

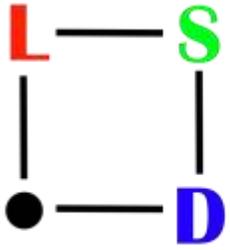
# From Walking Gauge Theory to Higgs Compositeness Using EFT

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In collaboration with James Ingoldby (Yale) and the Lattice Strong Dynamics (LSD) Collaboration

35<sup>th</sup> International Symposium on Lattice Field Theory  
June 20, 2017  
Granada

*Based on:*  
*arXiv:1702.00480 [1]*  
*and in progress work*



# Lattice Strong Dynamics Collaboration



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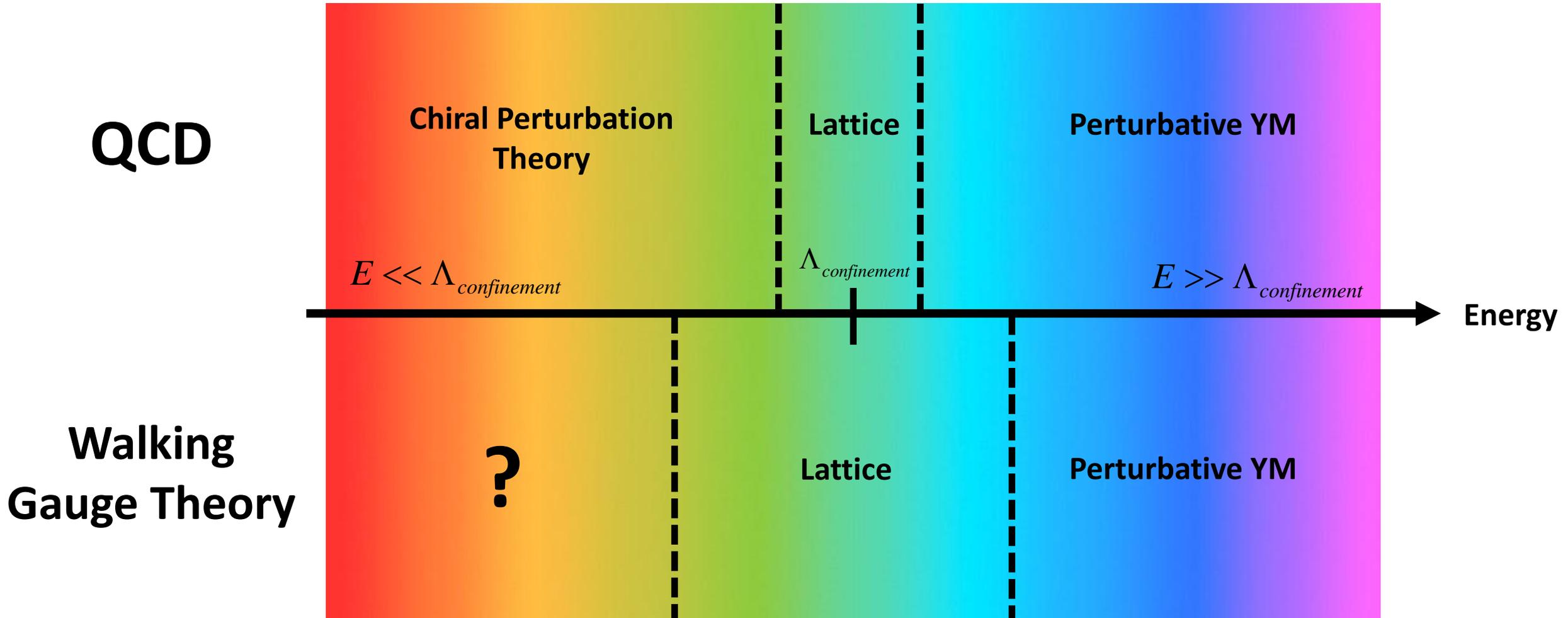
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# Confining Gauge Theories at High and Low Energies



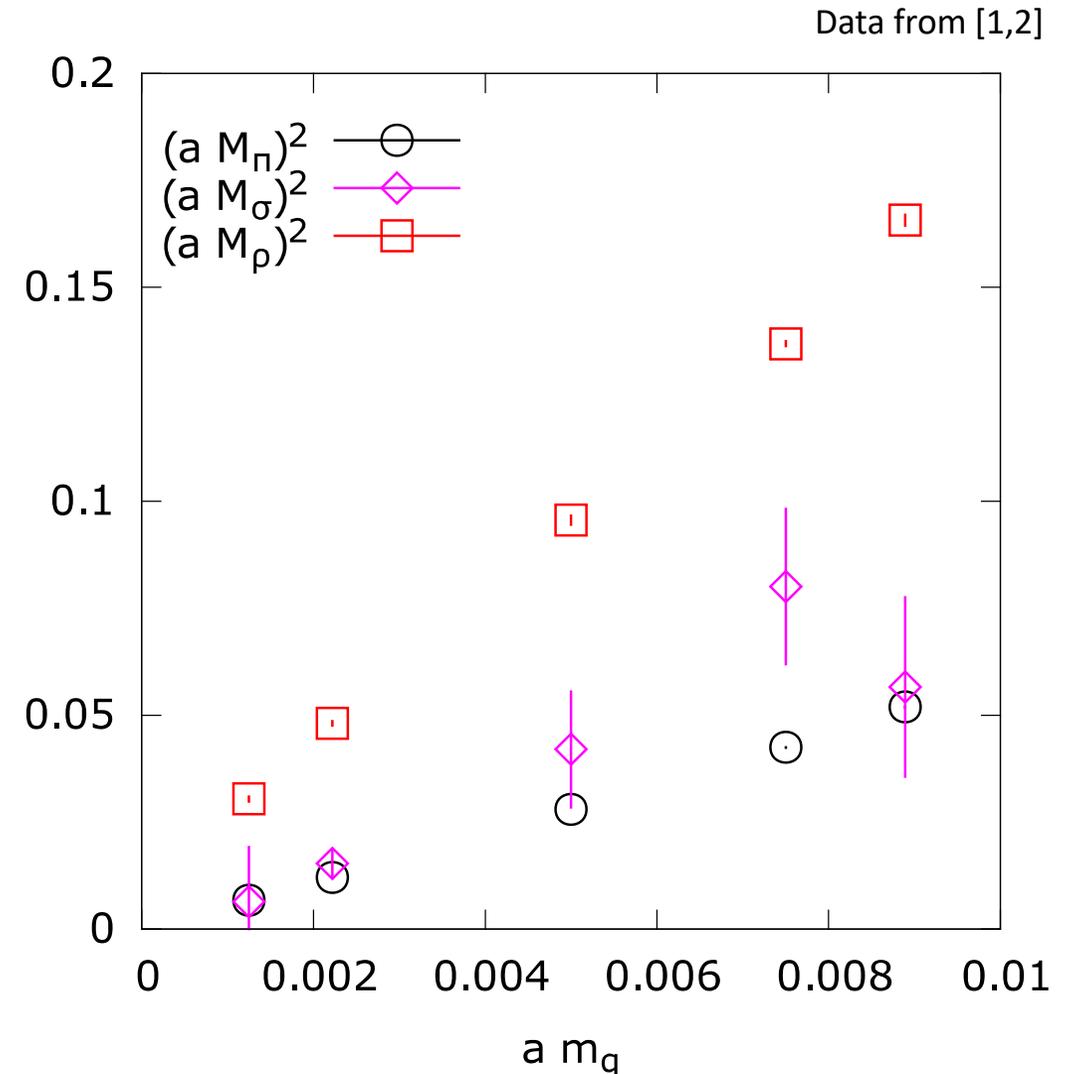
# Distinction between Walking EFT and Higgs EFT

- To make connection to Standard Model, eventually have to break flavor groups down to Standard Model global symmetry
- Many walking theories have a large flavor group with many PNGBs [2]
- Some other walking models do not come with additional PNGBs. For example, “minimal composite Higgs” [3][4]



# Spectrum of $N_f=8$ QCD and implications for EFT

- SU(3) gauge theory with 8 flavors in fundamental rep
  - Recent studies of 8f spectrum by latKMI [5] and LSD [1,2]
- Pions and Sigma nearly degenerate and well separated from the rho (and other heavier states, not shown)
- Sigma cannot be “integrated out” of effective theory at these distances from chiral limit
- Chiral extrapolations cannot be performed without EFT
  - Particularly important:  $\frac{M_\sigma}{F_\pi}$
- Seek EFT description that is good up to next resonance,  $M_\rho$
- Improved studies of systematic errors in 8f spectrum have been discussed in talk by George Fleming at 15:20 today
- The most recent progress on studying the sigma mass in 8f QCD has been discussed in talk by Enrico Rinaldi at 15:00 today



# Approximate Empirical Laws

- Light states adhere approximately to simple empirical laws:

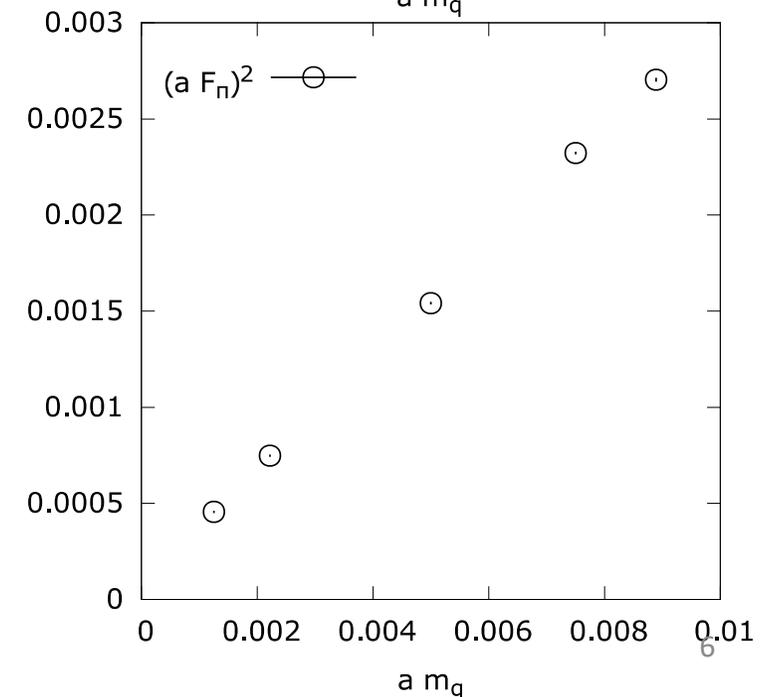
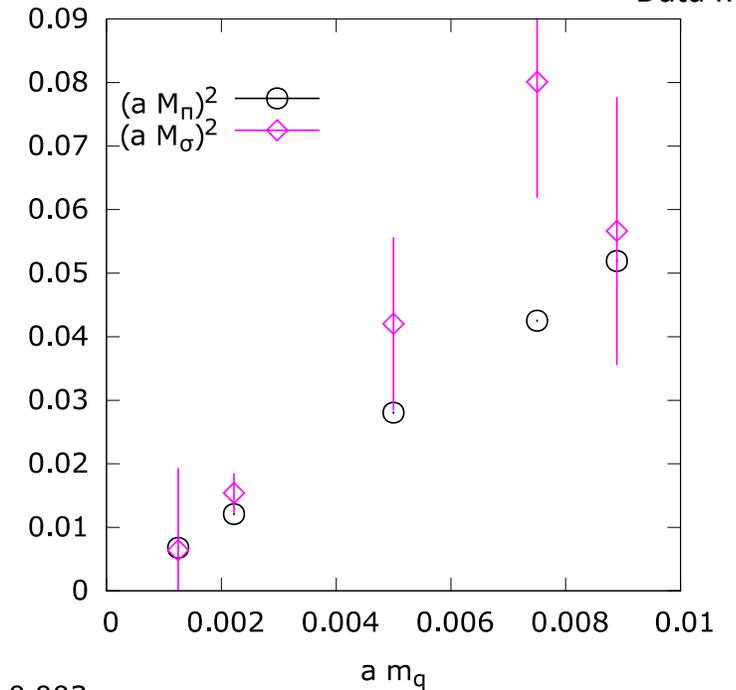
$$M_\pi^2 = B_\pi m_q \quad F_\pi^2 = f_\pi^2 + B_f m_q \quad M_\sigma^2 = m_\sigma^2 + B_\sigma m_q$$

- In these mass ranges,

$$f_\pi^2 \ll B_f m_q$$

- Expect significant qualitative changes when this inequality is saturated
- Empirical “law” for sigma somewhat dubious because of large errors
- Compare to LO Chiral Perturbation Theory
 
$$M_\pi^2 = B_\pi m_q \quad F_\pi^2 = f_\pi^2$$
- XPT cannot reproduce significant linear behavior in decay constant
  - This was discussed in detail in proceedings [1]

Data from [1,2]



# EFT approaches Beyond XPT for Walking Gauge Theories

	No Scale Symmetry	Scale Symmetry
No Special Multiplet	<ul style="list-style-type: none"><li>• XPT + Light Isosinglet Scalar [6,7]</li><li>• Couplings between sigma and pions general<ul style="list-style-type: none"><li>• Many LECs, more difficult to constrain with limited lattice data</li></ul></li></ul>	<ul style="list-style-type: none"><li>• “Dilatonic XPT” [8,9,10]</li><li>• Sigma couples to XPT as a compensating field</li><li>• Couplings between sigma and pions constrained by scale invariance</li><li>• Recently tested against lattice data with promising initial results [11]</li></ul>
Linear Multiplet	<ul style="list-style-type: none"><li>• Linear Sigma EFT<ul style="list-style-type: none"><li>• SU(N) formulation by Schechter and Ueda[12] for QCD</li><li>• This work: application to walking gauge theory</li></ul></li><li>• Pions and Sigma transform together under flavor group in linear multiplet</li></ul>	<ul style="list-style-type: none"><li>• “Dilatonic Linear Sigma EFT”<ul style="list-style-type: none"><li>• This work</li></ul></li><li>• Pions and Sigma transform together under flavor group in linear multiplet AND scale invariance is softly broken</li></ul>

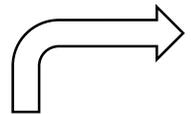
Linear Multiplet



# $SU_L(N_f) \times SU_R(N_f)$ Linear Multiplet: Field Content

- Fields live in bifundamental representation:  $M_a^{\bar{b}}$ 
  - $2N_f^2$  Real degrees of freedom
  - But one Goldstone degree of freedom must be excluded due to anomaly

- Transform Linearly under  $SU_L(N_f) \times SU_R(N_f)$ :  $M \rightarrow LMR^\dagger$

- Express degrees of freedom via a polar decomposition (CCWZ):   $T_a \in su(N_f)$

$$M = \Sigma S \quad \Sigma = \exp(i\pi_a T_a / F_\pi) \quad S = \frac{\sigma}{\sqrt{N_f}} + a_a T_a$$

- Full symmetry is nonlinearly realized in this basis
- Without flavored scalars, fields do not form a complete linear representation
  - $N_f = 2$  is special. Admits a real linear representation due to isometry with  $O(4)$

# $SU_L(N_f) \times SU_R(N_f)$ Linear Multiplet: Lagrangian

- First, write down most general action symmetric under flavor group and parity, including all relevant and marginal operators

$$\mathcal{L}_{LSM} = \frac{1}{2} \langle \partial_\mu M \partial^\mu M^\dagger \rangle - V_0(M^\dagger M) - V_{SB}(M, \mathcal{M})$$

$$V_0 = -\frac{1}{4} m_\sigma^2 \langle M^\dagger M \rangle + \frac{m_\sigma^2 - m_a^2}{8f^2} \langle M^\dagger M \rangle^2 + \frac{N_f m_a^2}{8f^2} \langle (M^\dagger M)^2 \rangle$$

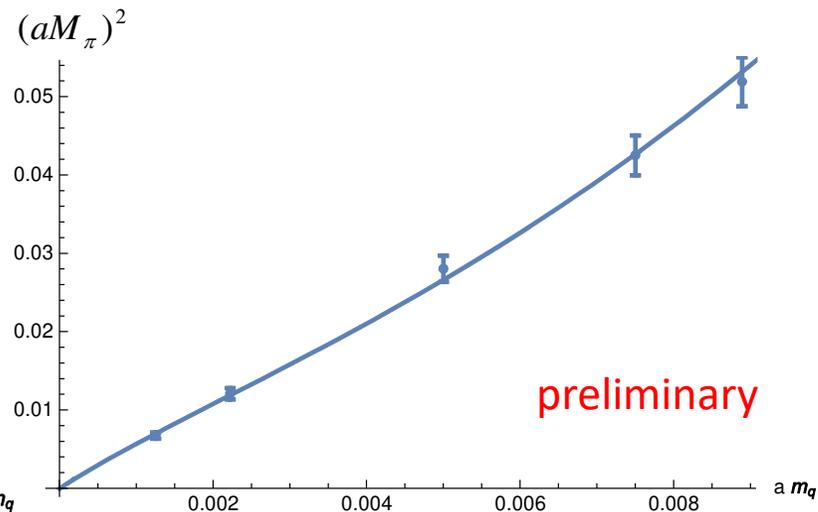
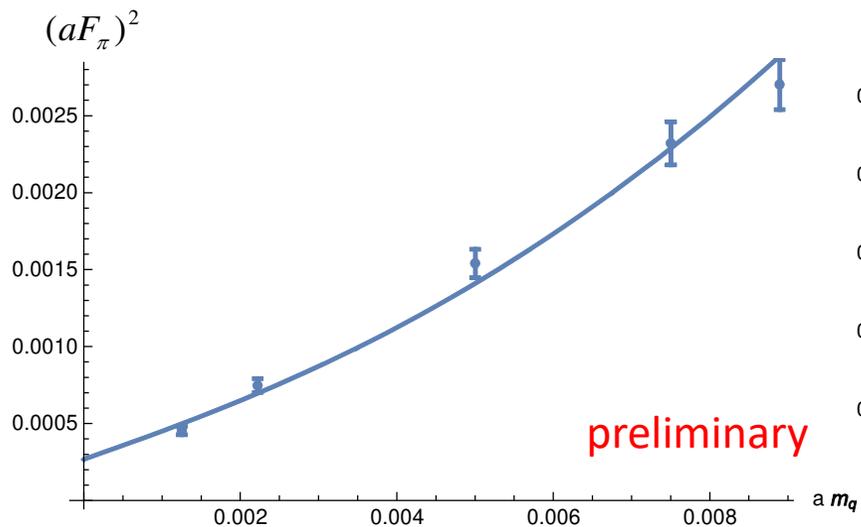
Note:

$$F_\pi = \frac{\langle 0 | \sigma | 0 \rangle}{\sqrt{N_f}}$$

- Breaking potential determined by spurion analysis (similar to XPT)

$$V_{SB} = \underbrace{\frac{-f}{\sqrt{N_f}} \langle X^\dagger M + XM^\dagger \rangle}_{O(X^1, M^1)} - \underbrace{\frac{c_1}{f\sqrt{N_f}} \langle M^\dagger M \rangle \langle X^\dagger M + XM^\dagger \rangle - \frac{c_2 \sqrt{N_f}}{f} \langle (M^\dagger M) X^\dagger M + (MM^\dagger) XM^\dagger \rangle}_{O(X^1, M^3)}$$

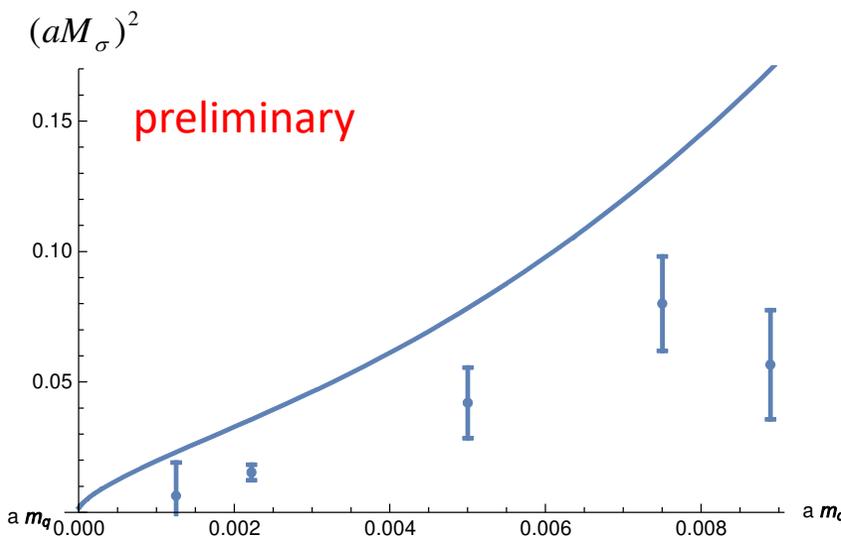
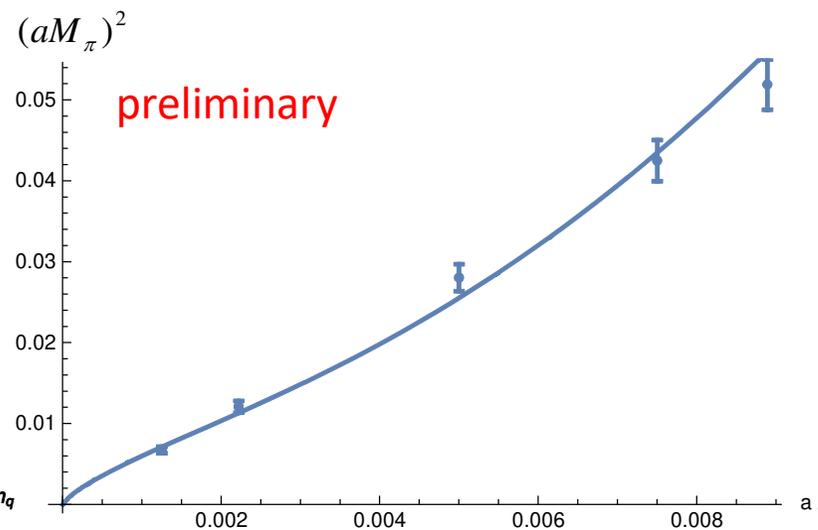
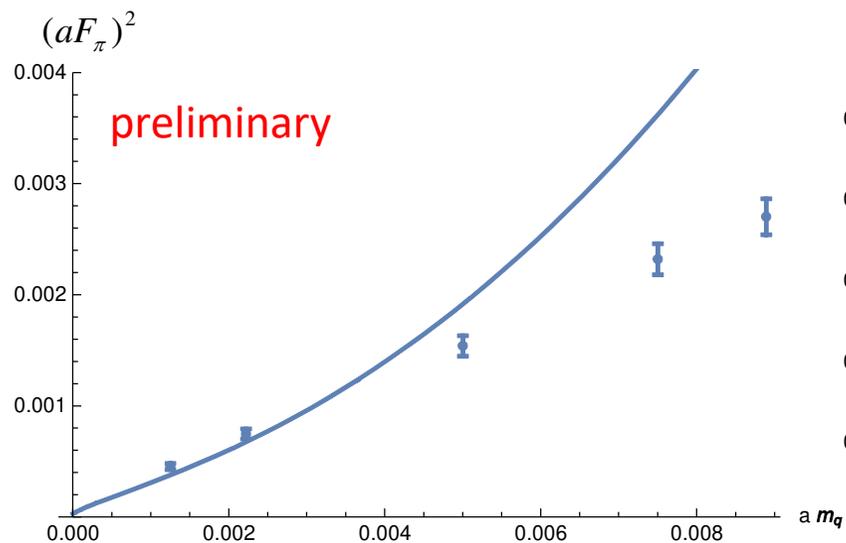
# Fits of $F_\pi$ , $M_\pi$ , and $M_\sigma$ to linear sigma EFT up to $O(X^1, M^3)$ , tree level



$\chi^2/\text{dof} = 1.2$   
With 3% Systematic Errors

*Fitting all 5 quark masses:  $\chi^2/\text{dof} = 12.4$   
With 3% systematic errors*

*Fitting first 3 quark masses:  $\chi^2/\text{dof} = 10.9$   
With 3% systematic errors*



# Linear Multiplet + Scale Invariance: Lagrangian

- Write down most general action symmetric under flavor group, parity, and classical scale invariance, including only marginal operators

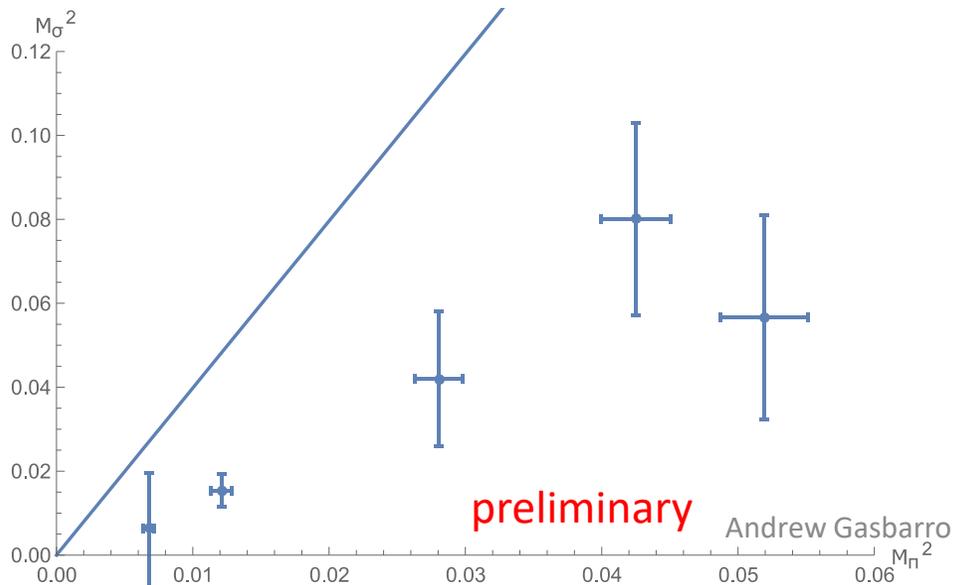
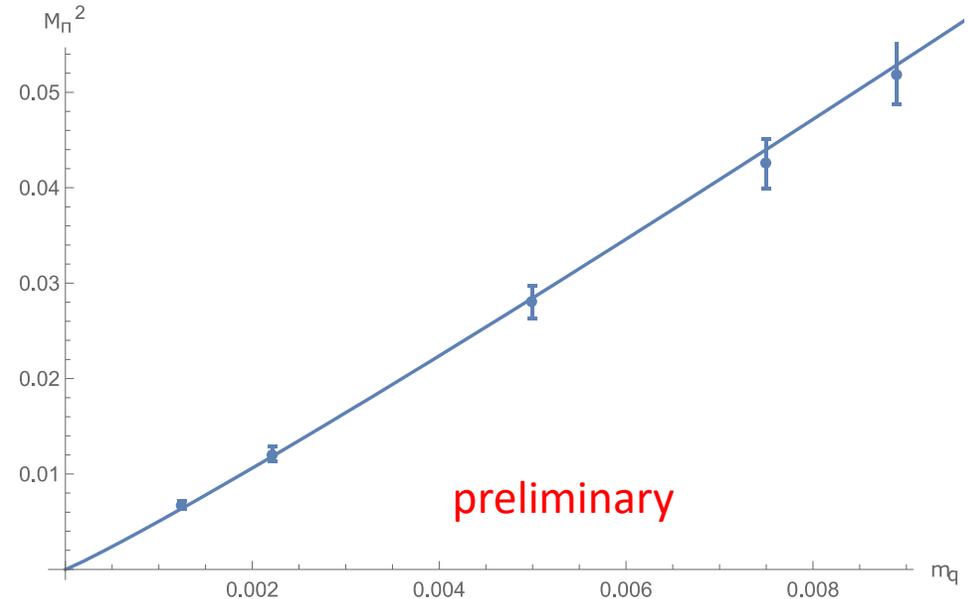
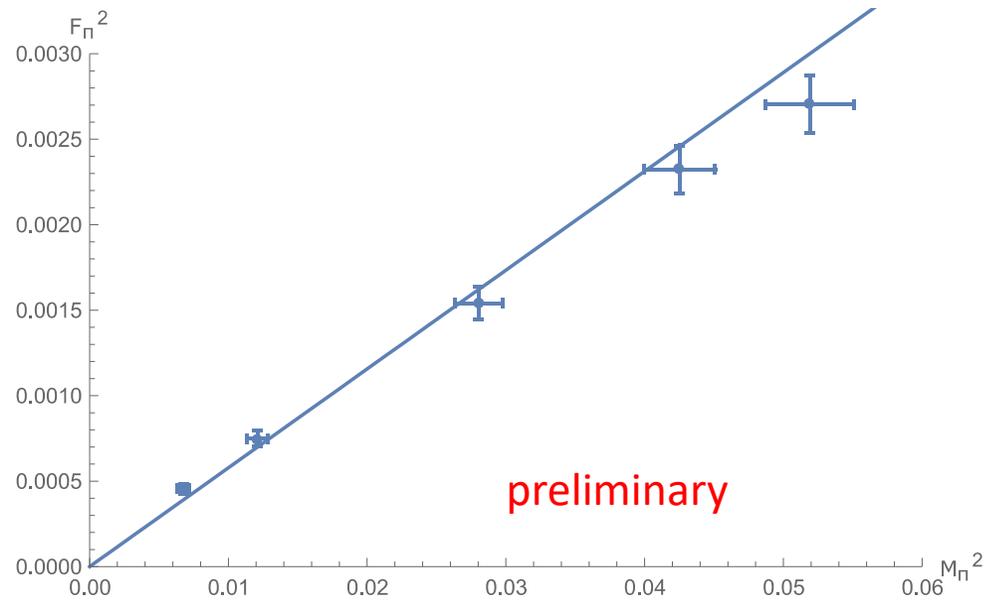
$$\mathcal{L}_{LSM} = \frac{1}{2} \langle \partial_\mu M \partial^\mu M^\dagger \rangle - V_0(M^\dagger M) - V_{SB}(M, \mathcal{M})$$

$$V_0 = \frac{m_\sigma^2 - m_a^2}{8f^2} \langle M^\dagger M \rangle^2 + \frac{N_f m_a^2}{8f^2} \langle (M^\dagger M)^2 \rangle$$

- Breaking potential determined by spurion analysis. Spurions transform under scale symmetry as well!

$$V_{SB} = -\frac{m_\sigma^2}{4} \langle M^\dagger M \rangle + f \left( \frac{\langle M^\dagger M \rangle}{f^2} \right)^b \langle X^\dagger M + XM^\dagger \rangle$$

# Fits of $F_\pi$ , $M_\pi$ , and $M_\sigma$ to scale invariant linear sigma, tree level



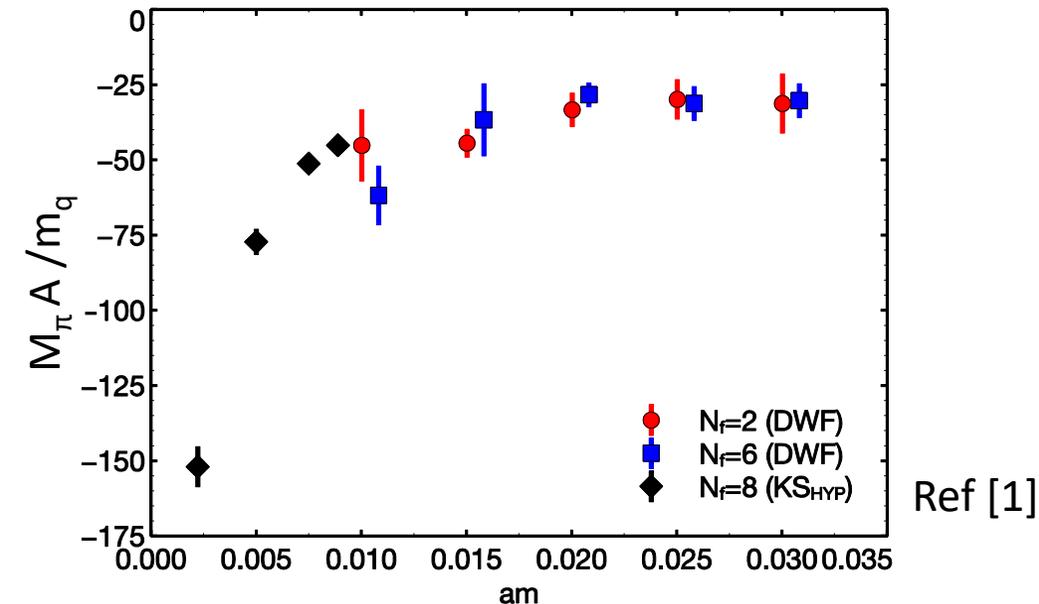
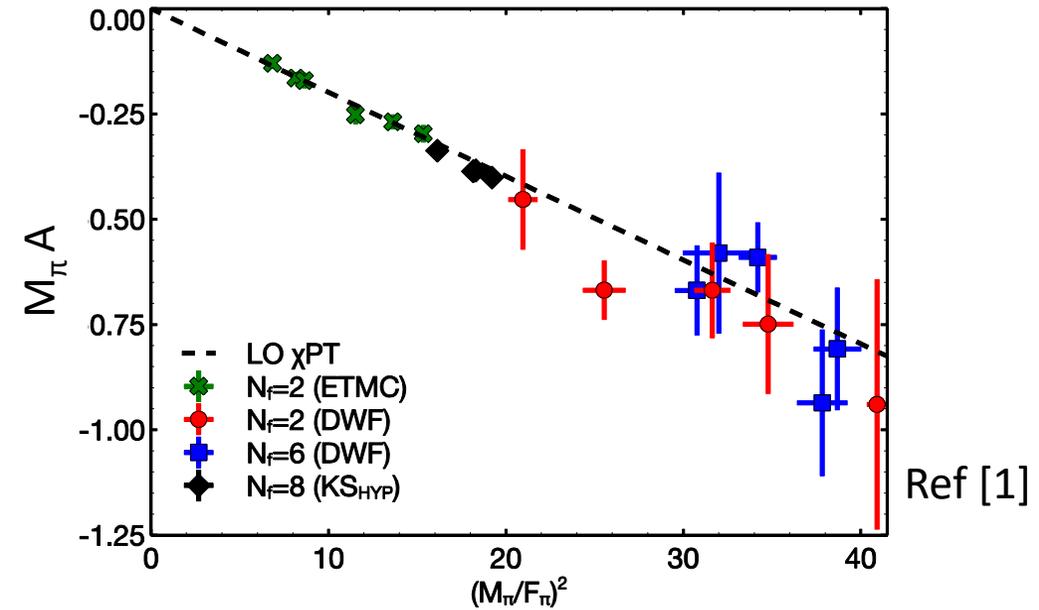
$\chi^2/\text{dof} = 7.5$   
With 3% Systematic Errors

## More Information from Pion Scattering

- Scattering length agrees well with LO XPT when plotted against physical (computed) values of  $M_\pi/F_\pi$

$$(M_\pi A)_{LO} = \frac{-1}{(16\pi)^2} \left( \frac{M_\pi}{F_\pi} \right)^2 = \frac{-1}{(16\pi)^2} \frac{2Bm_q}{F^2}$$

- Plotted against bare quark mass, very poor agreement with LO XPT
- Scattering data in process of being analyzed with various EFT options



# Conclusions & Outlook

- Linear Sigma EFT provides significant improvement over XPT in fitting PNGB data,  $F_\pi$  and  $M_\pi$
- Tree level expressions do not provide good combined fit to  $F_\pi$ ,  $M_\pi$  and  $M_\sigma$ 
  - Loops may not be negligible
- More work being done on the lattice to further constrain EFT
  - Pion scattering
  - Form factors (in progress)

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Backup Slides

# An Initial Analysis: Fitting to XPT

- At leading order in Chiral Perturbation Theory

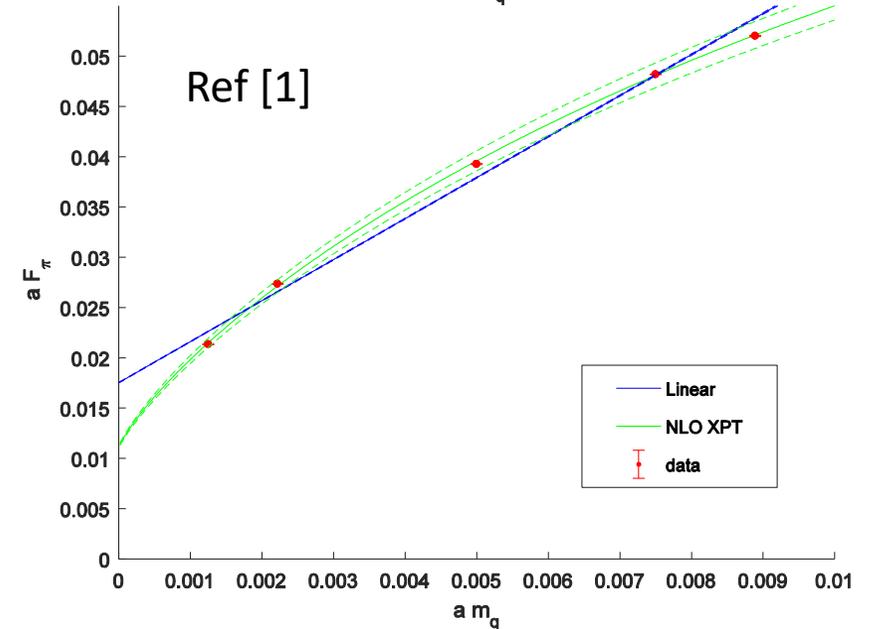
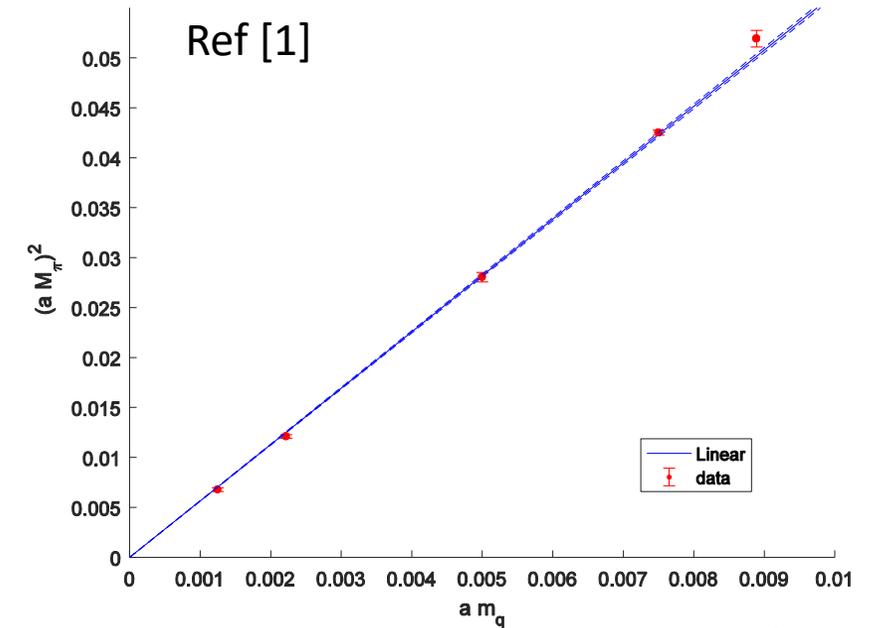
$$M_\pi^2 = 2Bm_q \quad F_\pi = F \text{ (constant)}$$

- Mass looks strikingly like LO XPT
- F deviates significantly from LO XPT
- Combined fit to NLO XPT fails

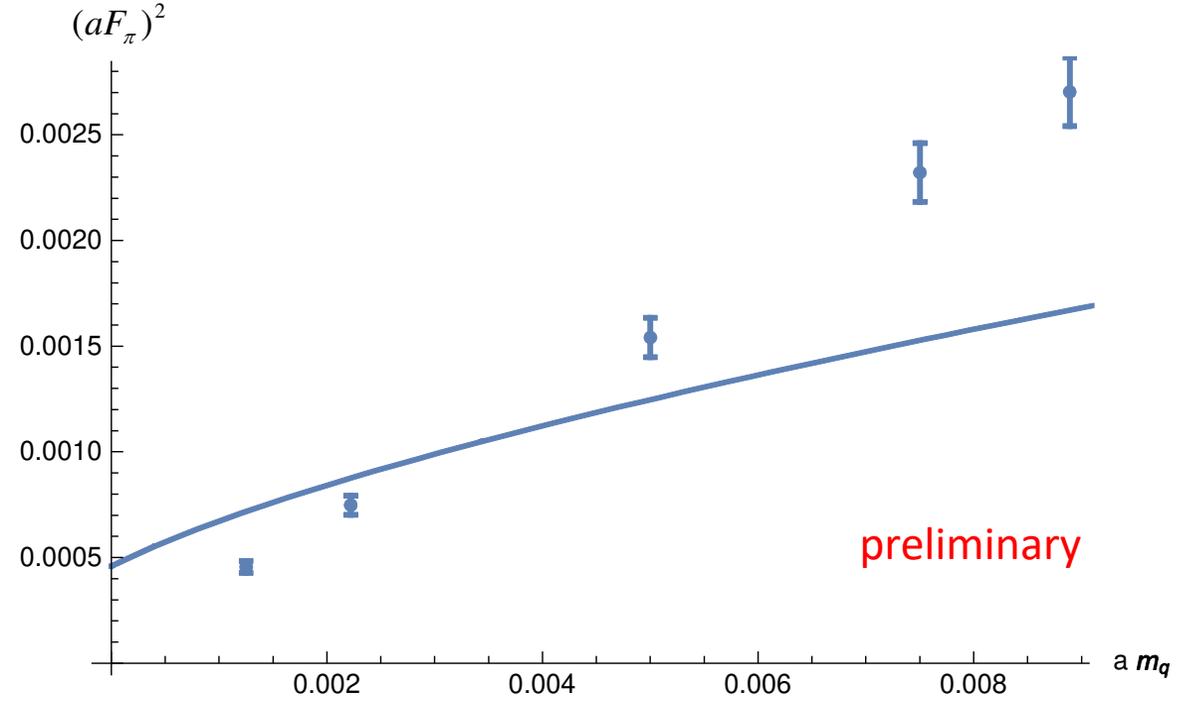
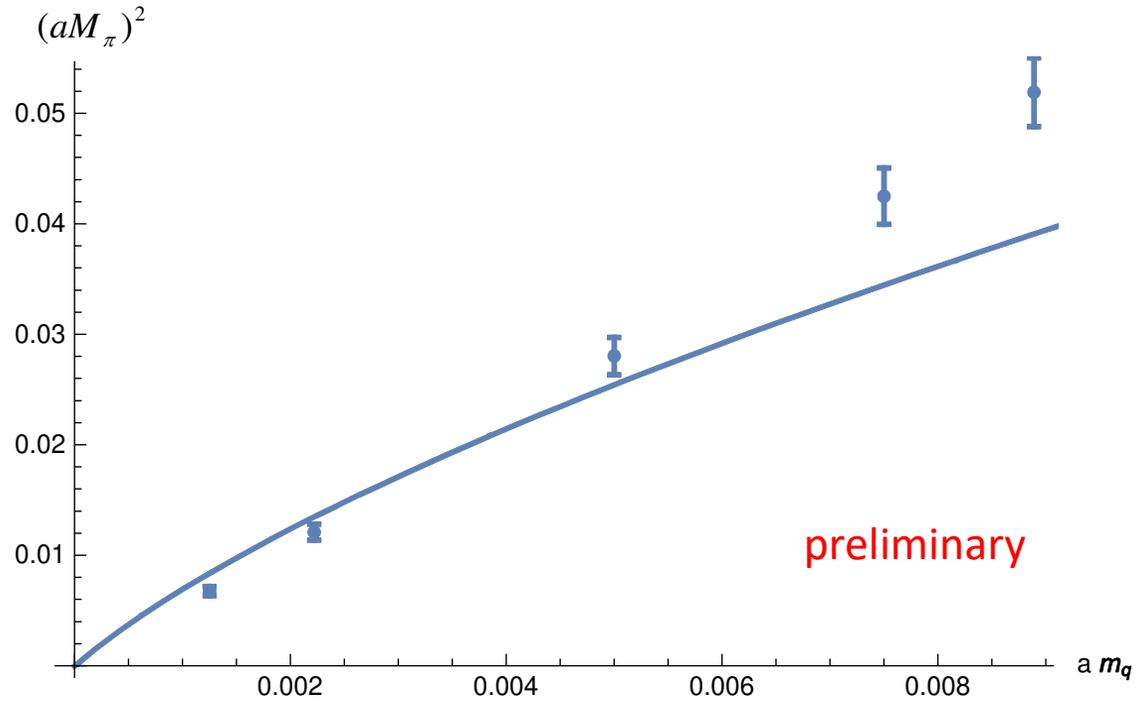
- Loop expansion parameter:
  - Determine “B” from slope of pion mass
  - Determine “F” from intercept of decay constant

$$\frac{2N_f B m_q}{(4\pi F)^2} \approx 5.8$$

- Incompatibility with XPT was pointed earlier in Nf=6 study [14]



# Fits of $M_\pi$ and $F_\pi$ to linear sigma up to $O(X^1, M^1)$



$\chi^2/\text{dof} = 37.6$   
With 3% systematic errors

# Explicit Breaking in Underlying Gauge Theory

- Introduce operators to explicitly break symmetry to  $SU_V(N_f)$

$$V_{SB} = c_1 \langle M + M^\dagger \rangle + c_2 \langle M^2 + (M^\dagger)^2 \rangle + c_3 \left( \langle M \rangle^2 + \langle M^\dagger \rangle^2 \right) + c_4 \langle M \rangle \langle M^\dagger \rangle + c_5 \langle M^3 + (M^\dagger)^3 \rangle + c_6 \langle M^2 M^\dagger + (M^\dagger)^2 M \rangle + \dots$$

- Can reduce number of breaking terms by looking to underlying gauge theory for guidance.

- Only source of explicit breaking in Yang Mills model is quark mass.

$$\mathcal{L}_{YM} \supset \psi_L^\dagger \mathcal{M} \psi_R + \psi_R^\dagger \mathcal{M}^\dagger \psi_L$$

- Spurion analysis: treat chiral breaking as an external source which transforms

$$X \rightarrow LX R^\dagger$$