

Baryonic and mesonic 3-point functions with open spin indices

M. Löffler for RQCD

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Universität Regensburg



What is LHA?

A new LQCD library for efficient calculations on modern parallel architectures and on CPUs

What can be done with LHA?

BDA, MDA

Baryon Spectrum

Baryon 3pt Functions

Meson 3pt Functions

Why use LHA?

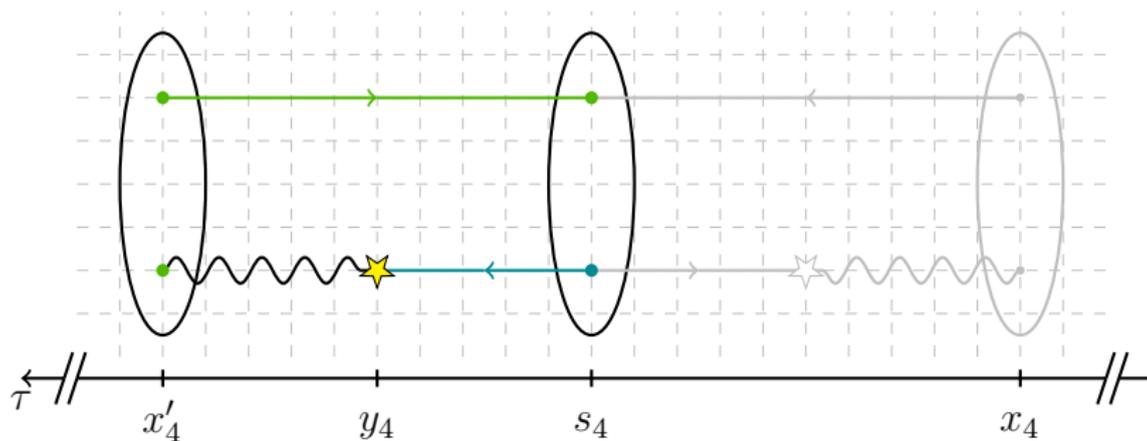
Standard implementations do not work efficiently on new architectures

Adaptable to every physical channel within data analysis

Computation per lattice site → SIMD data layout

Parallelization scheme

Meson 3pt Functions on the Lattice using stochastic Estimators



Operators

$$\text{Annihilator: } A(\mathbf{x}', x'_4) = \delta_{\mathbf{ab}} \bar{F}_{\mathbf{A}}(\mathbf{x}', x'_4)_{\mathbf{a}}^{\alpha} \Gamma'^{\alpha\beta} F_{\mathbf{B}}(\mathbf{x}', x'_4)_{\mathbf{b}}^{\beta}$$

$$\text{Creator: } C(\mathbf{s}, s_4) = \delta_{ba} \bar{F}_{\mathbf{B}}(\mathbf{s}, s_4)_{\mathbf{b}}^{\beta} (\gamma_4 \Gamma^{\dagger} \gamma_4)^{\beta\alpha} F_{\mathbf{A}}(\mathbf{s}, s_4)_{\mathbf{a}}^{\alpha}$$

$$\text{Current: } I(\mathbf{y}, y_4) = \delta_{\mathbf{ab}} \bar{F}_{\mathbf{A}}(\mathbf{y}, y_4)_{\mathbf{a}}^{\alpha} \Gamma_{\mathbf{J}}^{\alpha\beta} F_{\mathbf{B}}(\mathbf{y}, y_4)_{\mathbf{b}}^{\beta} \quad (+ \text{ derivative Operators})$$

with $F_i \in \{L, S, C\}$ (light, strange, charm).

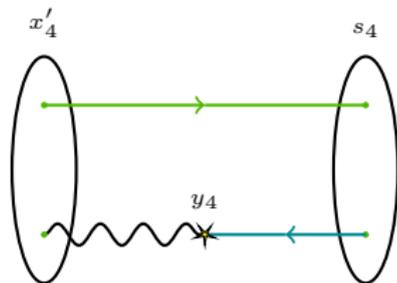
 Γ -Structure examples:

Scalar: $\mathbb{1}, \gamma_4$

Pseudoscalar: $\gamma_5, \gamma_4 \gamma_5$

Vector: $\gamma_i, \gamma_4 \gamma_i$

Axial Vector: $\gamma_i \gamma_5$

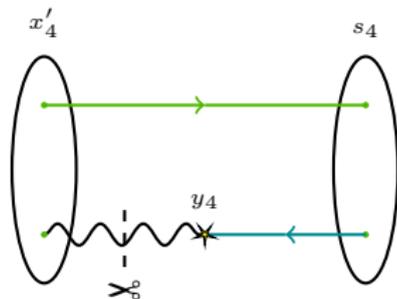


Performing the generic Wick contraction

$$\begin{aligned}
 \langle A(x') I(y) C(s) \rangle &= \\
 &= \text{tr} [G(F_A, s, x') \Gamma' G(F_B, x', y) \Gamma_{\mathcal{J}} G(F_{\mathfrak{B}}, y, s) \Gamma] \\
 &= \delta_{\mathbf{ab}} \delta_{ab} \delta_{\mathbf{ab}} (\Gamma')^{\alpha\beta} (\Gamma_{\mathcal{J}})^{\mathbf{ab}} (\Gamma)^{\beta\alpha} \times \\
 &\quad \times G(F_A, s, x')_{\mathbf{aa}}^{\alpha\alpha} G(F_B, x', y)_{\mathbf{ba}}^{\beta\mathbf{a}} G(F_{\mathfrak{B}}, y, s)_{\mathbf{bb}}^{\mathbf{b}\beta}
 \end{aligned}$$

The stochastic all-to-all propagator

$$\begin{aligned}
 G(F_A, y, x')_{\mathbf{ba}'}^{\beta\alpha'} &\approx \\
 \frac{1}{N} \sum_{i=1}^N (s_i)(F_A, y)_b^{\beta} (\eta_i^*) (x')_{\mathbf{a}'}^{\alpha'}
 \end{aligned}$$



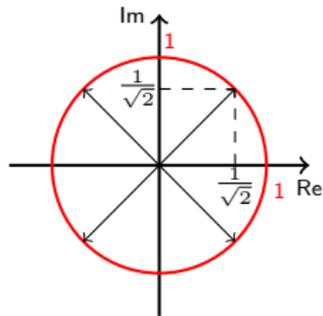
Vector – Definition¹

$$(\eta_i)(\mathbf{x}, x_4)_a^\alpha \in \begin{cases} (\mathbb{Z}_2 \otimes i\mathbb{Z}_2)/\sqrt{2} & , x_4 = \text{sink fwd/bwd} \\ 0 & , \text{otherwise} \end{cases}$$

Vector – Properties

$$\frac{1}{N} \sum_{i=1}^N (\eta_i)(x)_a^\alpha = 0 + \mathcal{O}\left(\frac{1}{\sqrt{N}}\right)$$

$$\frac{1}{N} \sum_{i=1}^N (\eta_i)(x)_a^\alpha (\eta_i^*)(x')_{a'}^{\alpha'} = \delta_{xx'} \delta_{\alpha\alpha'} \delta_{aa'} + \mathcal{O}\left(\frac{1}{\sqrt{N}}\right)$$



¹S.-J. Dong and K.-F. Liu, "Stochastic estimation with Z(2) noise," Phys. Lett. B328 (1994) 130, hep-lat/9308015

The 3pt Function becomes

Use that:

$$\langle A(x') I(y) C(s) \rangle =$$

$$G(x, y)_{ab}^{\alpha\beta} = (\gamma_5^*)^{\beta\gamma} G^*(y, x)_{ba}^{\gamma\delta} (\gamma_5^*)^{\delta\alpha}$$

$$\begin{aligned} &= \delta_{\mathbf{ab}} \delta_{ab} \delta_{\mathbf{ab}} (\Gamma')^{\alpha\beta} (\Gamma_{\mathcal{J}})^{\mathbf{ab}} (\Gamma)^{\beta\alpha} \times \\ &\quad \times G(F_A, s, x')_{\mathbf{aa}}^{\alpha\alpha} G(F_B, x', y)_{\mathbf{ba}}^{\beta\mathbf{a}} G(F_{\mathfrak{B}}, y, s)_{\mathbf{bb}}^{\mathbf{b}\beta} \\ \\ &= \delta_{\mathbf{ab}} \delta_{ab} \delta_{\mathbf{ab}} (\Gamma')^{\alpha\beta} (\Gamma_{\mathcal{J}})^{\mathbf{ab}} (\Gamma)^{\beta\alpha} \times \\ &\quad \times \underbrace{[\gamma_5^* G(F_A, x', s)^* \gamma_5^*]_{\mathbf{aa}}^{\alpha\alpha} [\eta_i(x') \gamma_5^*]_{\mathbf{b}}^{\beta}}_{\hat{=} \text{Spectator}} \underbrace{[\gamma_5^* s_i(F_B, y)]_{\mathbf{a}}^{\mathbf{a}} G(F_{\mathfrak{B}}, y, s)_{\mathbf{bb}}^{\mathbf{b}\beta}}_{\hat{=} \text{Insertion}} \end{aligned}$$

At the end we achieve an expression for the mesonic 3pt function

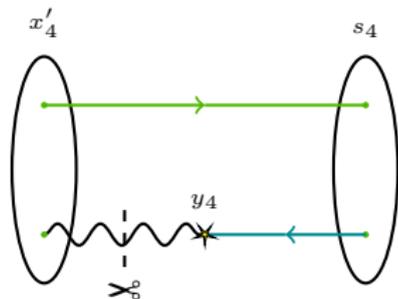
$$C^{F_A F_B F_{\mathfrak{B}} \alpha \alpha \beta a b \beta}(\mathbf{p}, x'_4, \mathbf{q}, y_4) = \frac{1}{N} \sum_i S_i^{F_A \alpha \alpha \beta}(\mathbf{p}, x'_4) I_i^{F_B F_{\mathfrak{B}} a b \beta}(\mathbf{q}, y_4)$$

$$S_i^{F_A \alpha \alpha \beta}(\mathbf{p}, x'_4) = \sum_{\mathbf{x}'} \delta_{ab} [\gamma_5^* G(F_A, x', s)^* \gamma_5^*]_{aa}^{\alpha \alpha} [\eta_i(x') \gamma_5^*]_{\mathbf{b}}^{\beta} \cdot e^{-i\mathbf{p} \cdot \mathbf{x}'}$$

$$I_i^{F_B F_{\mathfrak{B}} a b \beta}(\mathbf{q}, y_4) = \sum_{\mathbf{y}} \delta_{ab} \delta_{ab} [\gamma_5^* s_i(F_B, y)]_{\mathbf{a}}^{\mathbf{a}} G(F_{\mathfrak{B}}, y, s)_{\mathbf{b}\mathbf{b}}^{\mathbf{b}\beta} \cdot e^{i\mathbf{q} \cdot \mathbf{y}}$$

Contracting the spectator and insertion parts with the correct Γ -structure gives the desired result

What does this mean for computational purposes?



Spectator size (upper bound)

Spectator size =

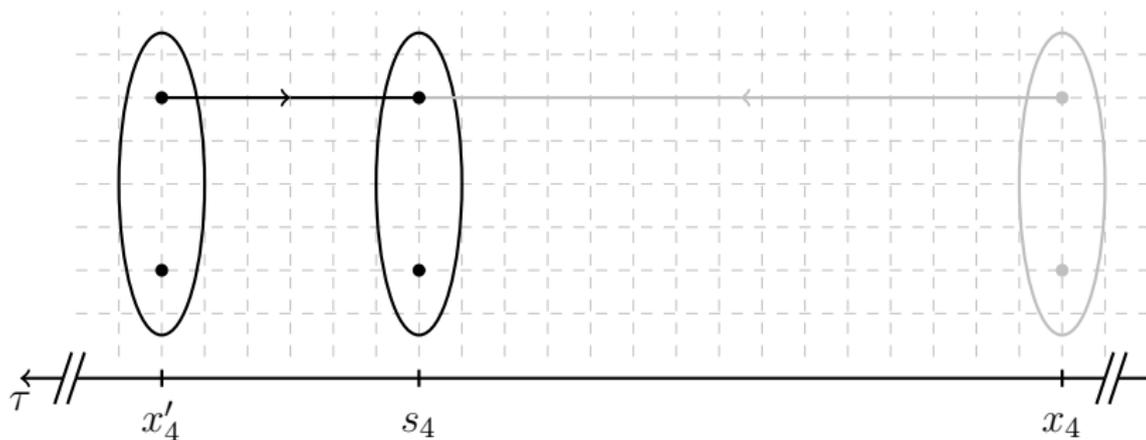
$$\begin{aligned}
 & \#spin \cdot \#color \cdot \#stoch \cdot \#momenta \cdot \#src \\
 & \cdot fwd/bwd \cdot \#flavorcombinations \cdot \#noisemearing = \\
 & 64 \cdot 3 \cdot 100 \cdot 60 \cdot 4 \cdot 2 \cdot 3 \cdot 3 = \\
 & 8.29 \cdot 10^7 \text{ Complex Numbers} \hat{=} 1.30 \text{ GB}
 \end{aligned}$$

Insertion size (upper bound)

Insertion size =

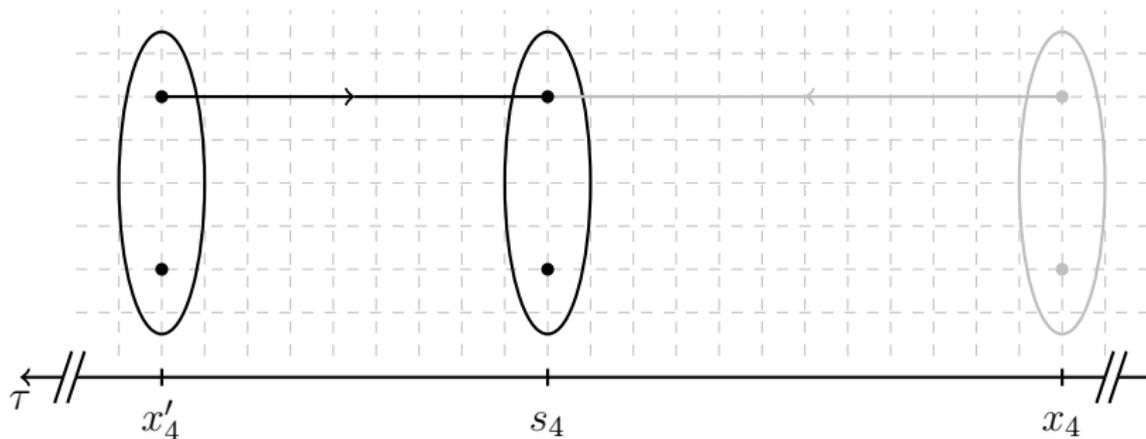
$$\begin{aligned}
 & \#spin \cdot \#color \cdot \#stoch \cdot \#momenta \cdot \#src \cdot \#ins \\
 & \cdot \#derivatives \cdot \#flavorcombinations = \\
 & 64 \cdot 3 \cdot 100 \cdot 60 \cdot 4 \cdot 20 \cdot (1 + 4) \cdot 9 = \\
 & 4.15 \cdot 10^9 \text{ Complex Numbers} \hat{=} 64.8 \text{ GB}
 \end{aligned}$$

Computation of the Spectator part (serial)

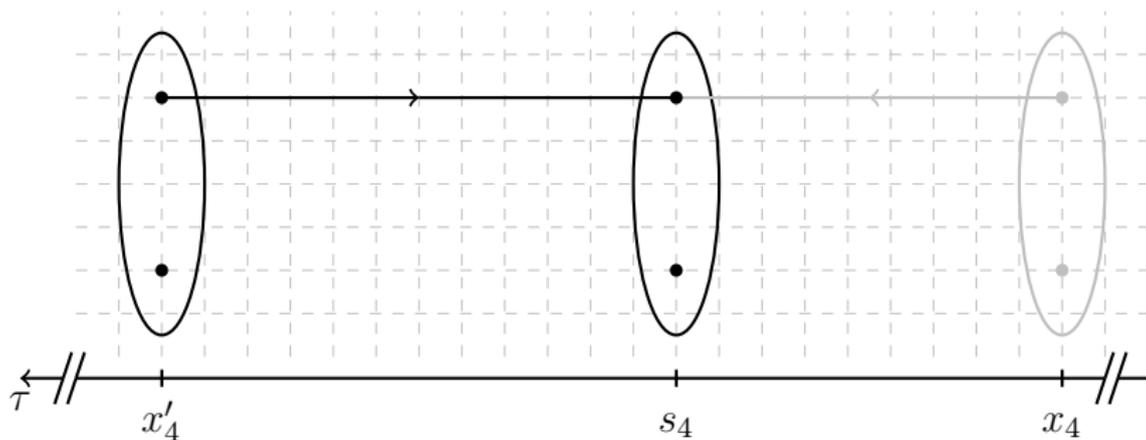
Serial loop over Source time slice s_4 

Computation of the Spectator part (serial)

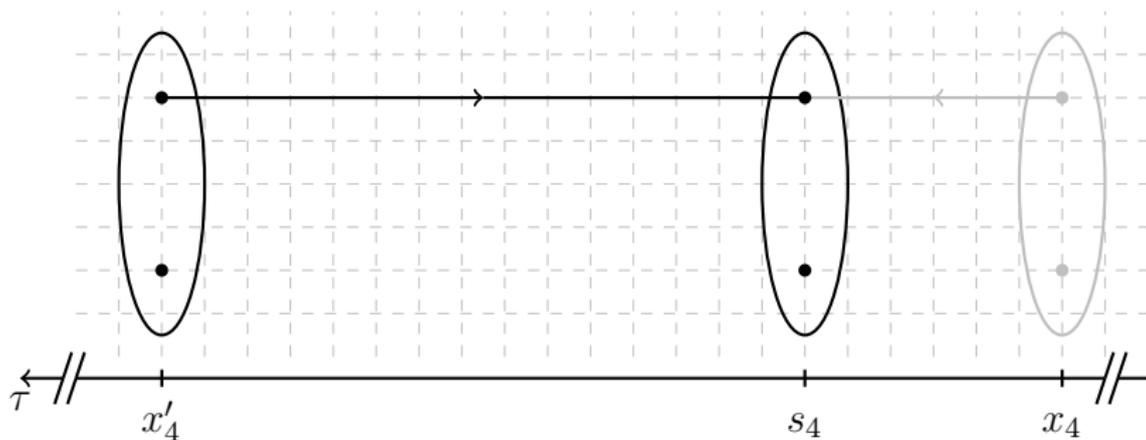
Serial loop over Source time slice s_4



Computation of the Spectator part (serial)

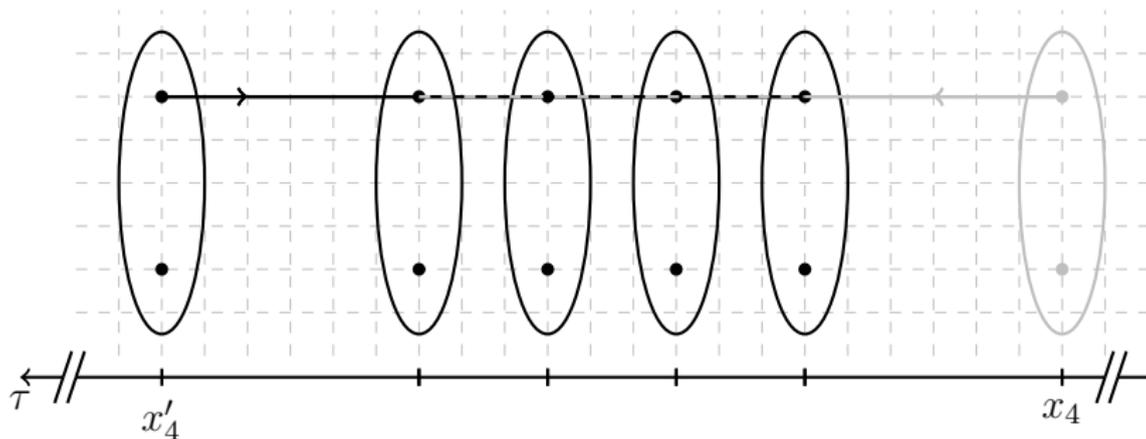
Serial loop over Source time slice s_4 

Computation of the Spectator part (serial)

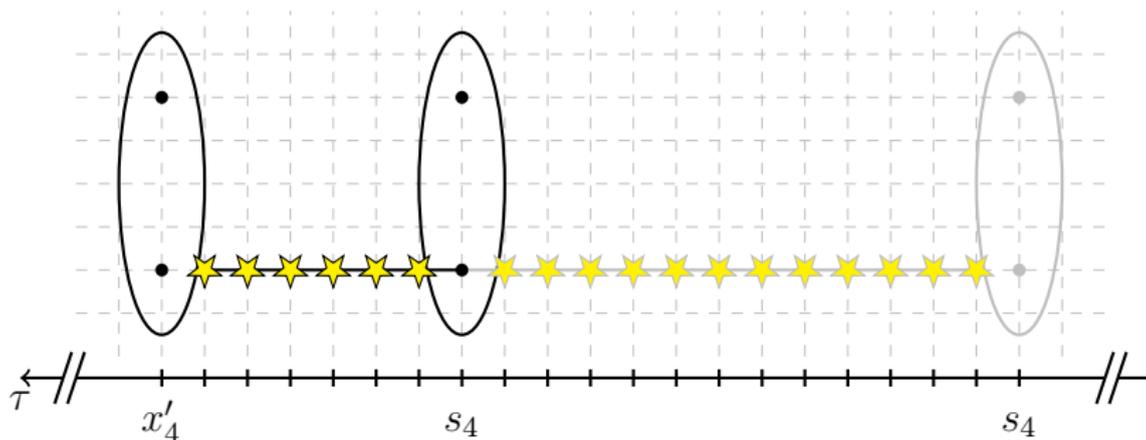
Serial loop over Source time slice s_4 

Computation of the Spectator part (parallel)

Parallel loop over Source time slice s_4



Computation of the Insertion part (parallel)

Serial outer Loop over Source time slice s_4 Parallel inner Loop over Current time slice y_4 

Reference implementation:

QDP++ library takes care of parallelism over lattice sites
 → SIMD vectorization difficult due to its internal data layout.

LibHadronAnalysis:

Reorder loops in order to use two of the spin indices for
 auto-vectorization.

```

for 3 spin indices do
  |
  for color indices a,b do
    |
    for all sites in the local lattice do
      |
      end
    end
  end
end
  
```

Algorithm 1: Reference implementation

```

for all sites in the local lattice do
  |
  for 1 spin index do
    |
    for color indices a,b do
      |
      for 2 spin indices SIMD vect. do
        |
        end
      end
    end
  end
end
  
```

Algorithm 2: LHA implementation

Machine: QPACE 3, Xeon Phi (KNL), Omni-Path, 8 nodes

Lattice Size: $32^3 \times 96$

Spec.: $S_i^{L,S} = \tilde{G}(L, x', s) \tilde{G}(S, x', s) \tilde{\eta}_i(S, x')$ ($0 \leq i < 30$)

Ins.: $I_i^{S,S} = s_i(S, y) G(S, y, s) + \text{deriv.}$ ($0 \leq i < 30$)

Src./Snk.: $s = 30, x' = 16/44$

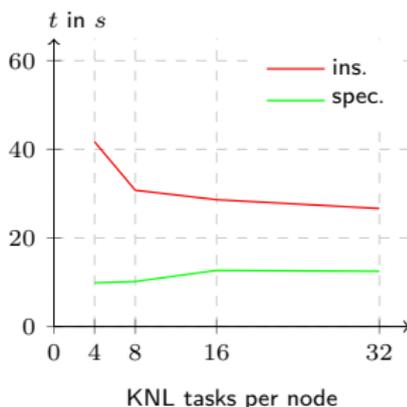


Figure 1: Contraction

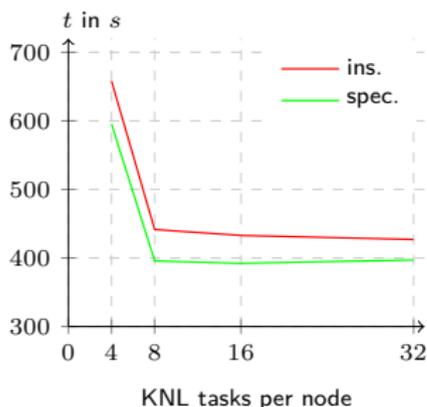


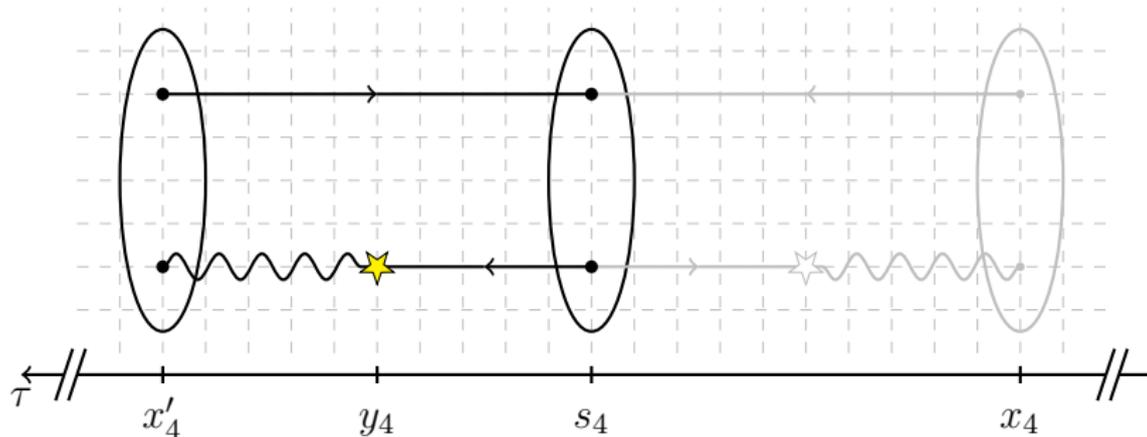
Figure 2: Total time (propagators etc.)

²Peter Georg: An in-depth evaluation of the Intel Omni-Path network for LQCD applications. 19/6/2017 17:00

Computation of mesonic 3pt functions by factorization into spectator and insertion parts

Highly parallelized

Using smart ordering of operator and site index loops \rightarrow SIMD
'Open 3pt Function'. Analysis software is also available for the
off-line contractions.



Thank you for your attention