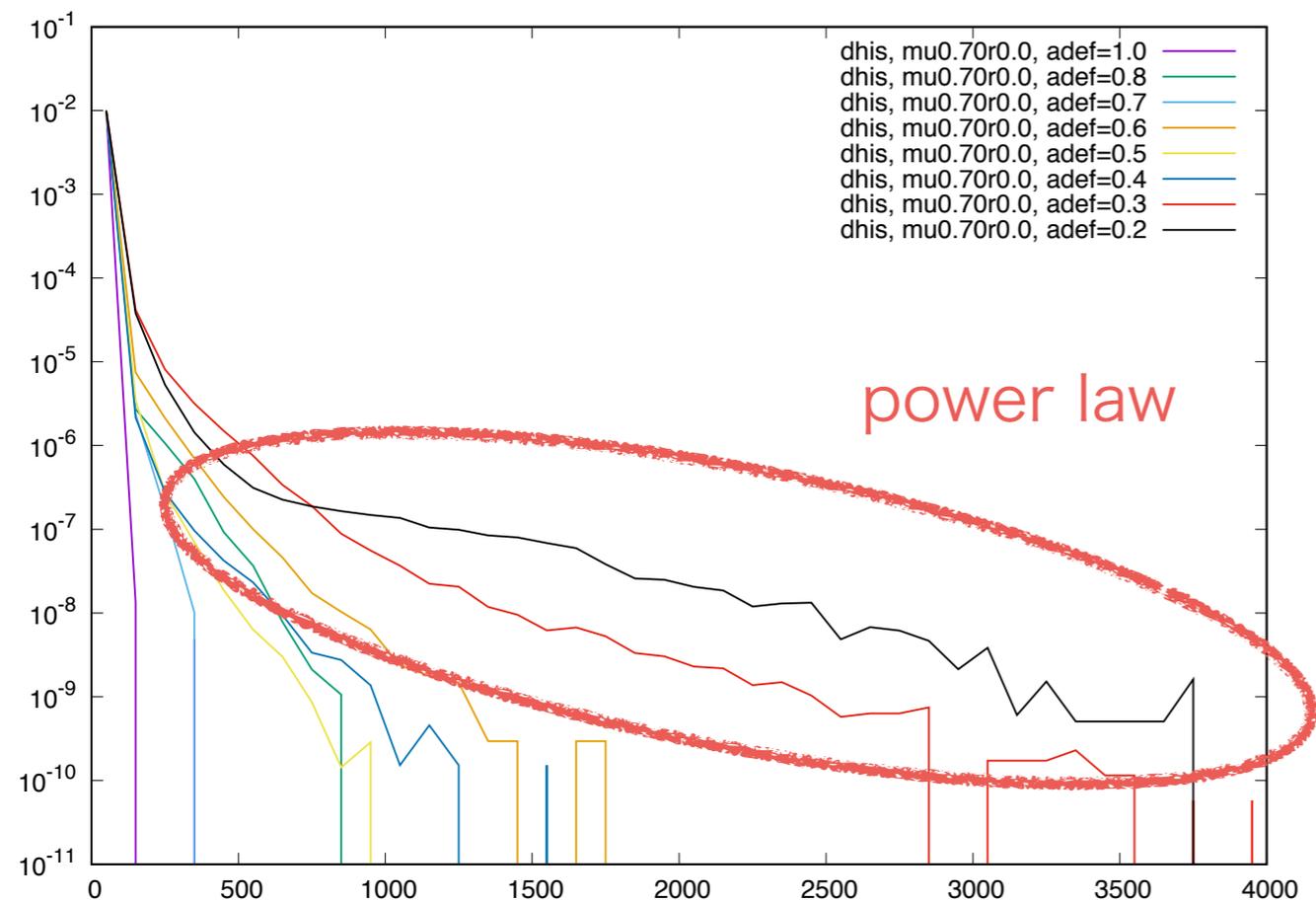
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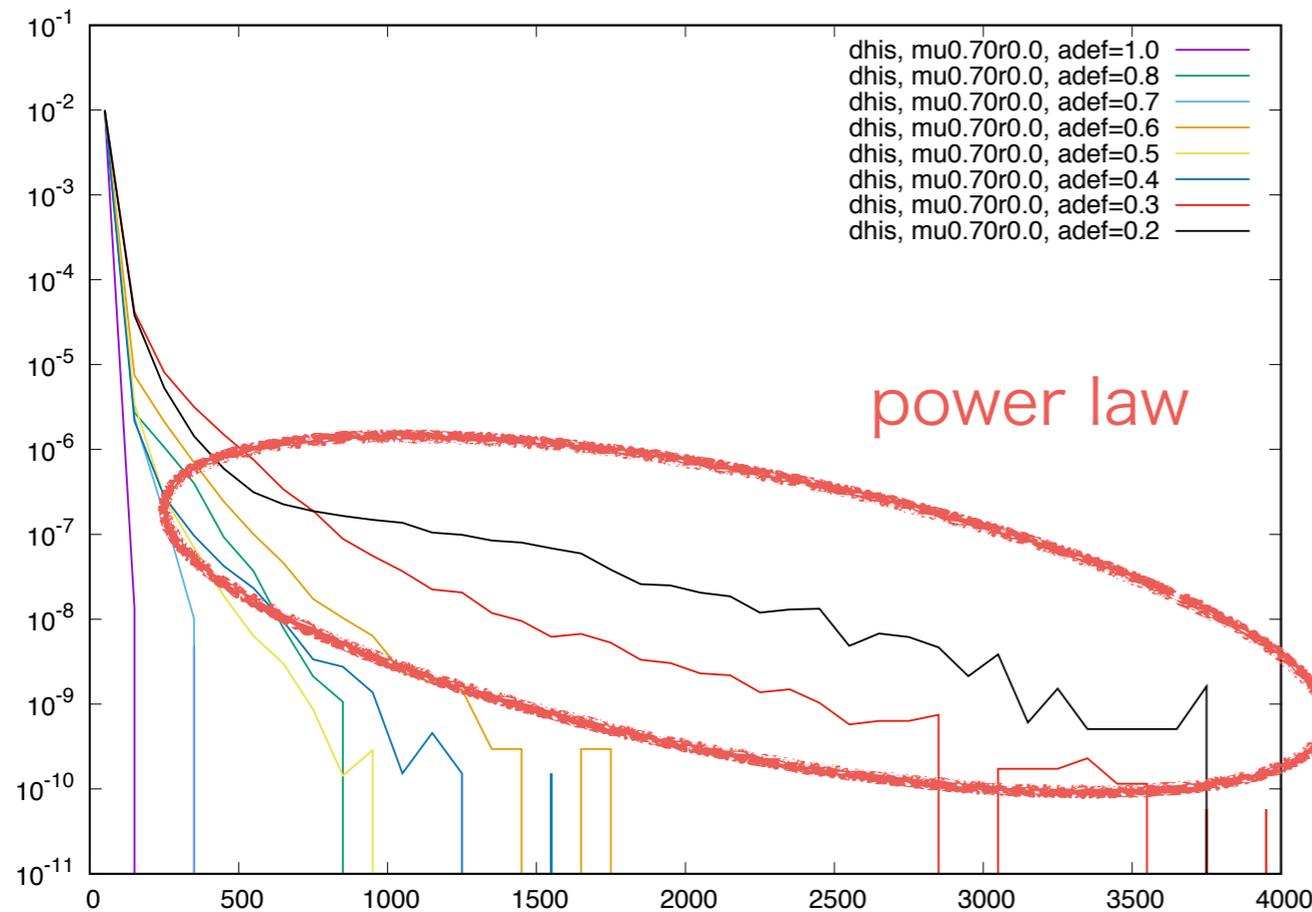
Prob. dist. (semilog) of the magnitude of the drift ($\mu=0.7$)



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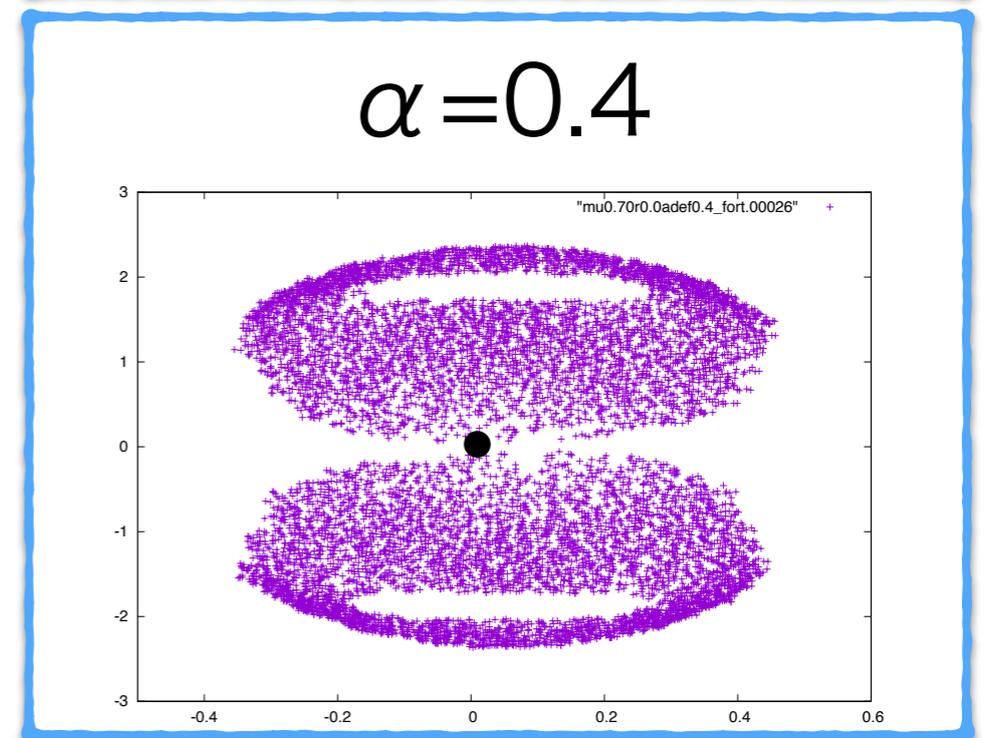
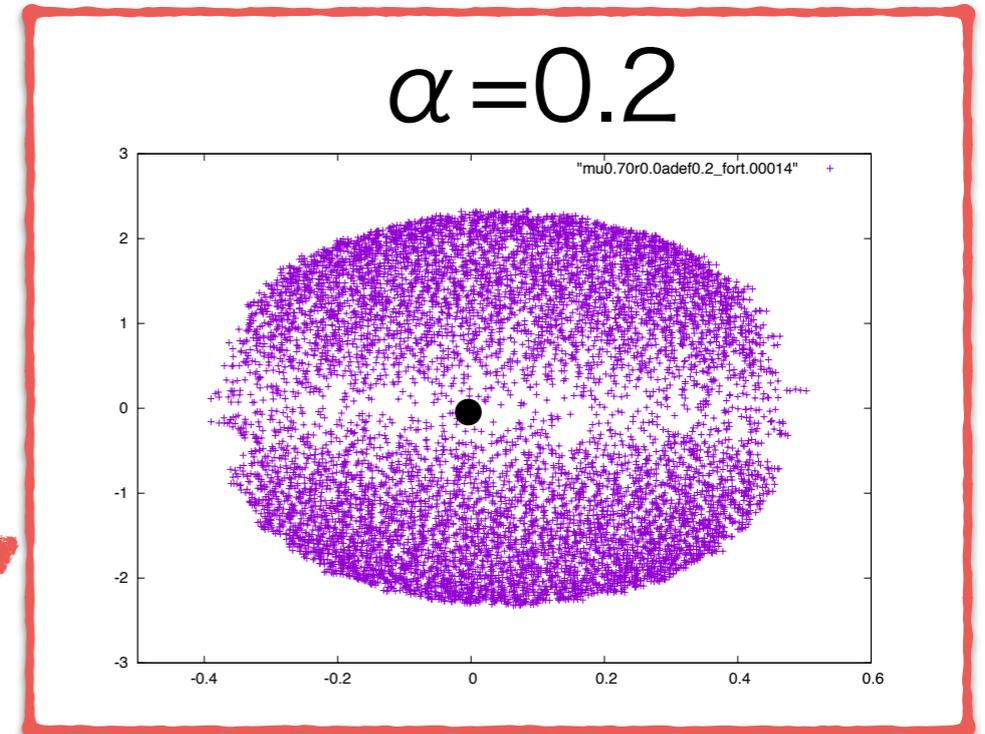
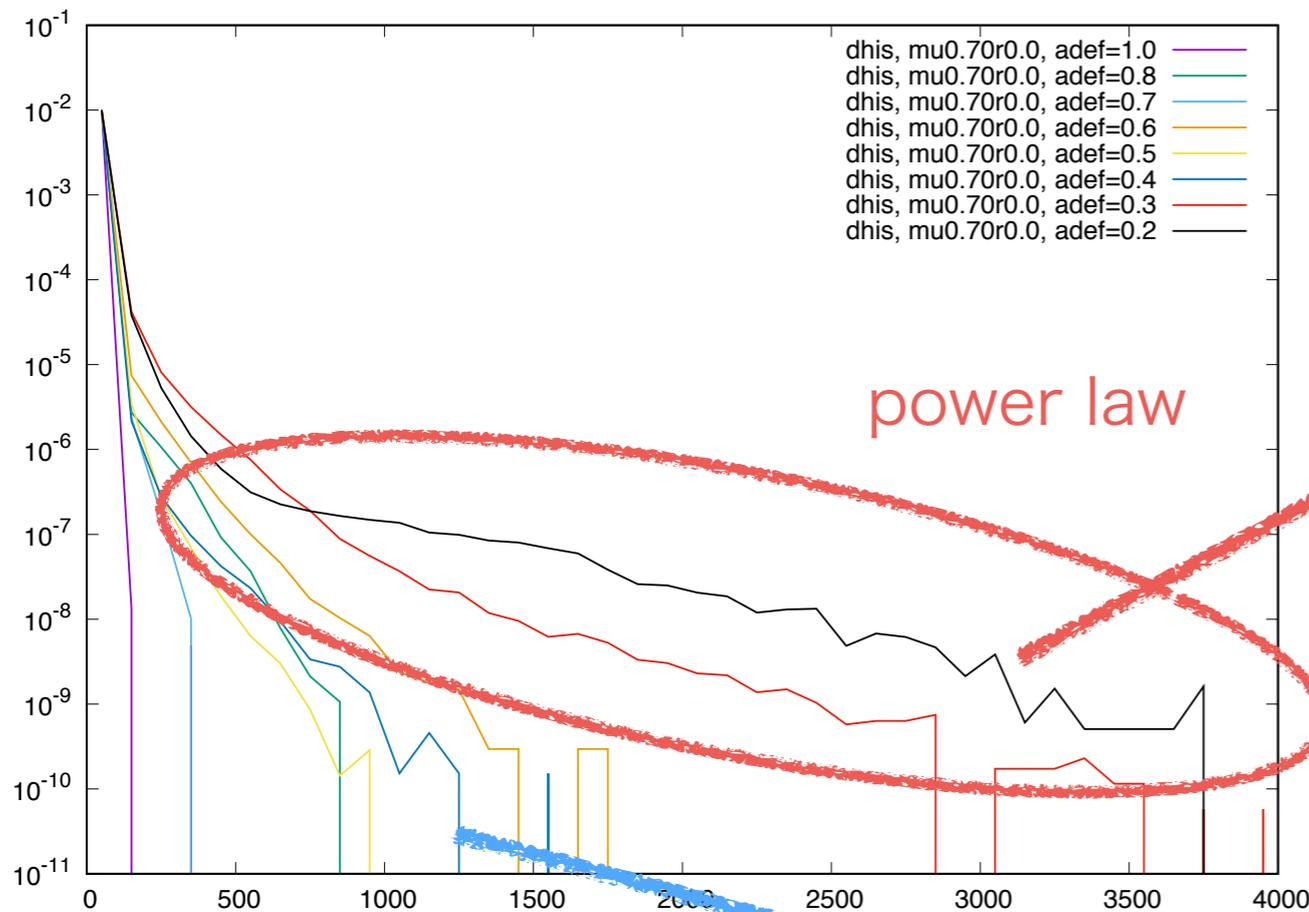
Prob. dist. (semilog) of the magnitude of the drift ($\mu=0.7$)



A drastic change of the dist. can be observed as α is varied.
→ indicates a transition from the exp. fall-off to the power law

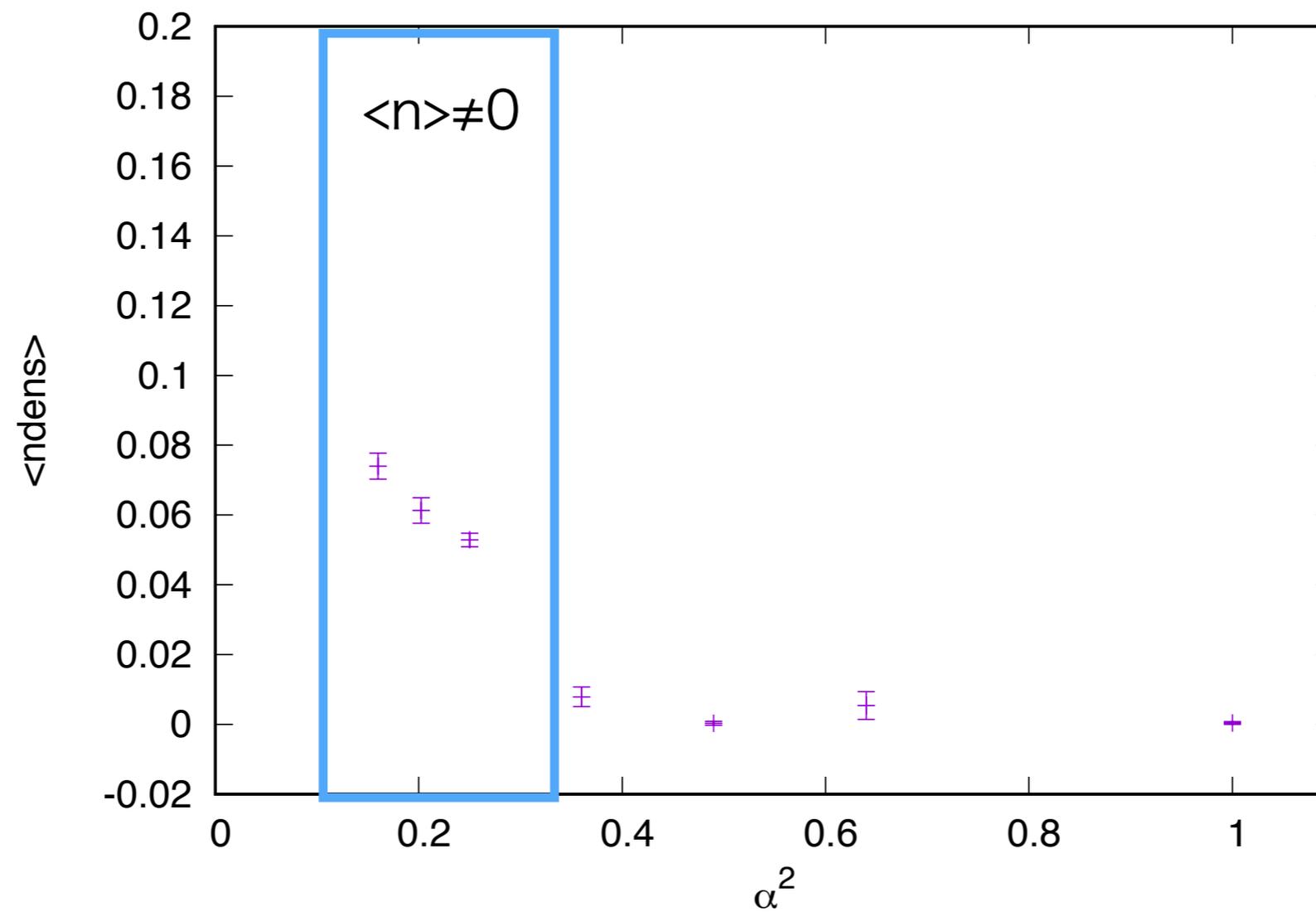
Practically, the fall-off behavior can be judged from its α -dependence

Eigenvalue distribution of $M=D+m$ ($\mu=0.7$)

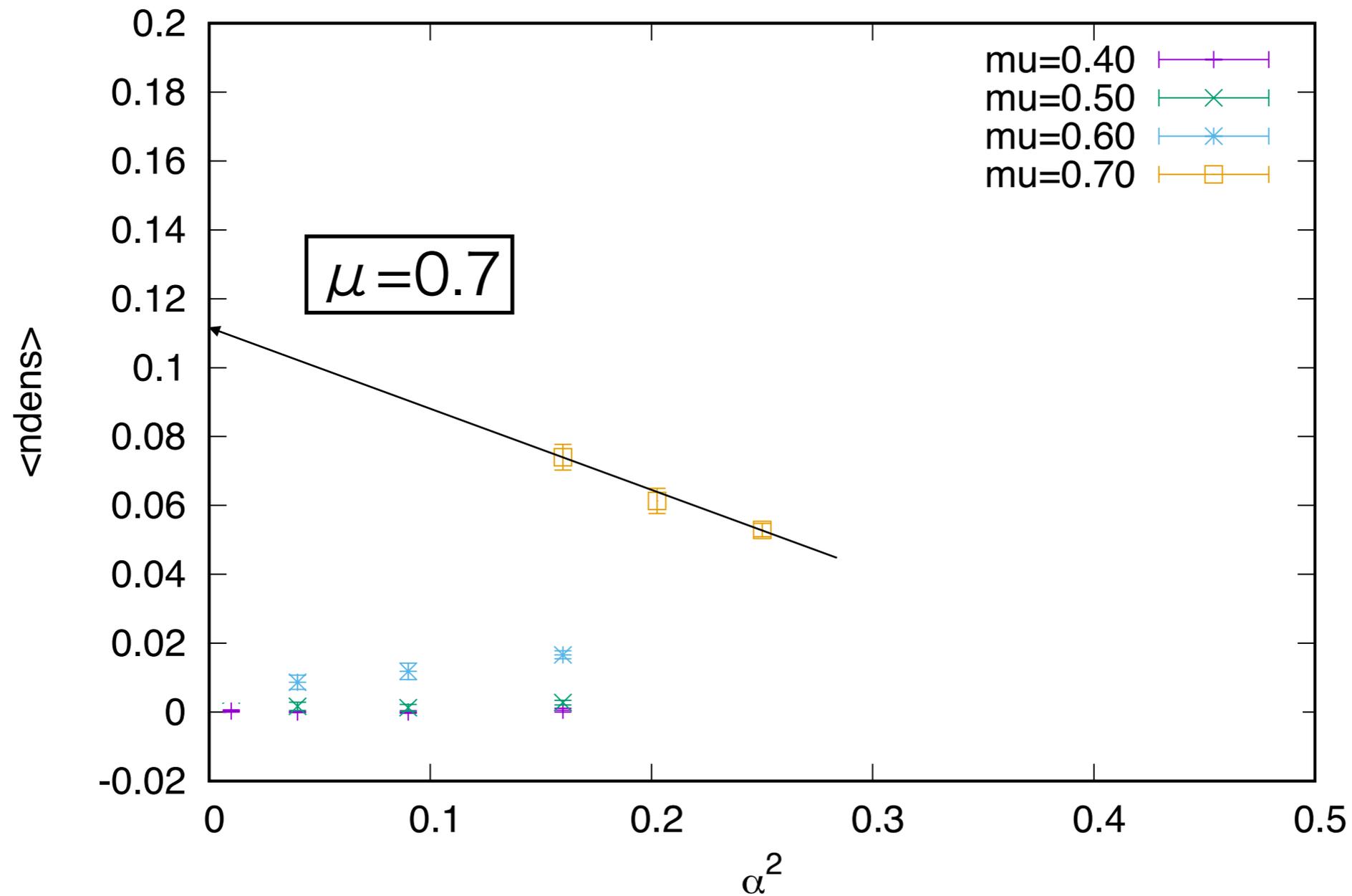


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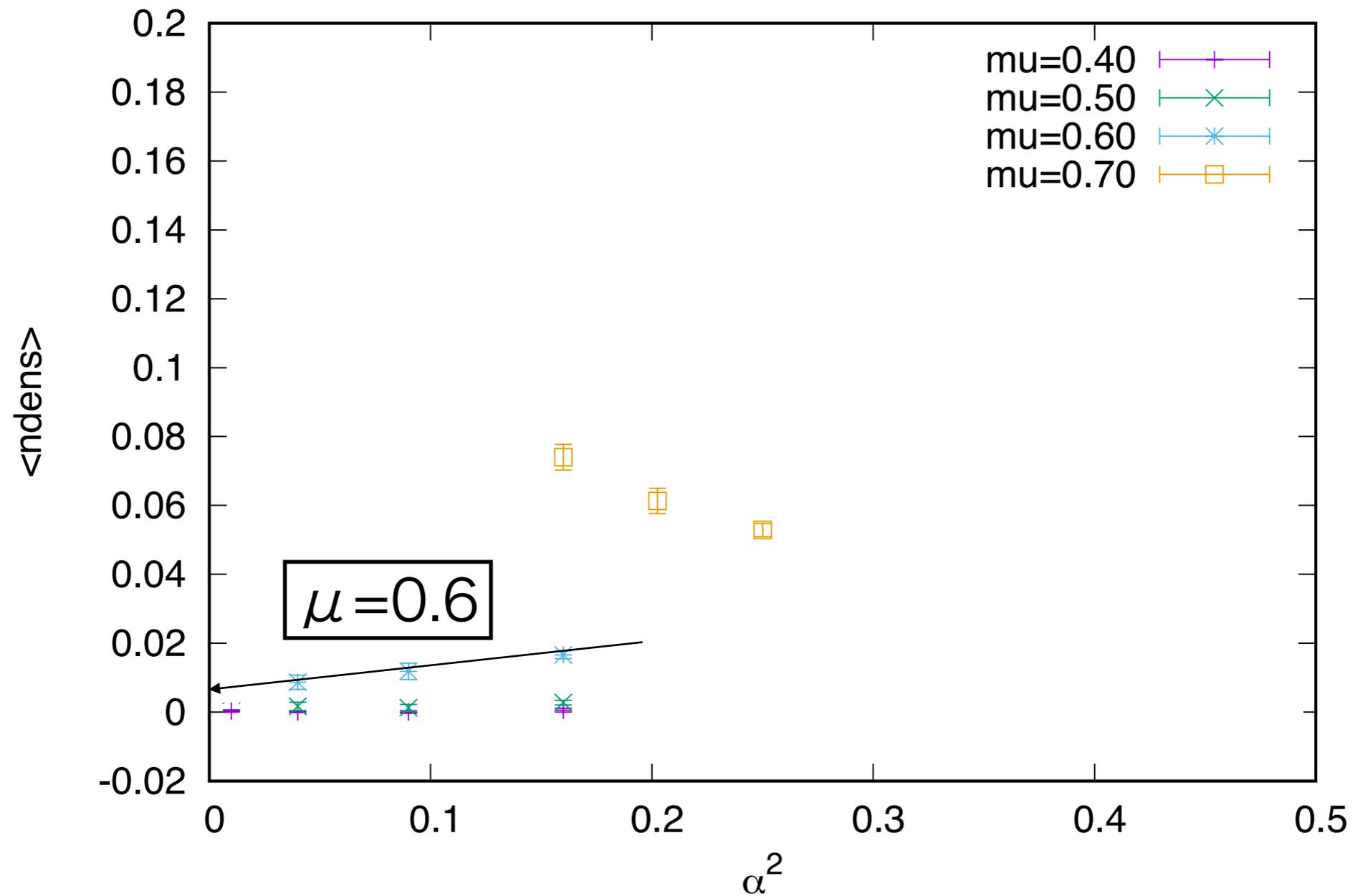
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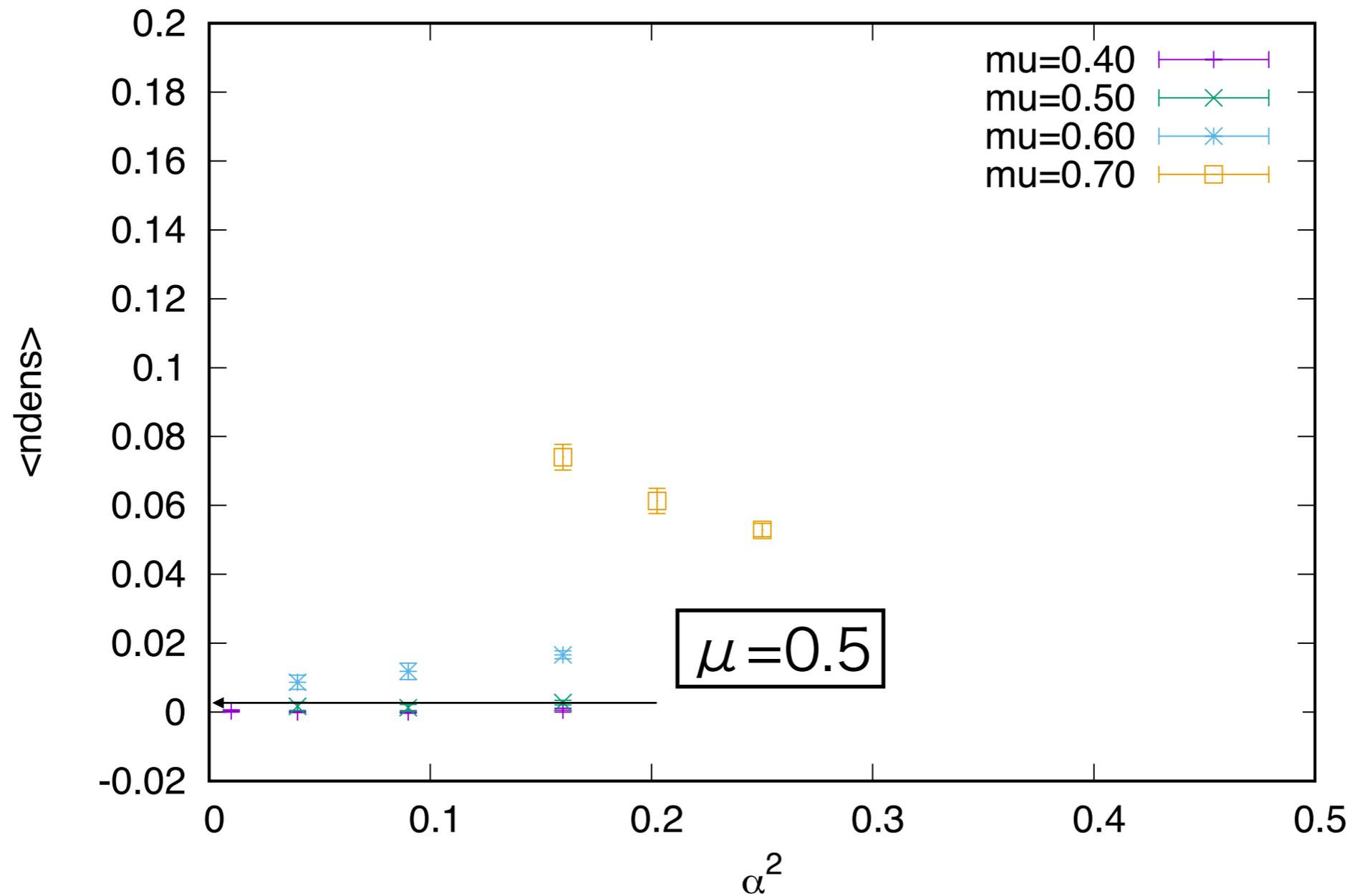
“ $\langle n \rangle$ vs α^2 ” at $\mu=0.4, 0.5, 0.6, 0.7$
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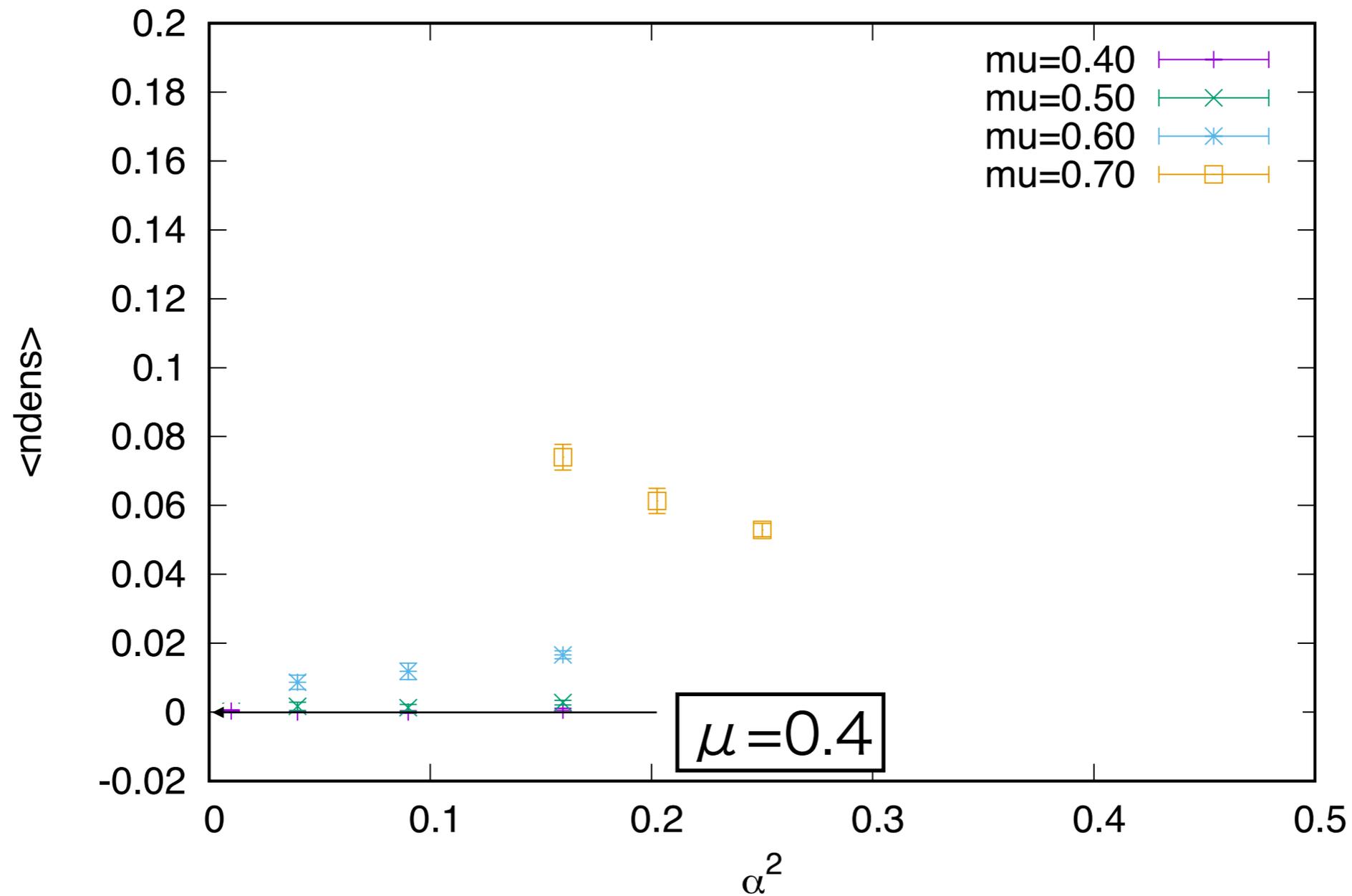
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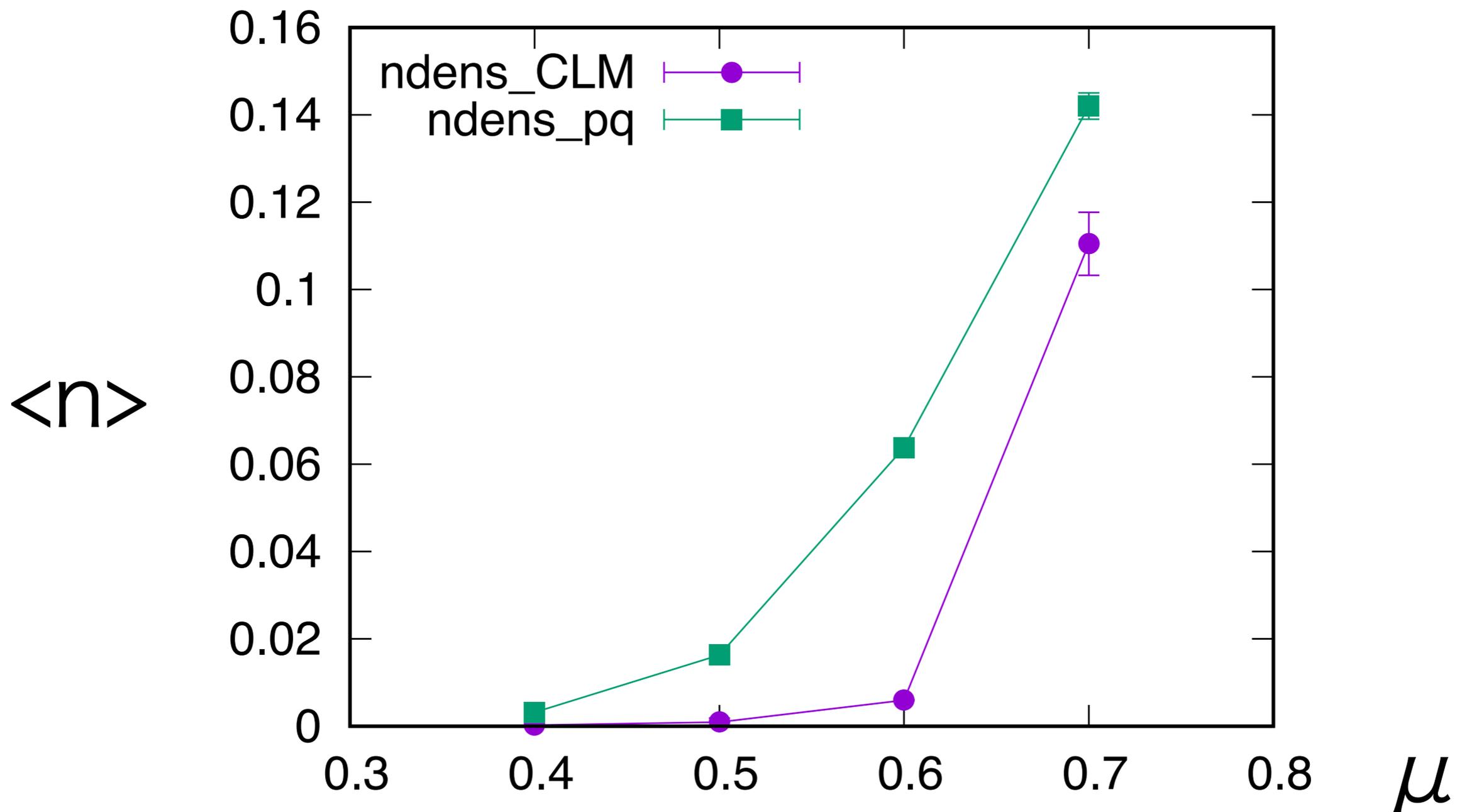
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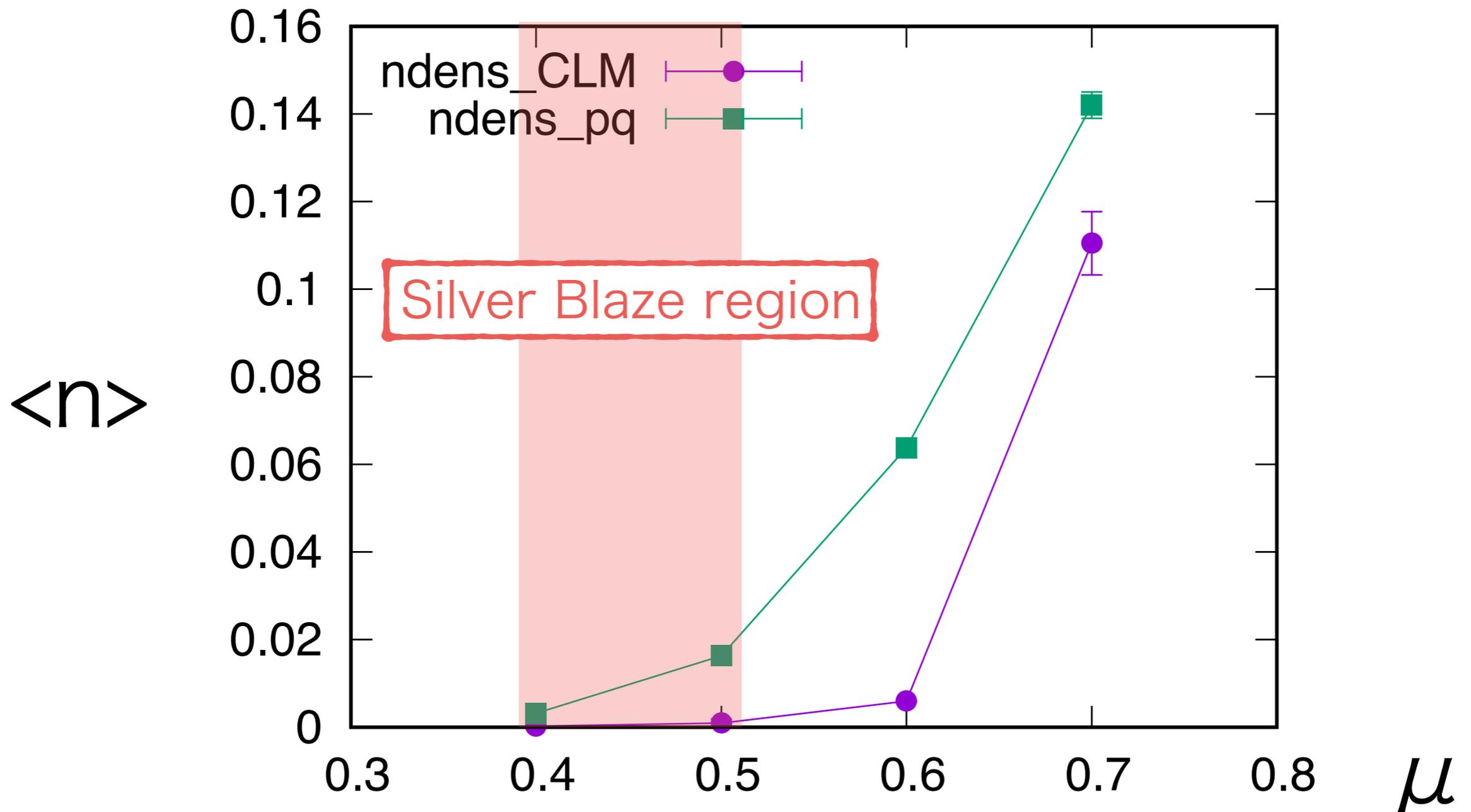
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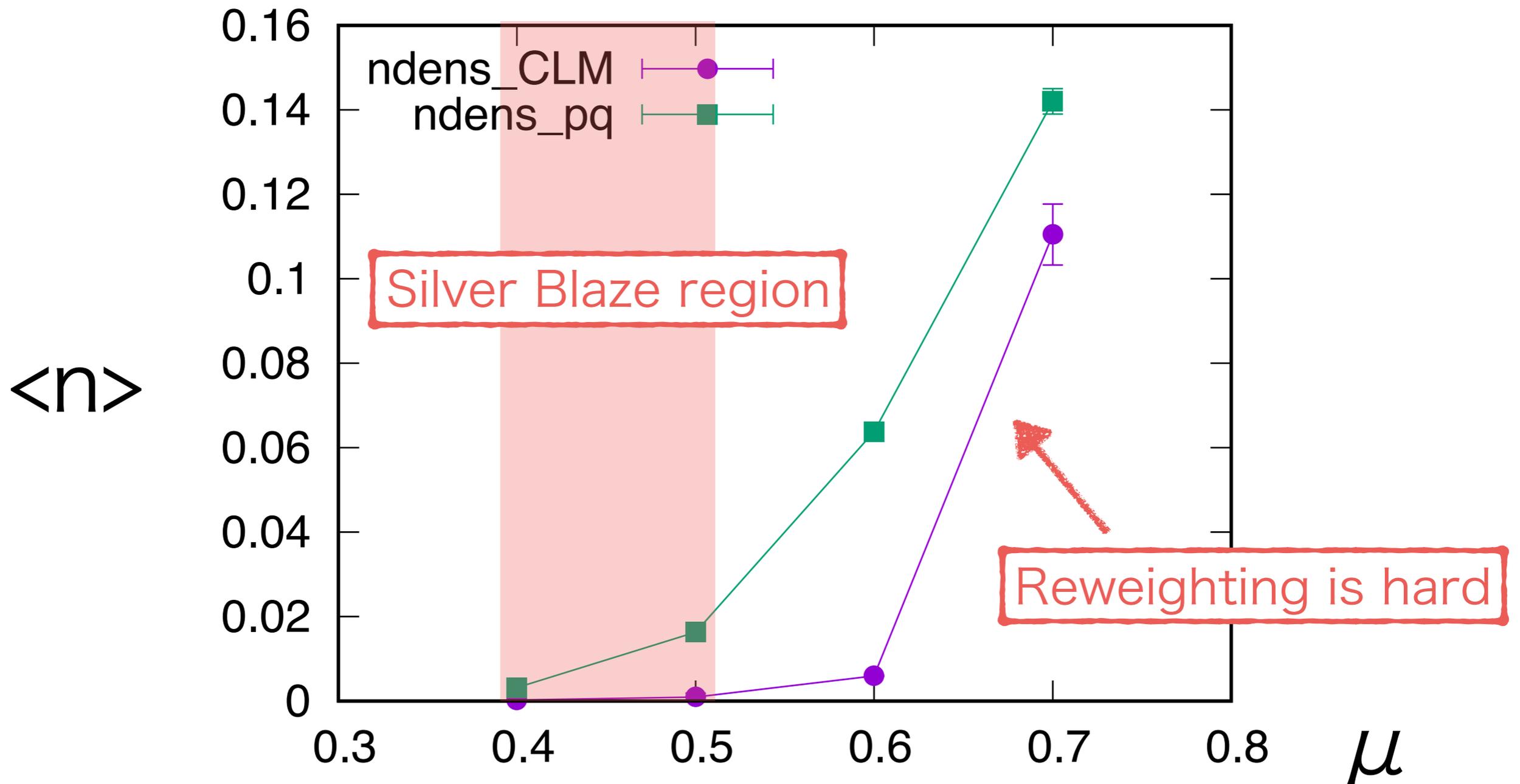
$\langle n \rangle$ vs μ ($4^3 \times 8$)



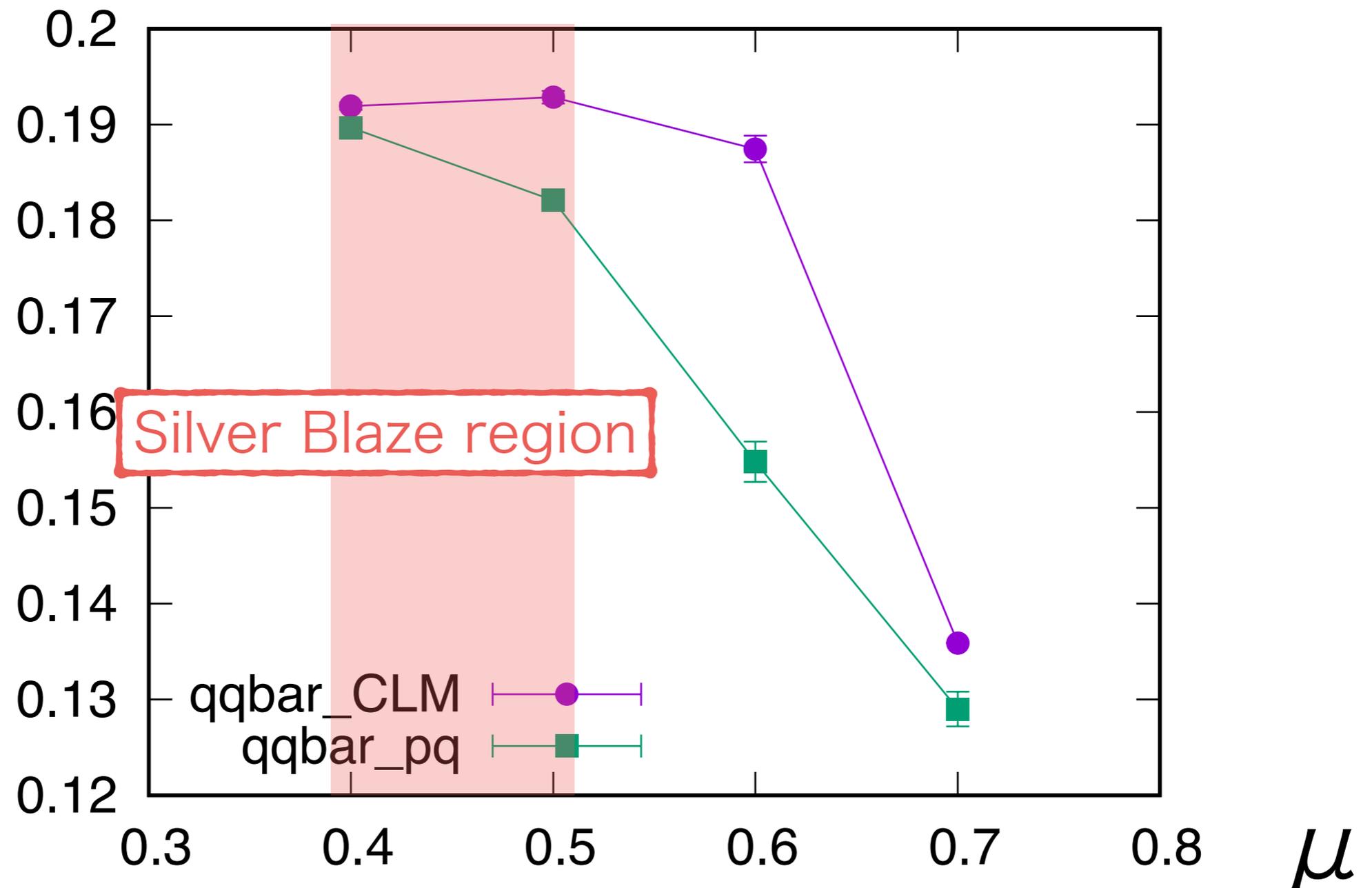
$\langle n \rangle$ vs μ ($4^3 \times 8$)



$\langle n \rangle$ vs μ ($4^3 \times 8$)



$\langle \Sigma \rangle$ vs μ ($4^3 \times 8$)



Preliminary results for

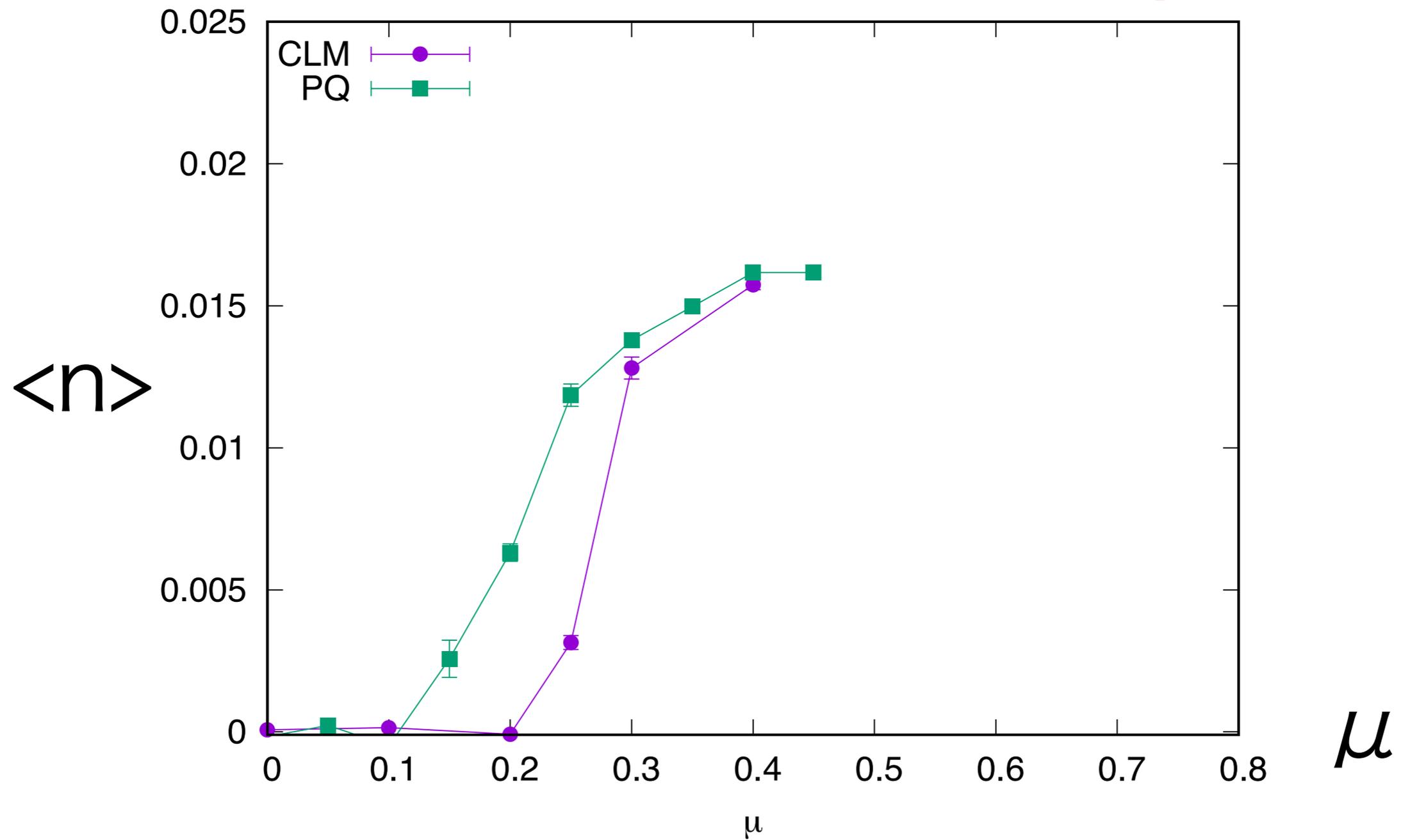
$$N_x=N_y=N_z=8, N_t=16$$

parallel computation employed

[Ito, Matsufuru, Moritake, Nishimura, SS, Tsuchiya, Tsutsui]

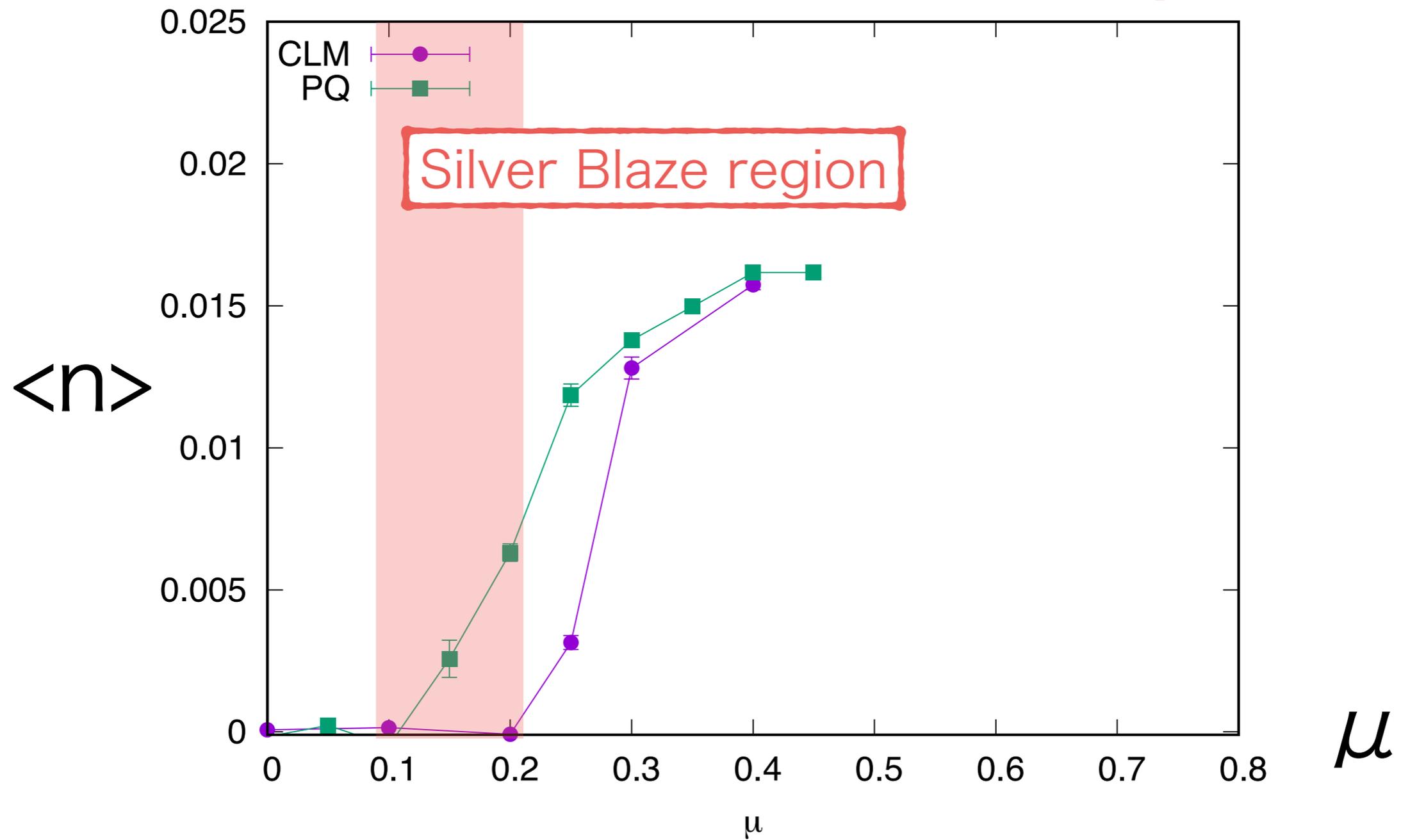
$\langle n \rangle$ vs μ ($8^3 \times 16$)

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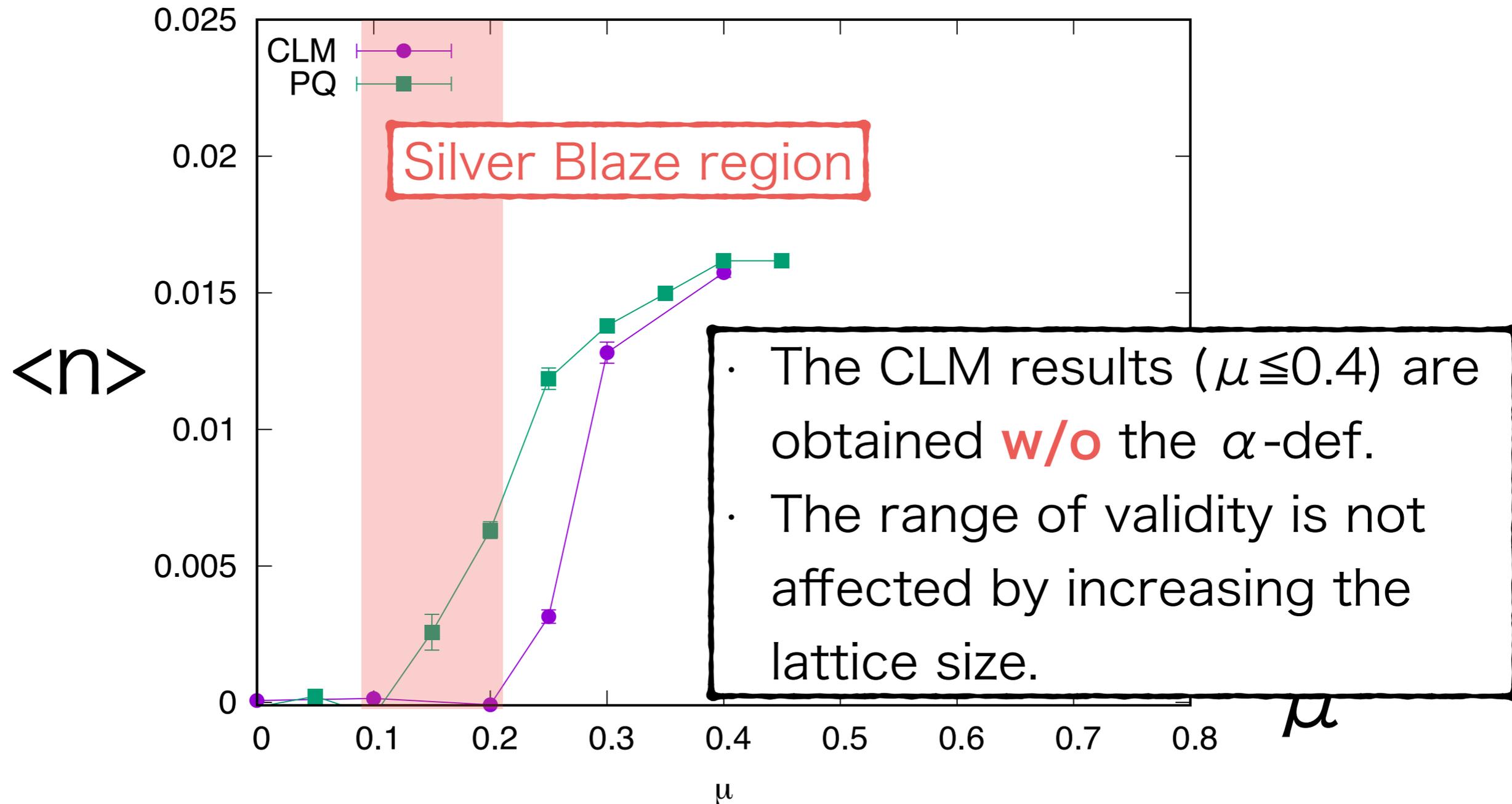
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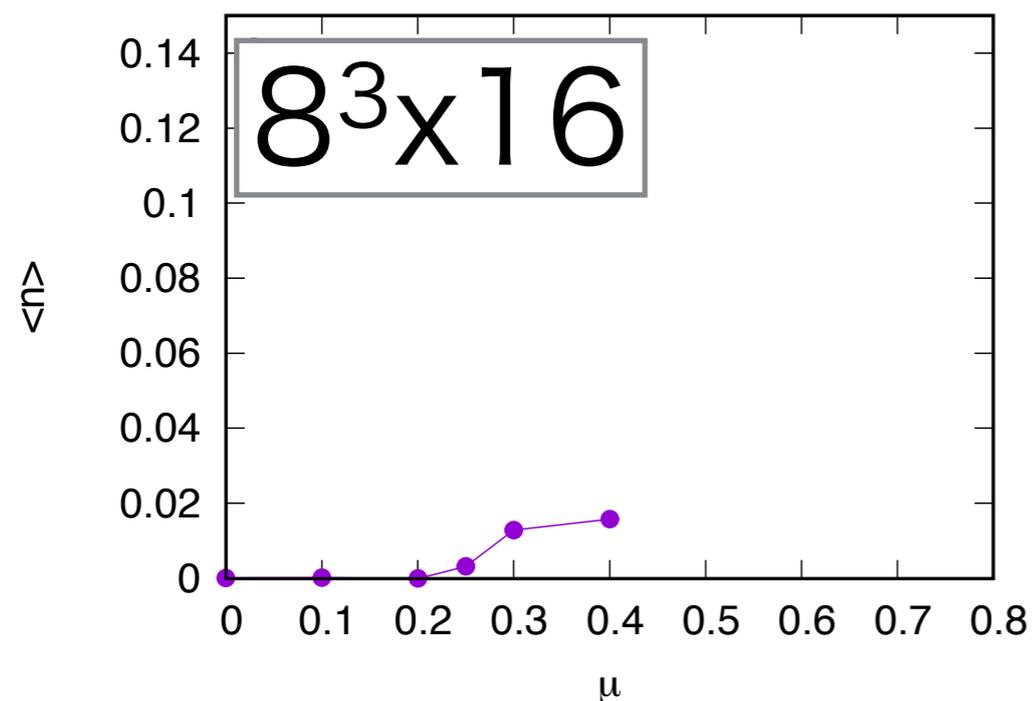
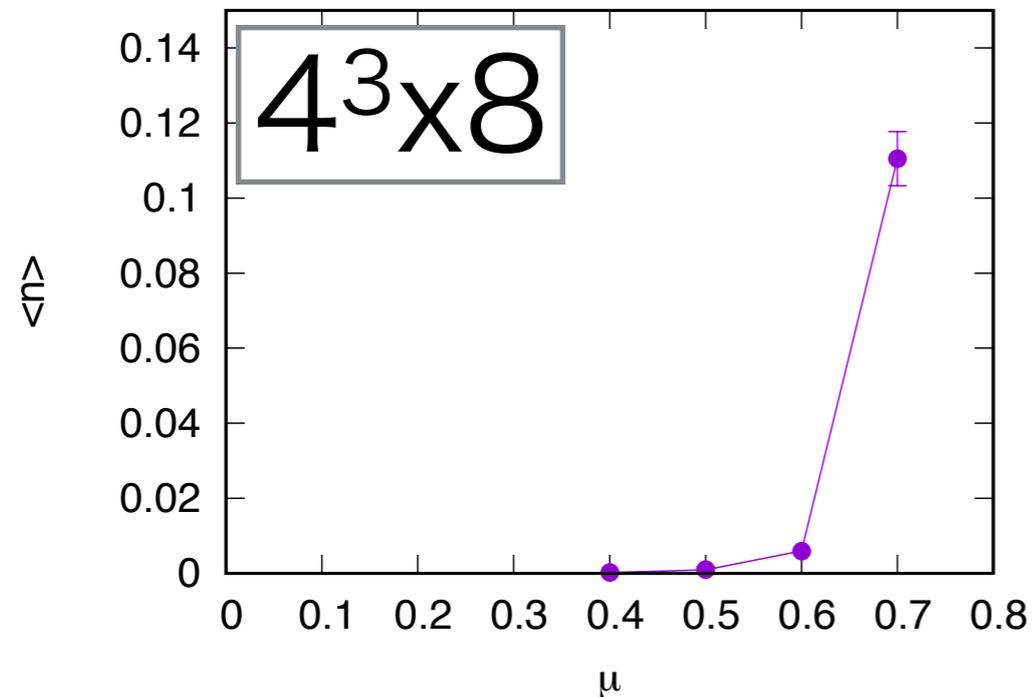
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Comparison of $8^3 \times 16$ with $4^3 \times 8$

preliminary



• $4^3 \times 8 \rightarrow 8^3 \times 16$

$V \rightarrow 8 \times V$

$T \rightarrow T/2$

• The onset of $\langle n \rangle$ for $8^3 \times 16$ is earlier than that for $4^3 \times 8$

← finite volume effect

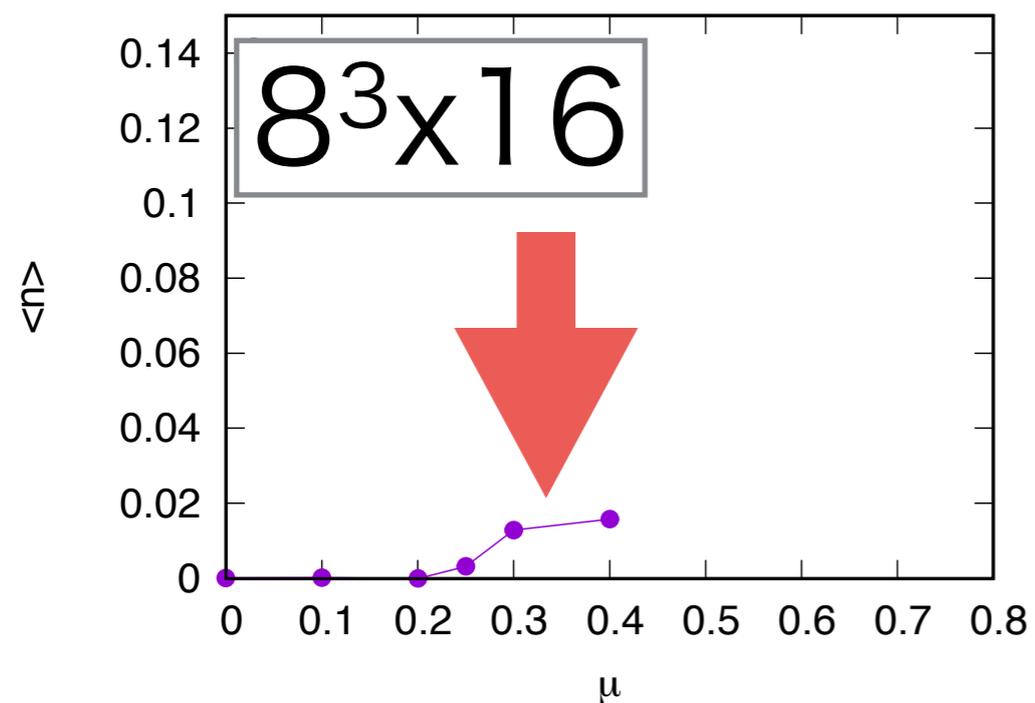
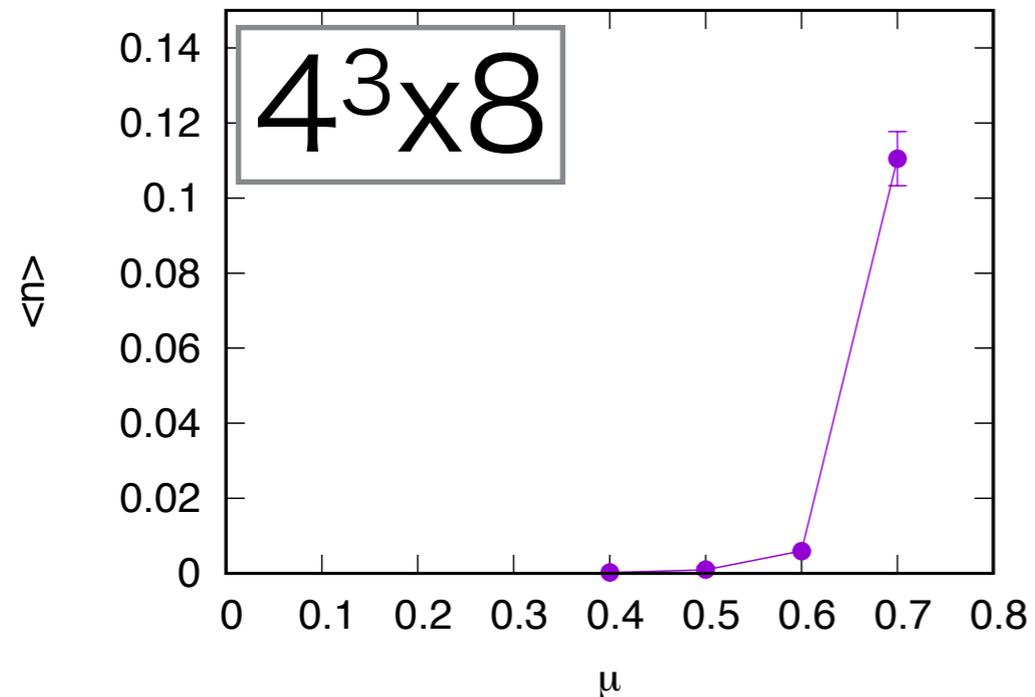
• In $8^3 \times 16$, a saturation-like behavior is observed.

• the nuclear matter phase?

• the second phase transition to quark matter exists?

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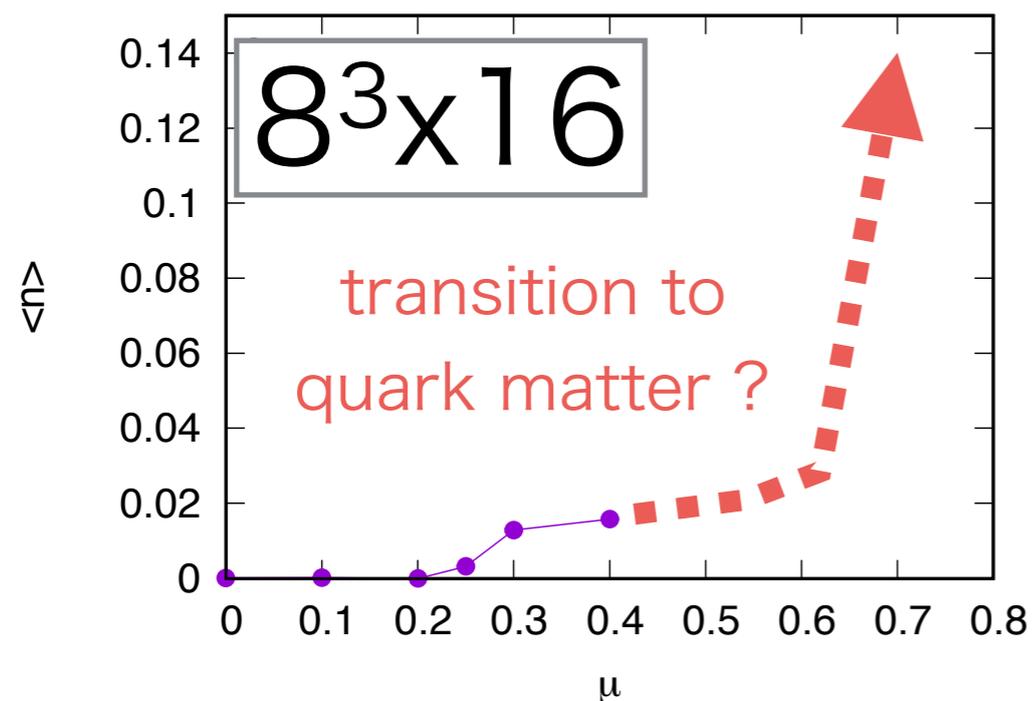
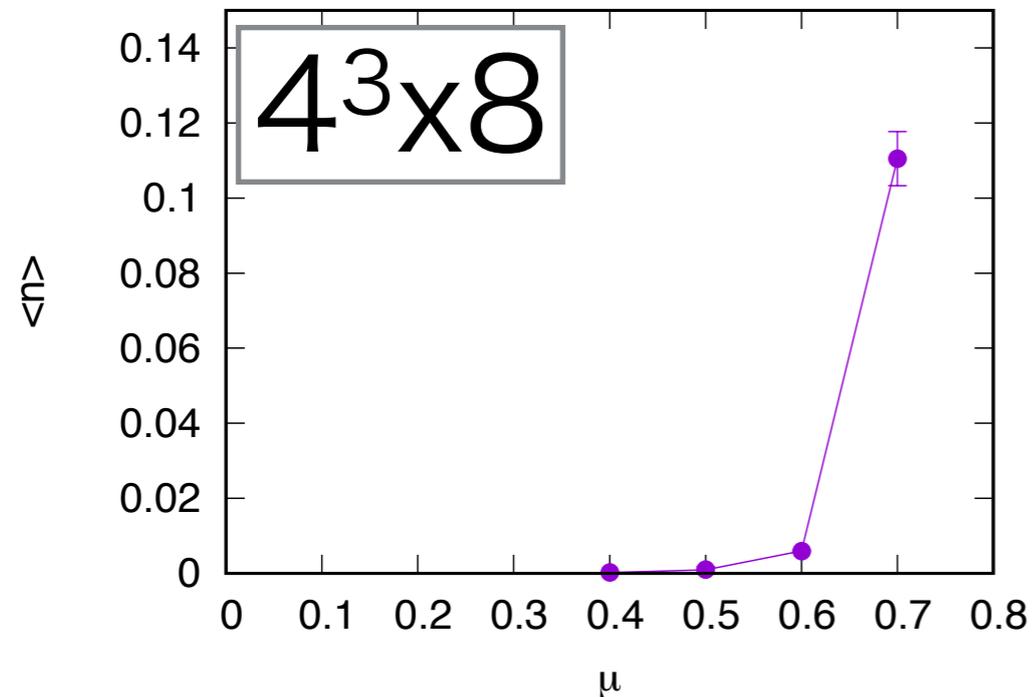
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Summary

- We study QCD at finite density and low temperature by using the CLM with
 - ✓ Gauge cooling
 - ✓ Deformation of the Dirac operator
 - ✓ Reliability check of the obtained results by the prob. dist. of the drift term
- We have found a clear difference between the results obtained by the CLM and the PQ simulation
 - We have succeeded in observing the Silver Blaze phenomenon and the phase transition from $\langle n \rangle = 0$ to $\langle n \rangle \neq 0$.
 - The CLM works well even in the region where the reweighting-based methods are extremely hard.
 - The range of validity is unaffected much by increasing lattice size.

Future work

- Simulation of $8^3 \times 16$ lattice for larger μ by making use of the α -deformation of the Dirac operator.
- Can we observe the second phase transition to quark matter?