

# Determination of QCD phase transition line in the canonical approach

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*D.L. Boyda - speaker*

# Mission of this talk

## This talk is about:

- Calculation of observables at nonzero chemical potential using LQCD simulations at imaginary  $\mu$

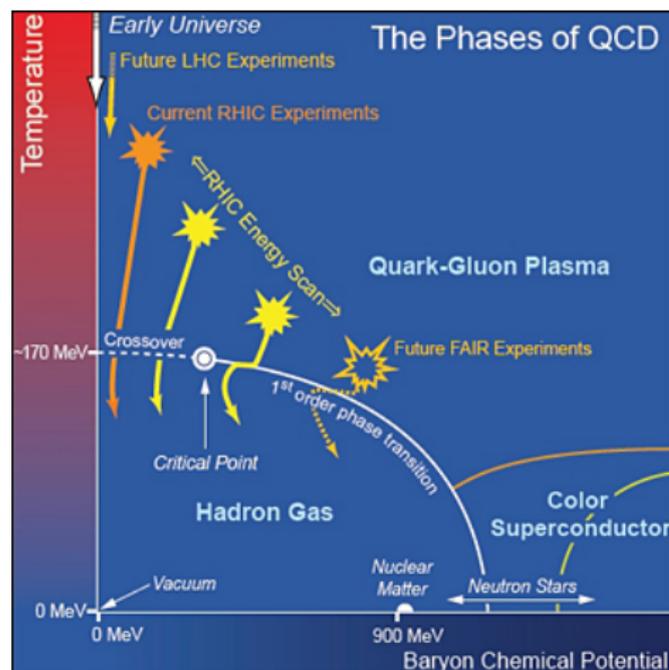
## Main goals:

- Reliability range of our results for real chemical potential values
- Propagation of errors from imaginary  $\mu$  calculations to real  $\mu$  values
- (*Preliminary*) crossover line

# The outline

- Introduction: QCD Phase Diagram, Sign Problem, Canonical Approach, Number Density Formulation
- Lattice Setup
- Number density fits and extrapolation to real  $\mu$
- Canonical Approach: Integration method
- Comparison of our results with RHIC data
- Temperature dependence
- Crossover line
- Conclusion

# Introductory (1): QCD Phase Diagram



**Experiments** on heavy ion collisions:  
RHIC, LHC, J-PARC, NICA

**Theory**  
LQCD is only tool based on first principles calculation to study QCD Phase Diagram

but ...

## Introductory (2): Sign problem

Standard Monte Carlo techniques

$$Z_{GC}(\mu, T, V) = \int \mathcal{D}U (\det \Delta(\mu))^{N_f} e^{-S_G}$$

Fermion determinant satisfies relation

$$[\det \Delta(\mu)]^* = \det \Delta(-\mu^*)$$

- $\mu$  is real  $\rightarrow \det \Delta(\mu)$  is complex
- $\mu$  is imaginary  $\rightarrow \det \Delta()$  is real

!!!

**Monte Carlo calculations for imaginary  $\mu$  is free from Sign problem**

But how can we extract data for real  $\mu$ ?

## Introductory (3): Canonical approach

By definition

$$\begin{aligned} Z_{GC}(\mu, T, V) &= \text{Tr}(e^{-\frac{\hat{H}-\mu\hat{N}}{T}}) = \sum_{n=-\infty}^{\infty} \langle n | e^{-\frac{\hat{H}}{T}} | n \rangle e^{\frac{\mu n}{T}} = \\ &= \sum_{n=-\infty}^{\infty} Z_C(n, T, V) e^{\frac{\mu n}{T}} = \sum_{n=-\infty}^{\infty} Z_n \xi^n \end{aligned}$$

$Z_C(n, T, V)$  - canonical partition function ( $Z_n$  in following)

$\xi = e^{\mu/T}$  - fugacity

$\hat{N}$  - operator of any conserved quantum number (baryon, charge, etc.)

For imaginary  $\mu$  we can calculate  $Z_n$  by inverse Fourier transformation (A. Hasenfratz and D. Toussaint, Nucl. Phys. B 371 (1992))

$$Z_n = \int_0^{2\pi} \frac{d\theta}{2\pi} e^{-in\theta} Z_{GC}(\mu = i\theta T, T, V).$$

## Introductory (4): Number Density Formulation

Number density on a **lattice**:

$$n_B^{lat} = \frac{1}{3N_s N_t} \frac{\partial \ln Z_{GC}}{\partial (a\mu)} = \frac{N_f}{3N_s^3 N_t} \int \mathcal{D}U e^{-S_G \text{tr}} \left[ \Delta^{-1} \frac{\partial \Delta}{\partial (a\mu)} \right] (\det \Delta(\mu))^{N_f}$$

Number density using **canonical formula** ( $\frac{\mu}{T} = i\frac{\mu_l}{T} = i\theta$ )

$$\begin{aligned} n_B &= \frac{1}{3V} T \frac{\partial}{\partial \mu} \ln \left( \sum_{n=-\infty}^{\infty} Z_n e^{(\mu/T)n} \right) = C \frac{\sum_{n>0}^{\infty} n Z_n (e^{i\theta n} - e^{-i\theta n})}{Z_0 + \sum_{n>0}^{\infty} Z_n (e^{i\theta n} + e^{-i\theta n})} \\ &= iC \frac{2 \sum_{n>0}^{\infty} n Z_n \sin(n\theta)}{1 + 2 \sum_{n>0}^{\infty} Z_n \cos(n\theta)} \end{aligned}$$

where we used  $Z_n = Z_{-n}$   
(we write  $Z_n/Z_0$  as  $Z_n$ )

# Lattice Setup

- clover improved Wilson action
- Iwasaki gauge action
- Lattice  $4 \times 16^3$  ( $L \approx 3.2fm$ ,  $a \approx 0.2fm$ )
- $\frac{m_\pi}{m_\rho} = 0.8$  ( $m_\pi = 0.7GeV$ )
- $T/T_c = 0.93(5)$  (0.84, 0.99, 1.08, 1.20, 1.35)
- 40 values  $\mu_l$ , 1800 - 3800 configurations (10 trajectories separated)

Parameters were taken from  
S. Ejiri et. al., PRD 82, 014508 (2010)

Our cluster: Vostok1 (20 GPU K40)

# Number density fits and Extrapolation to real $\mu$

**Idea:** number density -  $n_B(\mu/T)$  can be parametrized by some function  $f^{pol}(x)/f^{sin}(x)$  at imagine  $\mu_I$  and extrapolated to real region

**Taylor series** (Polynomial fit)

$$f^{pol}(x) = \sum_k^{N_{max}} a_k x^{2k+1}$$

**Fourier series**

$$f^{sin}(x) = \sum_k^{N_{max}} f_k \sin(k x), \quad x = \mu_I/T$$

**Extrapolation** ( $x \rightarrow ix$ ):

$$\rightarrow \sum_k^{N_{max}} (-1)^k a_k x^{2k+1}$$

$$\rightarrow \sum_k^{N_{max}} f_k \sinh(k x)$$

M. D'Elia, et. al., Phys. Rev. D 95, 094503 (2017)

J. Gunther, et. al., arXiv:1607.02493 (2016)

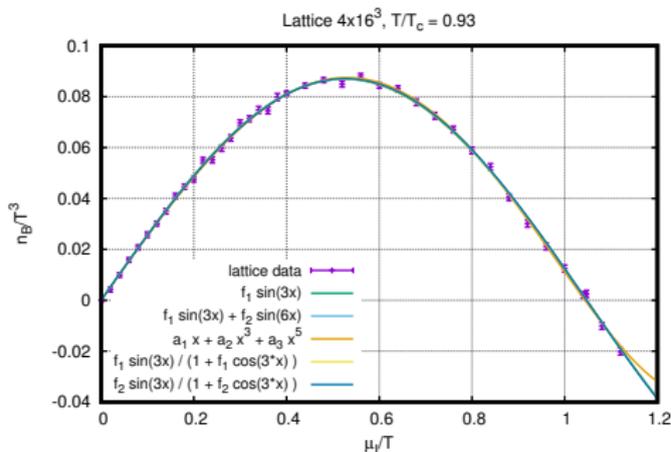
J. Takahashi et. al., Phys. Rev. D 91, 014501 (2015)

J. Takahashi et. al., arXiv:1002.0890 (2010)

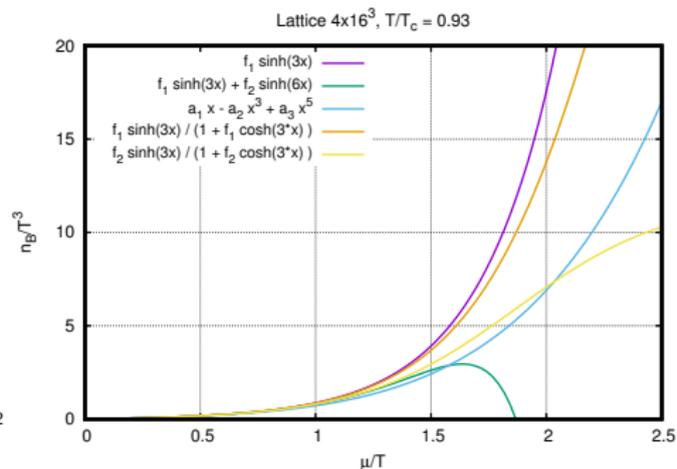
M. D'Elia, et. al., Phys. Rev. D 80, 014502 (2009)

# Number density fits and Extrapolation to real $\mu$ (2)

## Fit Imagine $\mu$



## Extrapolation to Real $\mu$



### Questions:

- How can we estimate **reliability range** of  $\mu$ ?
- Is this method able to **find transition line** ?

# Canonical approach: Integration method

**Idea:** Using observable - baryon number density calculate  
Grand Canonical Partition Function -  $Z_{GC}(\mu/T)$   
(V. Bornyakov, et. al. Phys. Rev. D 95, 094506 (2017))

working at pure imaginary chemical potential  $\mu = i\mu_I$ :

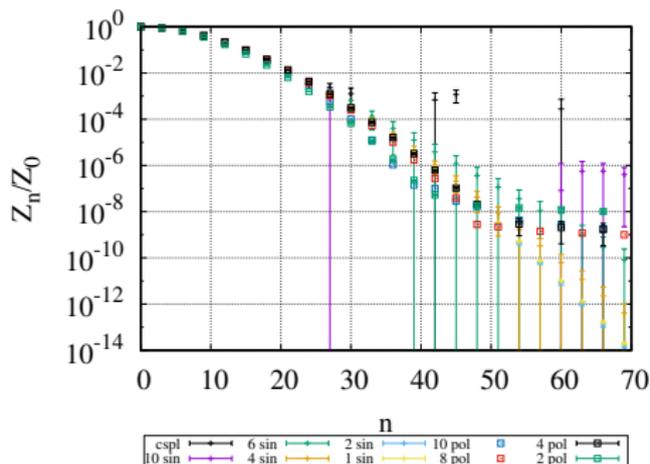
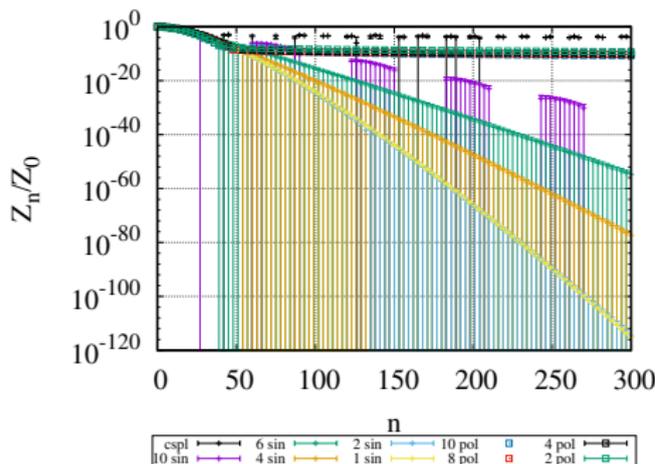
$$n_B = \frac{1}{V} \frac{\partial(\ln Z_G)}{\partial(\mu/T)} \quad \rightarrow \quad \ln Z_{GC}(\theta) - \ln Z_{GC}(0) = V \int_0^\theta d(i\tilde{\theta}) i \operatorname{Im}[n_B(\tilde{\theta})]$$
$$\frac{Z_{GC}(\theta)}{Z_{GC}(0)} = \exp\left(-V \int_0^\theta dx \operatorname{Im}[n_B(x)]\right)$$

- numerical integration (quadrature formula/cubic spline)
- use of some parametrization ( $f^{pol}(x)$  or  $f^{sin}(x)$ ) for  $n_B(x)$

Strategy:

- 1 Calculate  $n_B(\mu/T)$  using LQCD simulation at imaginary  $\mu$
- 2 Calculate  $Z_{GC}$
- 3 Calculate  $Z_n$  as Fourier transformation of  $Z_{GC}$
- 4 Using  $Z_n$  calculate any thermodynamic observable

# Canonical approach: Integration method (2)

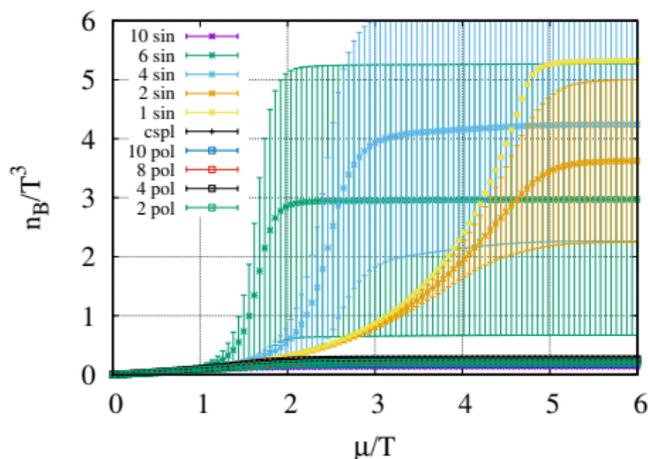
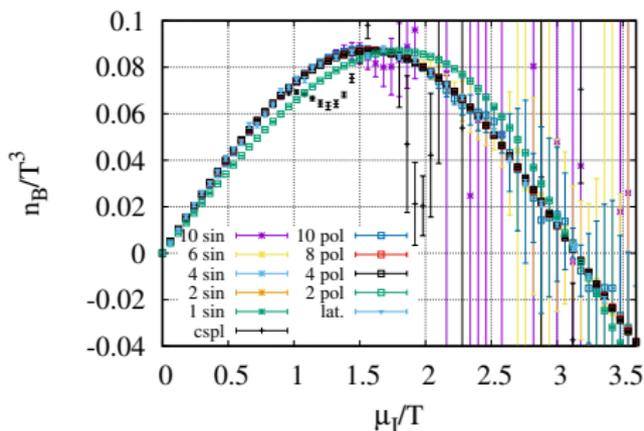


$Z_n$  must be positive!

$$n_B = \sum_k^{N_{max}} \sin(kx)$$

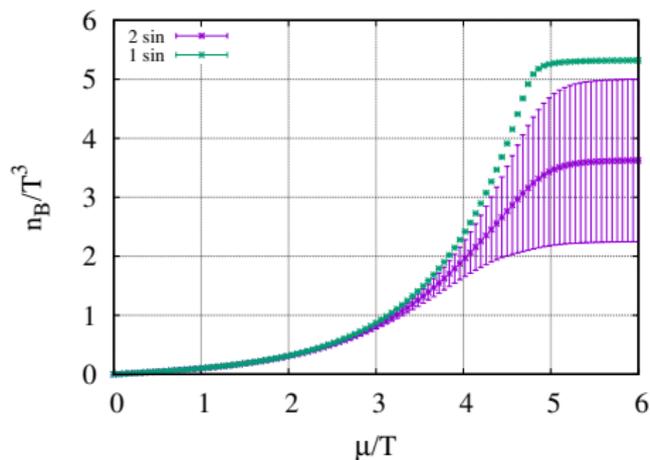
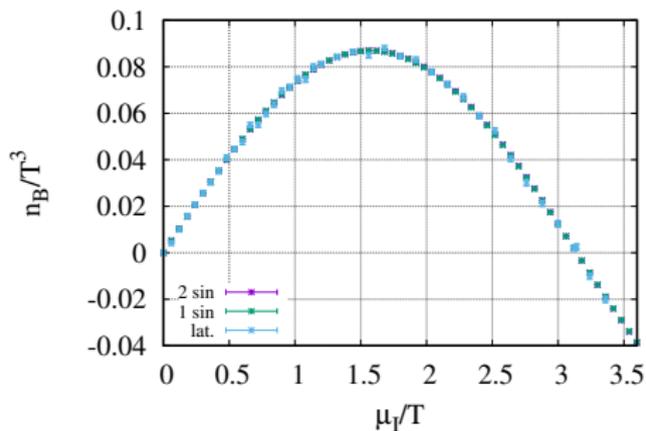
$N_{max}$	$f_1 * 10^6$	$f_2 * 10^6$	$f_3 * 10^6$	$f_4 * 10^6$	$\chi^2/d.o.f.$
1	1359(4)	-	-	-	0.94
2	1360(4)	-5(4)	-	-	0.93
4	1360(4)	-4(4)	-3(4)	-2(4)	0.96

# Canonical approach: Integration method (3)



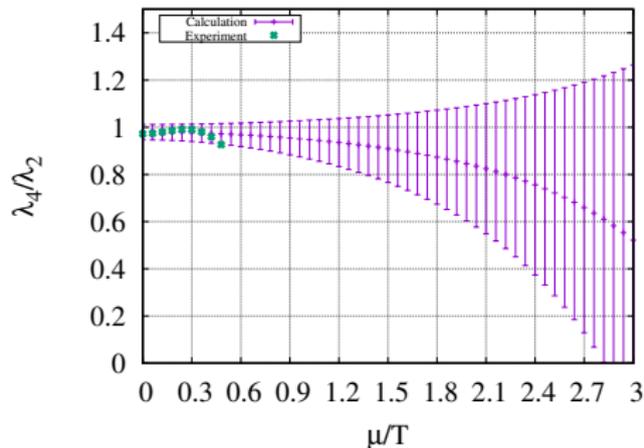
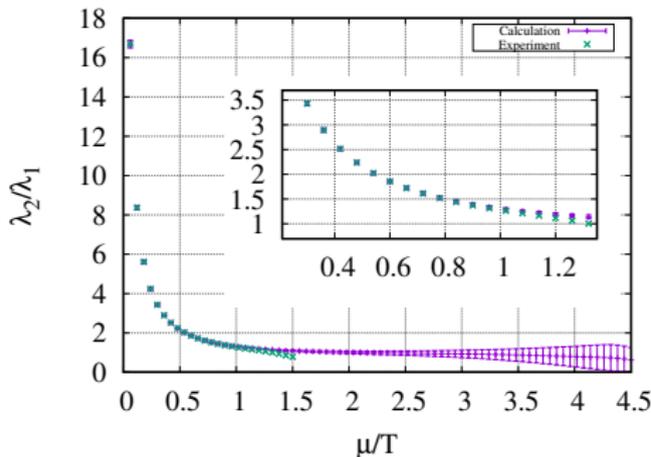
$$n_B = C \frac{2 \sum_n^{N_{max}} n Z_n \sinh(n\theta)}{1 + 2 \sum_n^{N_{max}} Z_n \cosh(n\theta)} \xrightarrow{\theta \rightarrow \infty} C \frac{2 N_{max} Z_{N_{max}} \sinh(N_{max}\theta)}{2 Z_{N_{max}} \cosh(N_{max}\theta)} \sim C N_{max}$$

# Canonical approach: Integration method (4)



# Comparison of our results with RHIC data

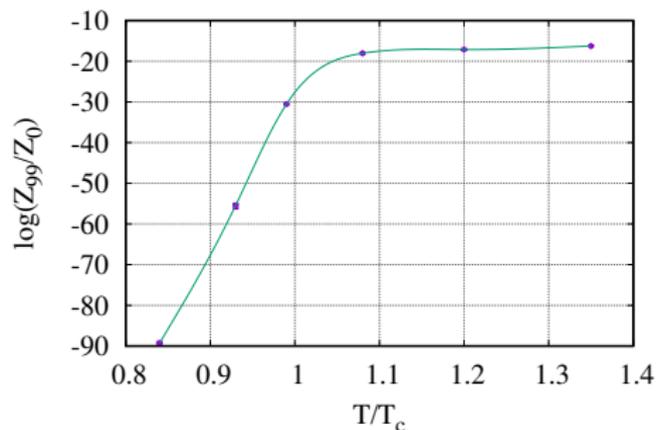
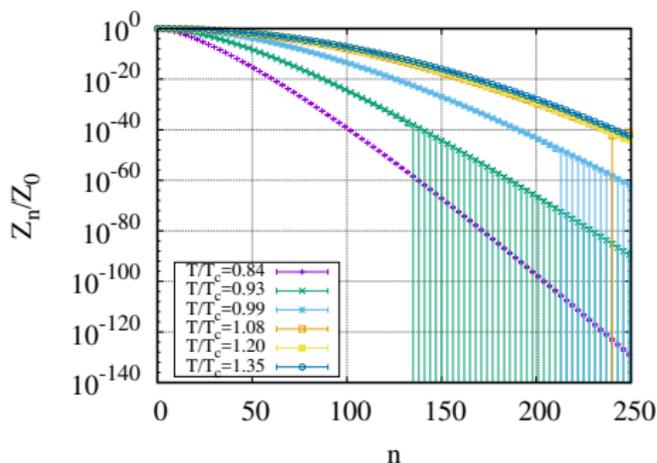
$$\lambda_k(\mu) = \left( T \frac{\partial}{\partial \mu} \right)^k \log \left( \sum_n Z_n \xi^n \right)$$



$Z_n$  constructed for proton multiplicity RHIC data ( $\sqrt{s_{NN}} = 62.4$  GeV) were taken from A. Nakamura, K. Nagata, *PTEP* 3, 33D01 (2016)

# Temperature dependence

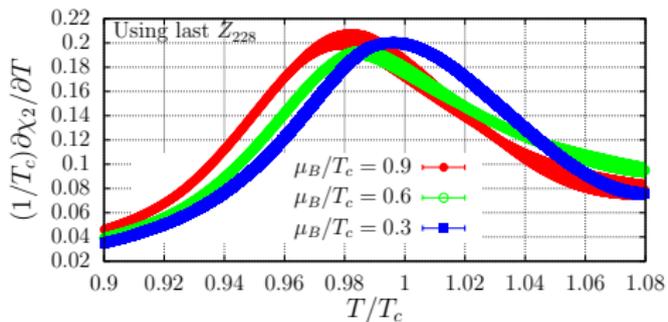
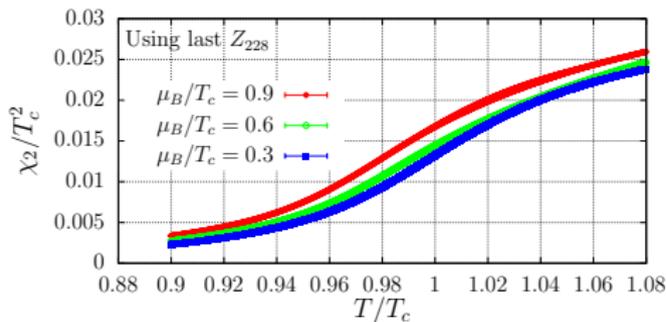
Using  $Z_n$  dependence on temperature we can calculate observables for any temperature and study crossover(transition) line



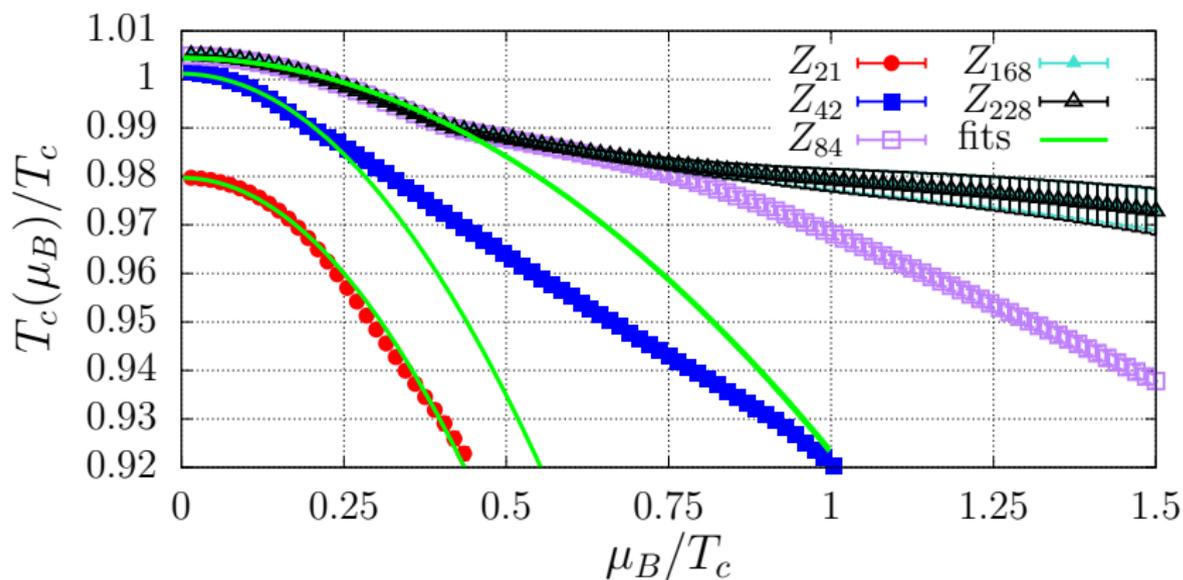
See talk by Vladimir Goy → next 2 slides

# (Preliminary) Dependence of susceptibility on T

$$\chi_2(\mu, T) = \frac{\partial}{\partial \mu} n_B(\mu, T) = \frac{T}{V} \left( \frac{\partial}{\partial \mu} \right)^2 \log \left( \sum_n^{N_{max}} Z_n(T) \xi^n \right), \quad \xi = e^{\mu/T}$$



## (Preliminary) Crossover line



for small  $\mu$  crossover line can be fitted by  $f(x) = a - bx^2$ ,  $a = 1.0049(3)$ ,  $b = 0.033(10)$

# Conclusions

- Canonical approach and analysis with  $Z_n$  allow to choose the best parametrization function and find reliability range
- $\lambda_2/\lambda_1$  and  $\lambda_4/\lambda_2$  were calculated and agreed with RHIC data
- Using  $Z_n$  dependence on temperature we can study any temperature
- (*Preliminary*) Crossover line was calculated
- Integration method allows study QCD at finite density without any ansatz (using numerical integration for calculation  $Z_{GC}$ )