

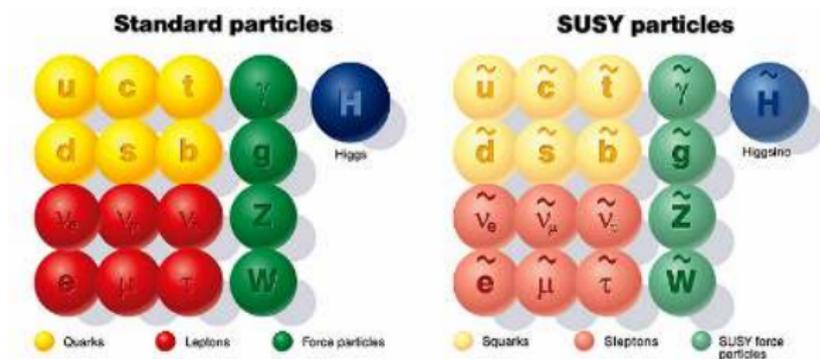
Supermultiplets in $\mathcal{N} = 1$ $SU(2)$ SUSY Yang-Mills Theory

Henning Gerber

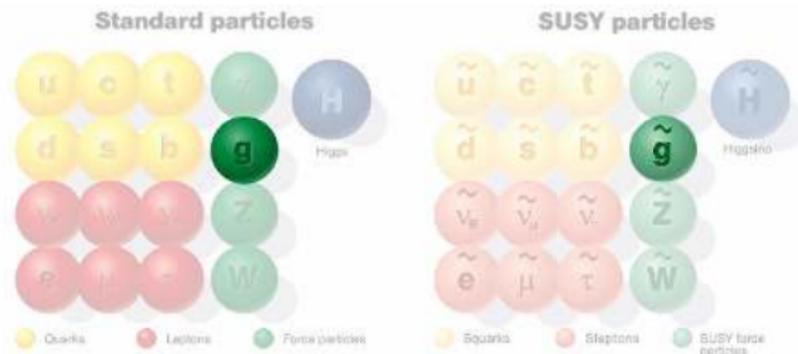
Münster-DESY(-Regensburg-Jena) collaboration

Lattice 2017 - Granada

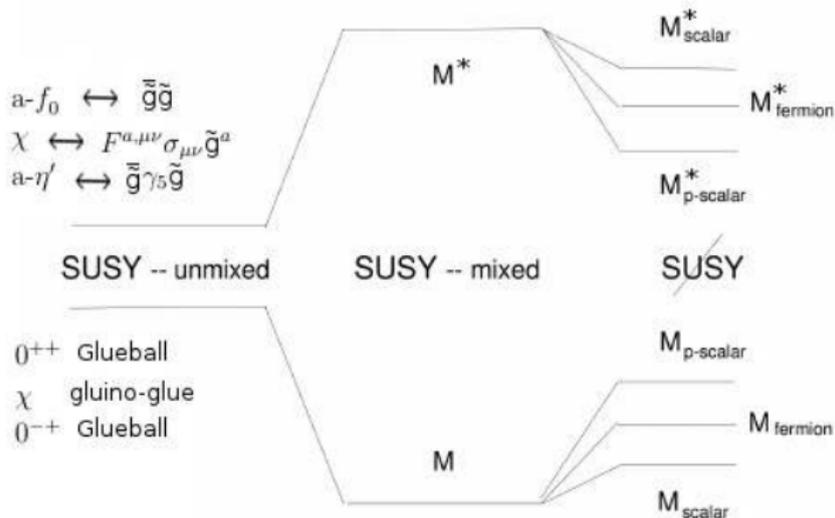




- $\mathcal{L} = -\frac{1}{4}\text{Tr} (F_{\mu\nu}^a F^{\mu\nu,a}) + \frac{i}{2}\tilde{g}\not{D}\tilde{g} + \frac{m_g}{2}\tilde{g}\tilde{g}$
- $A_\mu^a(x)$: gauge fields
- $\tilde{g}(x)$: gluino fields, Majorana fermions in adjoint representation
 - $(D_\mu \tilde{g})^a = \partial_\mu \tilde{g}^a + g f_{abc} A_\mu^b \tilde{g}^c$, here: $f_{abc} = \epsilon_{abc}$
- $\frac{m_g}{2}\tilde{g}\tilde{g}$: soft SUSY-breaking term



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G. R. Farrar, G. Gabadadze, M. Schwetz, "Phys. Rev. D60 (1999) 035002, arXiv:hep-th/9806204.

- Correlation Matrix $C_{ij}(\Delta t) = \langle O_i(\Delta t) O_j(0) \rangle$
- Generalized Eigenvalue Problem (GEVP):

$$C(t) \vec{v}^{(n)} = \lambda^{(n)}(t, t_0) C(t_0) \vec{v}^{(n)}$$
- $\lim_{t \rightarrow \infty} \lambda^{(n)}(t, t_0) \propto e^{-m_n(t-t_0)} (1 + \mathcal{O}(e^{-\Delta m_n(t-t_0)}))$
 - $\Delta m_n = \min_{l \neq n} |m_l - m_n|$

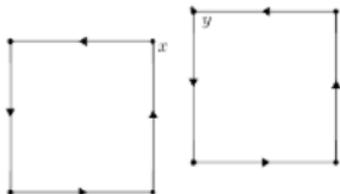
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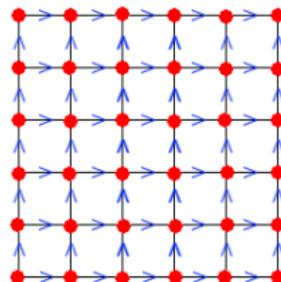
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- Glueballs: $O(x) = \sum_{i < j} P_{ij}(x)$



- Gluino-gluon: $O^a(x) = \sum_{i < j} \sigma_{ij}^{\alpha\beta} \text{Tr}_c [P_{ij}(x) \tilde{g}^\beta(x)]$

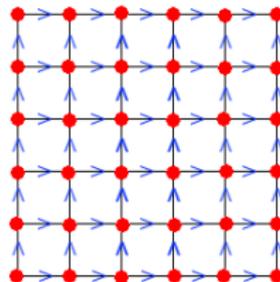
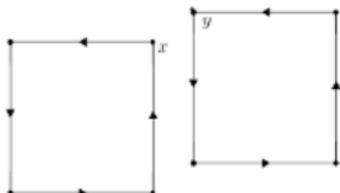
- Mesons: $O(x) = \bar{\tilde{g}}(x) \Gamma \tilde{g}(x)$



● : Fermion fields

→ : Gauge fields

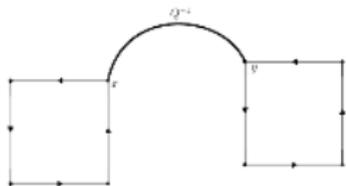
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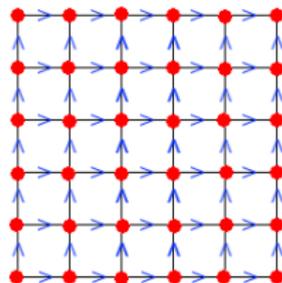
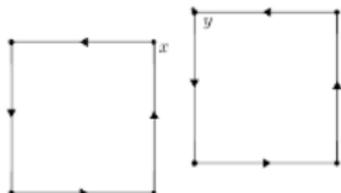
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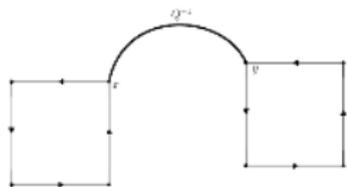
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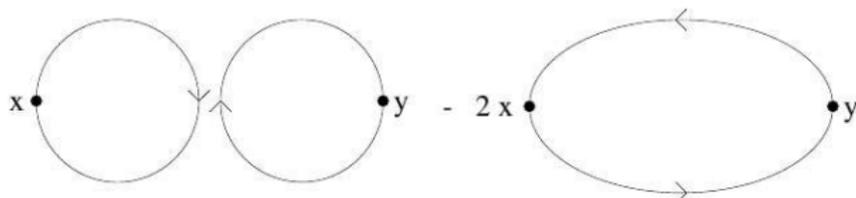
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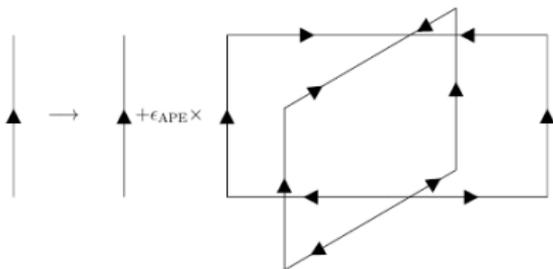
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- Build new operators by applying smearing techniques:
- For gluonic operators: APE Smearing

$$U'_\mu(x) = U_\mu + \epsilon_{\text{APE}} \sum_{\nu=\pm 1, \nu \neq \mu}^{\pm 3} U_\nu^\dagger(x + \mu) U_\mu(x + \hat{\nu}) U_\nu(x)$$

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- For mesonic operators: Jacobi Smearing

$$\tilde{g}'(\vec{x}, t) = \sum_{\vec{y}} F(\vec{x}, \vec{y}) \tilde{g}(\vec{y}, t)$$

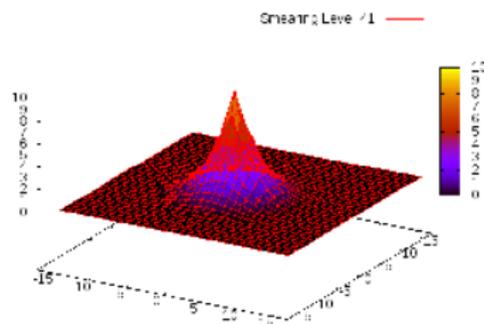
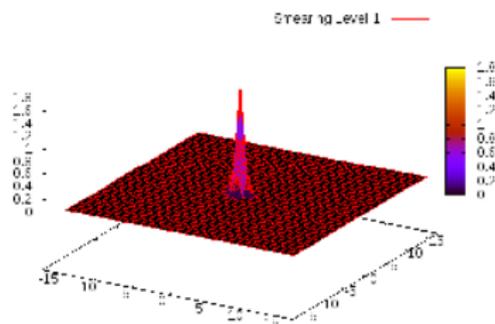
- $F(\vec{x}, \vec{y})$ = iterative solution of 3d Klein-Gordon equation for source and sinks :

$$F_{\beta b, \alpha a}(\vec{x}, \vec{y}) = \delta_{\vec{x}, \vec{y}} \delta_{\beta \alpha} + \delta_{\beta \alpha} \sum_{i=1}^{N_J} \left(\kappa_J \sum_{\mu=1}^3 \left[\delta_{\vec{y}, \vec{x} + \hat{\mu}} U_{\mu, ba}(x) + \delta_{\vec{y} + \hat{\mu}, \vec{x}} U_{\mu, ba}^\dagger(x) \right] \right)^i$$

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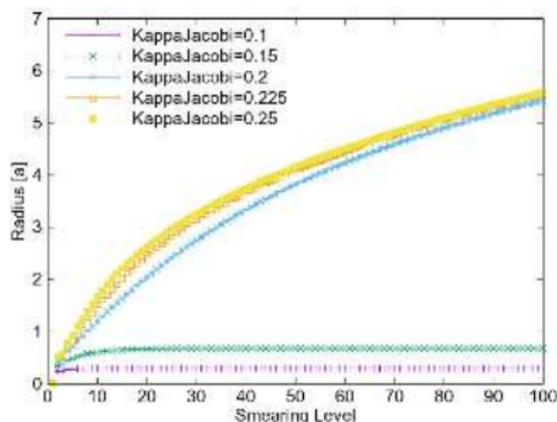


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Jacobi Smearing

- Smearing Radius:

$$R_J = \frac{\sum_{\vec{x}} |\vec{x}|^2 |F(\vec{x}, 0)|^2}{\sum_{\vec{x}} |F(\vec{x}, 0)|^2}$$



- ⇒ Use $\kappa_{\text{Jacobi}} = 0.2$
- ⇒ Use smearing levels up to 80
- ⇒ Optimizing the signal lead to choosing smearing levels 0, 40 and 80

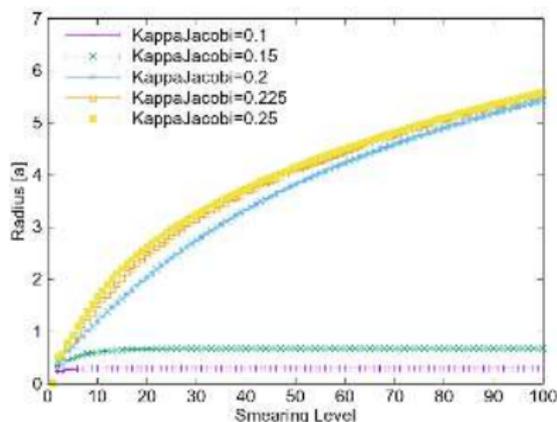
APE smearing:

- similar analysis as for Jacobi smearing
- $\epsilon_{\text{APE}} \sim 0.4$
- smearing levels up to 95 for $32^3 \times 64$ lattice
- i.e. $\{5, 15, \dots, 95\}$ for gluino-gluon

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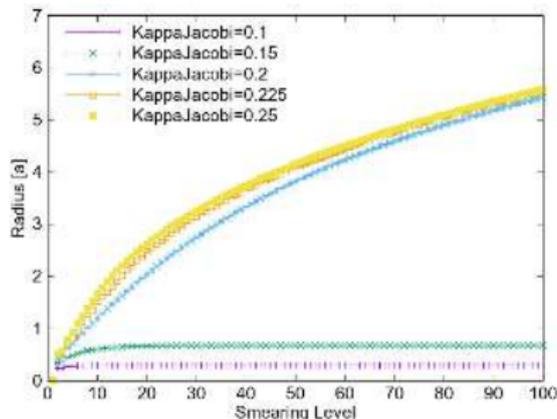
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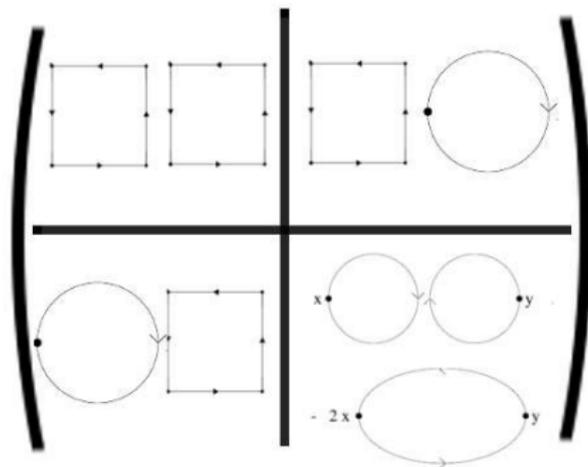
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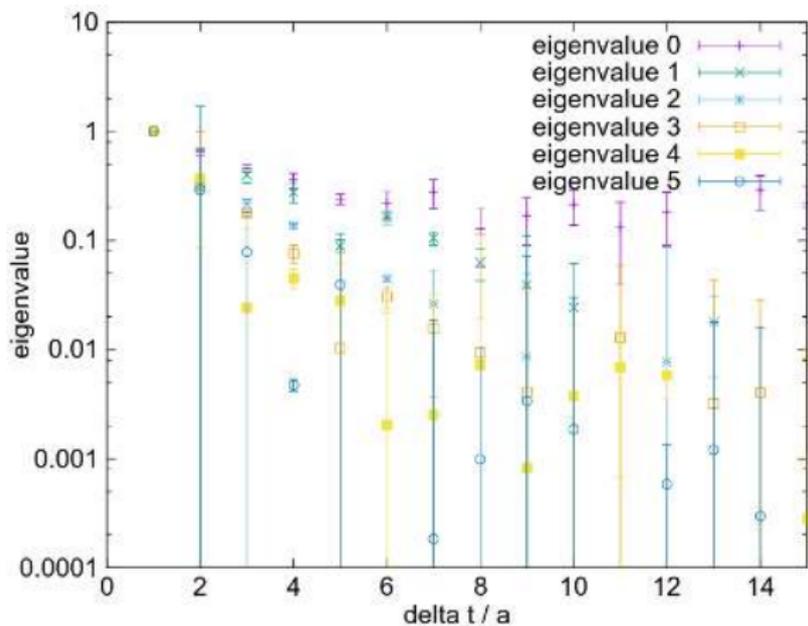
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- Mixing of meson and glueball states allows using a larger operator basis for determining the masses in the 0^{++} and 0^{-+} channel
 - e.g. $a-f_0$ and 0^{++} -glueball
- Build correlation matrix C from mesonic and gluonic operators:

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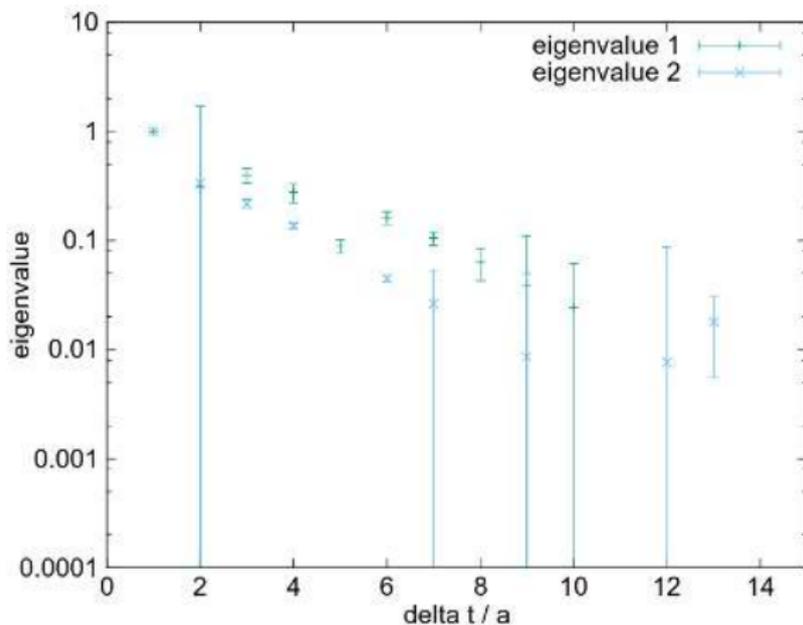


Ordering the eigenvalues



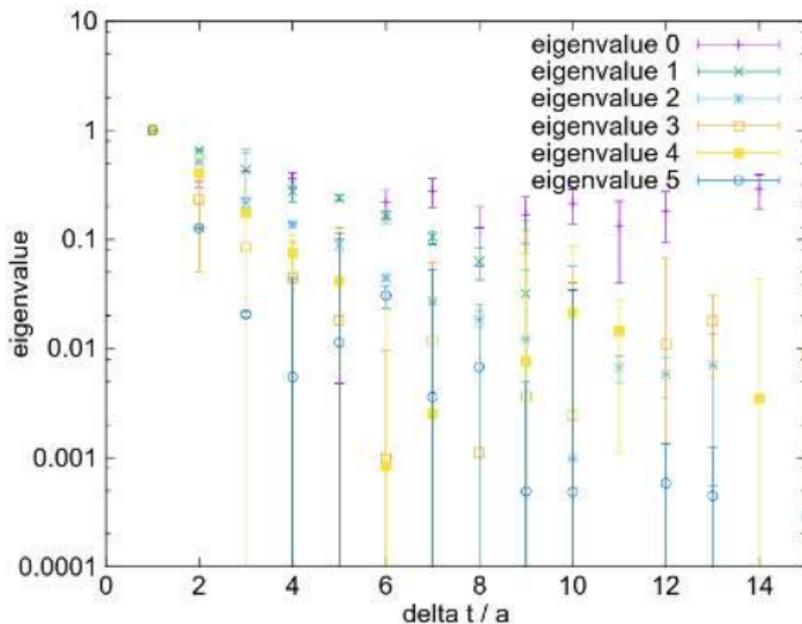
Generalized eigenvalues of 0^{++} -channel at $\beta = 1.9$, $\kappa = 0.1433$ sorted by value

Ordering the eigenvalues



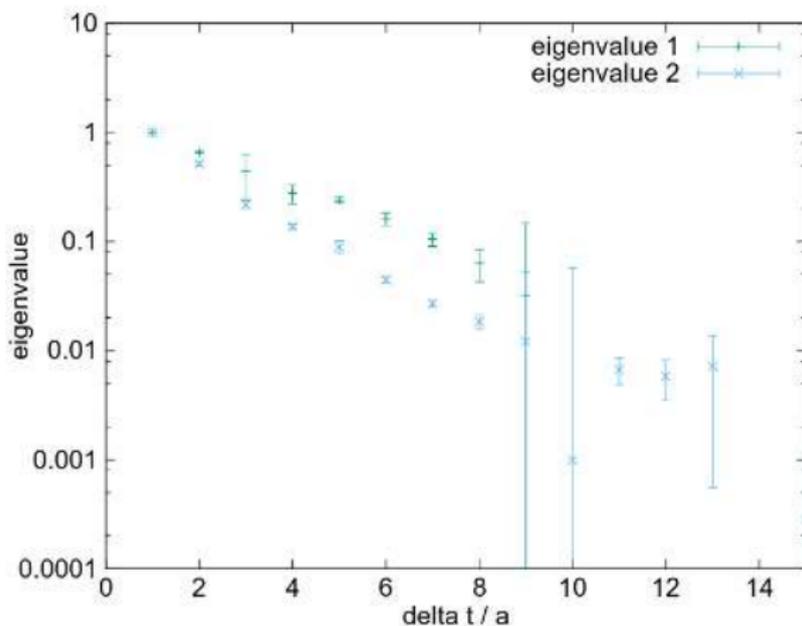
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Generalized eigenvalues of 0^{++} -channel at $\beta = 1.9$, $\kappa = 0.1433$ sorted by new method

- 1 Solve GEVP at some fixed t :
 - $C(t)\vec{w}^{(m)} = \lambda^{(m)}(t, t_0)C(t_0)\vec{w}^{(m)}$
- 2 Solve GEVP at all other t :
 - $C(t)\vec{v}^{(n)} = \lambda^{(n)}(t, t_0)C(t_0)\vec{v}^{(n)}$
- 3 Sort $\vec{v}^{(n)}$ by their scalar products with $\vec{w}^{(n)}$:
 - $M_{mn} = \vec{w}^{(m)} \cdot \vec{v}^{(n)}$
 - scan for largest entry in M , call the indices m_{\max}, n_{\max}
 - identify $\vec{v}_{\text{sorted}}^{(n_{\max})} = \vec{v}^{(m_{\max})}$
 - delete m -th row and n -th column from M
 - repeat 3. until all vectors are identified

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- Tree-level Symanzik improved gauge action
- 1 level of stout smearing
- Configurations produced on JUQUEEN - Juelich - Germany
- Measurements performed on
 - JURECA - Juelich - Germany
 - local HPC clusters in Münster, standard CPU-clusters and Xeon-Phi Knights Corner cluster

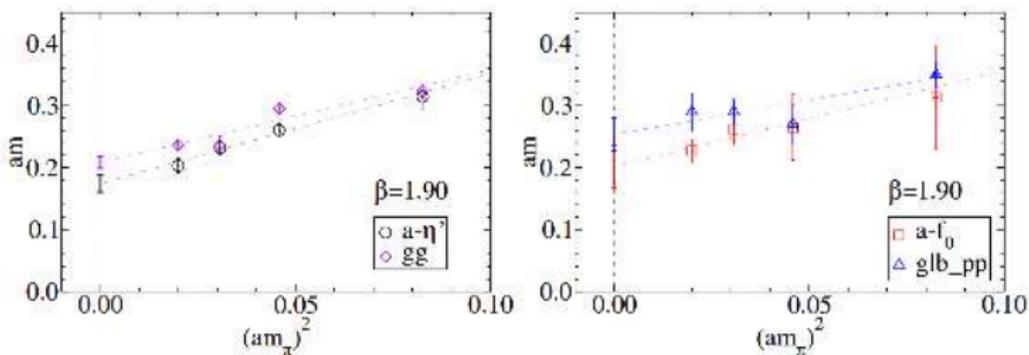
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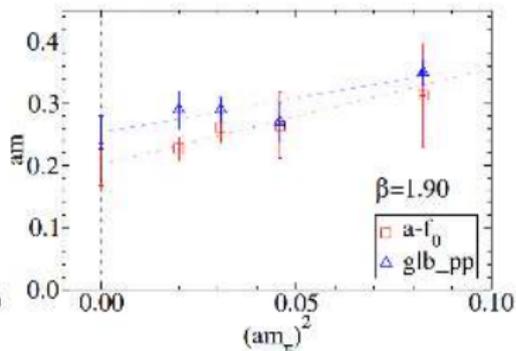
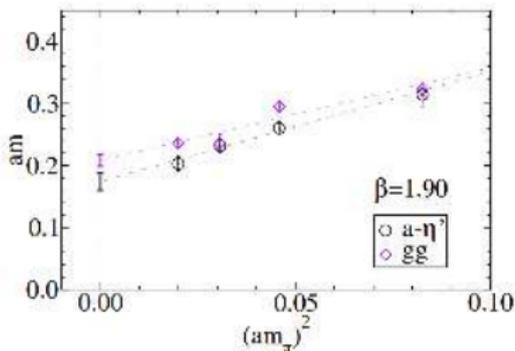
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Groundstate multiplet

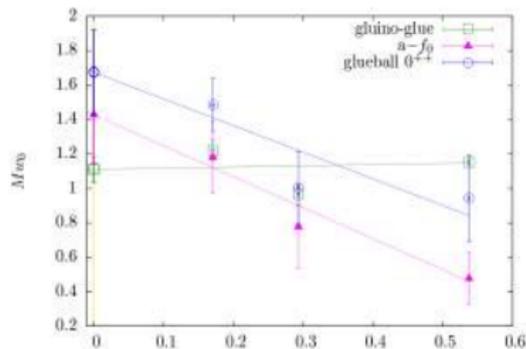
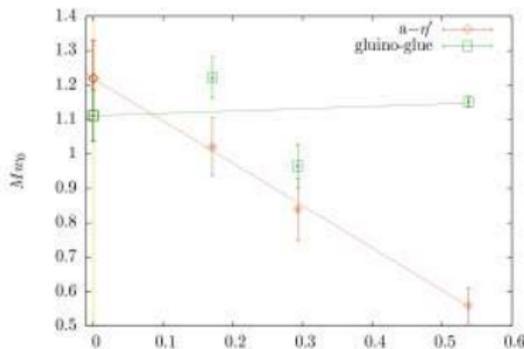


Extrapolation of the lightest supermultiplet to the chiral limit. Desy-Münster collaboration arXiv:1512.07014

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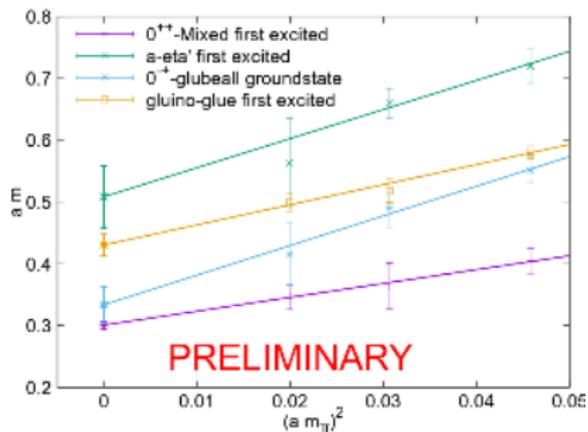
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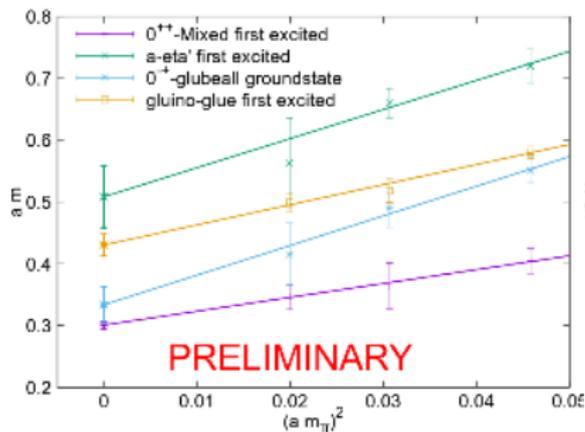
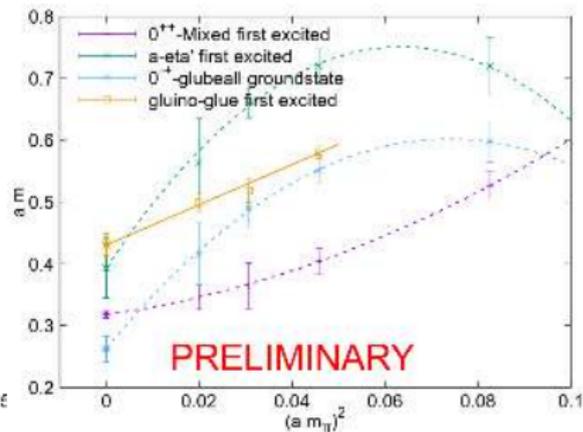
Excited Multiplet @ $\beta = 1.9$

PRELIMINARY!!!

Linear chiral extrapolation at $\beta = 1.9$

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PRELIMINARY!!!

Linear chiral extrapolation at $\beta = 1.9$ Quadratic chiral extrapolation at $\beta = 1.9$

- $\beta = 1.9$ (0.036fm)

κ	$N_{\text{measured}}^{\text{confs}}$ (mesos)	$N_{\text{available}}^{\text{confs}}$
0.1433	1151	10374
0.14387	1276	10237
0.14415	2814	21090
0.14435	1343	10680

- $\beta = 1.75$ (0.054fm)

κ	$N_{\text{measured}}^{\text{confs}}$ (mesons)	$N_{\text{available}}^{\text{confs}}$
0.1490	1163	9350
0.1492	1070	13017
0.14925	879	7224
0.1493	612	10280

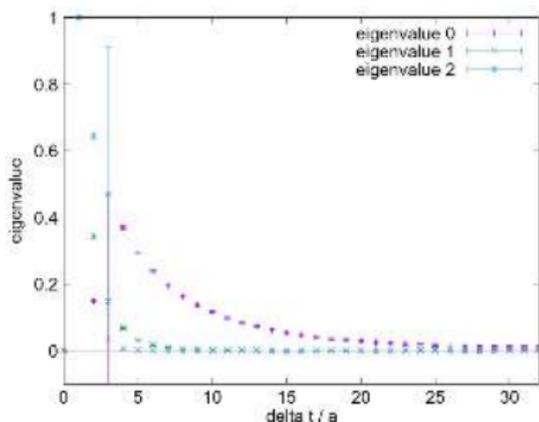
- still too noisy for extraction of first excited states

- Second order polynomial fit fits the data too well
 - errorbars overestimated?
- Within errors there seem so be no mixing in the 0^{-+} -channel
 - off-diagonal entries of the correlation matrix are zero within errors
- 0^{-+} -glueball and $a - \eta'$ masses differ

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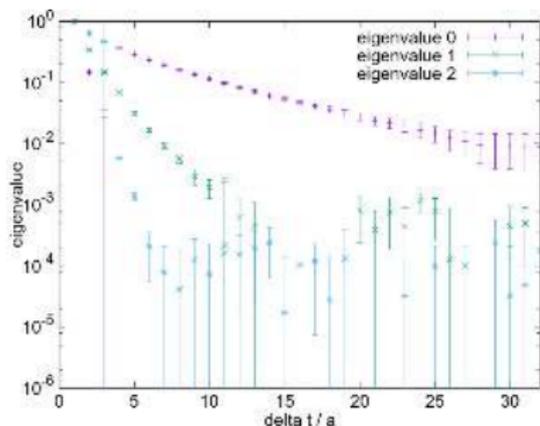
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- Large autocorrelation
- Talk C. Urbach, J. Simeth, Tuesday 18.10, 1830
 - Derivative trick
 \Rightarrow Use
 $\tilde{C}(t) = C(t) - C(t+1)$

$a-\eta'$ @ $\beta = 1.9$, $\kappa = 0.14435$

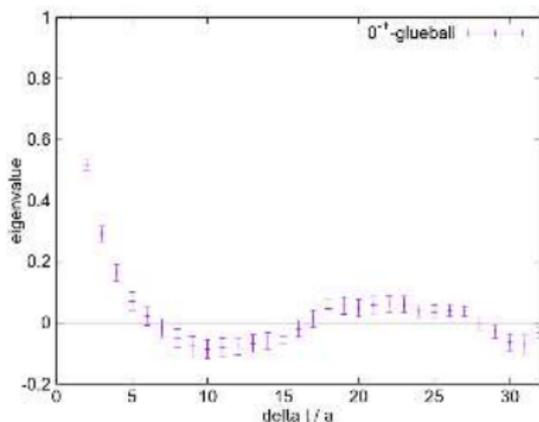
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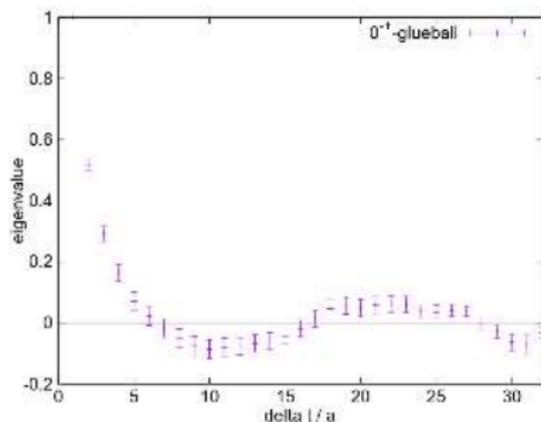
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- 0^{-+} -glueball and $a - \eta'$ masses differ



0^{-+} -glueball
 $@\beta = 1.9, \kappa = 0.14435$

- Large autocorrelation
- Talk C. Urbach, J. Simeth, Tuesday 18.10, 1830
 - Derivative trick
 \Rightarrow Use
 $\tilde{C}(t) = C(t) - C(t+1)$

- Second order polynomial fit fits the data too well
 - errorbars overestimated?
- Within errors there seem so be no mixing in the 0^{-+} -channel
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- Agreement with the predicted supermultiplet structure for the groundstate
⇒ SUSY is restored in the continuum limit
- 0^{-+} -glueball groundstate heavier than groundstate multiplet
⇒ might become lighter in the continuum limit
⇒ maybe it belongs to the first excited multiplet
- Work to be done:
 - Measurements for different β
 - Continuum extrapolation of the excited multiplet
 - reanalysis of groundstate multiplet using the full correlation matrix
 - Determine mixing between glueball and mesonic states using the variational method:
 - learn about mixing in the SUSY-phase

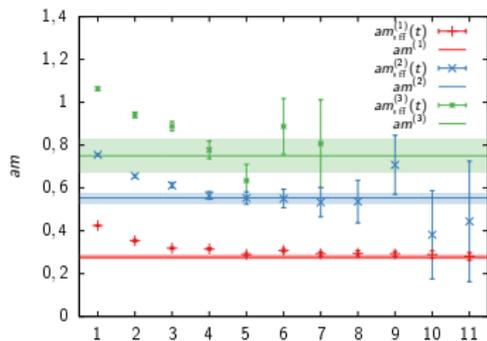
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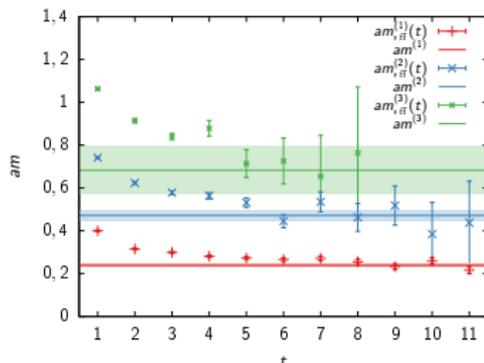
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Thank you for your attention!!

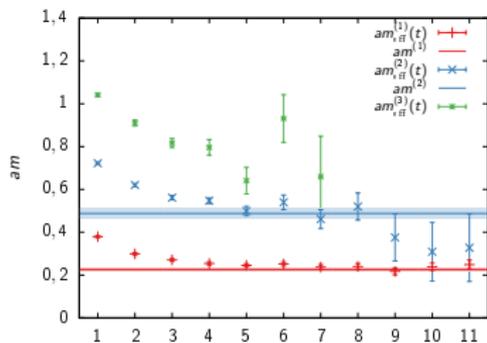




Effective gluino-gluon masses at
 $\kappa = 0.14387$



Effective gluino-gluon masses at
 $\kappa = 0.14415$



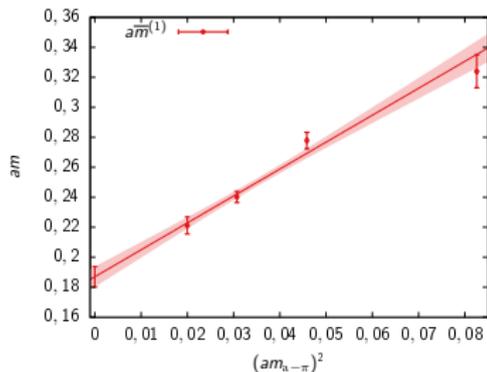
Effective gluino-gluon masses at
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*S. Kuberski, Masterthesis, Uni
Münster*

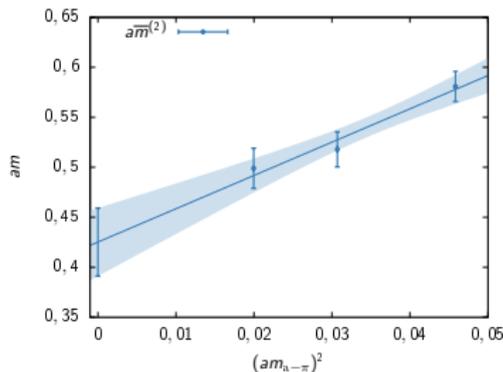
$$\kappa = \frac{1}{2am+8}$$

$$\beta = \frac{2N}{g^2}$$

$$m_{\text{eff}} = \ln \frac{C(\Delta t)}{C(\Delta t+1)}$$



Extrapolation to the chiral limit
(groundstate)



Extrapolation to the chiral limit (first
excited state)

S. Kuberski, Masterthesis, Uni Münster

- Calculate the lowest eigenvalues of Q and corresponding eigenvectors
 - using Arnoldi (ARPACK)
 - Chebyshev Polynomials of order 11
 - Even/Odd-Preconditioning
- Stochastic estimator technique for space orthogonal to the previously calculated eigenvectors:

- $\frac{1}{N_S} \sum_i^{N_S} |\eta^i\rangle \langle \eta^i| = \mathbb{1} + \mathcal{O}(\sqrt{N_S})$

- use \mathbb{Z}_4 -noise

- $Q |s^i\rangle = |\eta^i\rangle$

- $Q^{-1} = \frac{1}{N_S} \sum_i^{N_S} |s^i\rangle \langle \eta^i|$

- Conjugate gradient

- $N_S = 40$ for $\beta = 1.9$, $32^3 \times 64$