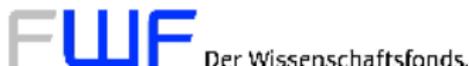


# Dual simulation of the massless lattice Schwinger model with a topological term and non-zero chemical potential

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22.06.2017



Partition sum of the massless one-flavor model with staggered fermions and a topological term in conventional representation:

$$Z = \int \mathcal{D}[U] \int \mathcal{D}[\bar{\psi}, \psi] e^{-S_G[U] - i\theta Q[U] - S_\psi[U, \bar{\psi}, \psi]},$$

$S_G[U] \dots$

Wilson Gauge action

$Q[U] \dots$

Topological charge

$S_\psi[U, \bar{\psi}, \psi] \dots$

Staggered Fermion action

→ Sign problem for  $\theta \neq 0$  or  $\mu \neq 0$ !

# The partition sum in dual representation

Dual partition sum for one flavor and  $\theta > 0$ :

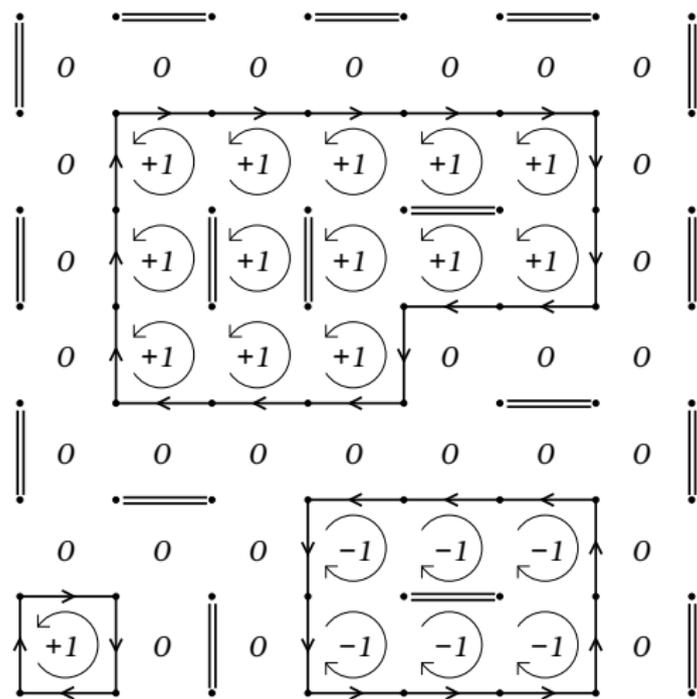
$$Z = \left(\frac{1}{2}\right)^V \sum_{\{l,d,p\}} \prod_n \mathbb{I}_{|p(n)|} (2\sqrt{\eta\bar{\eta}}) \left(\sqrt{\frac{\eta}{\bar{\eta}}}\right)^{p(n)}$$

$$\eta = \frac{\beta}{2} - \frac{\theta}{4\pi}, \quad \bar{\eta} = \frac{\beta}{2} + \frac{\theta}{4\pi}$$

$$l_{n,\nu} \in \{-1, 0, 1\}, \quad d_{n,\nu} \in \{0, 1\}, \quad p_n \in (-\infty, \infty)$$

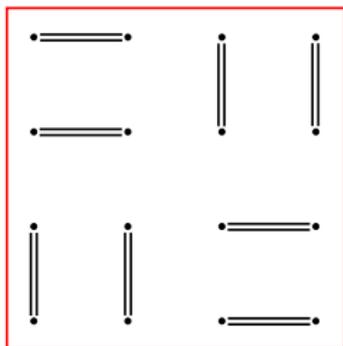
*C. Gattringer, T. Kloiber, V. Sazonov, Solving the sign problems of the massless lattice Schwinger model with a dual formulation, Nucl.Phys. B897 (2015) 732-748 arXiv:1502.05479 [hep-lat]*

# Example of an admissible configuration

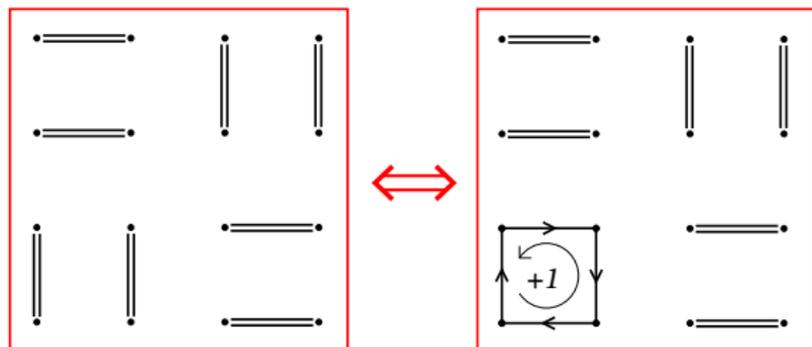


- **Pauli principle:**  
Each lattice point has to be occupied by a fermion exactly once.
- **Vanishing link flux:**  
 $\Rightarrow$  Sheets of plaquettes bounded by loops.

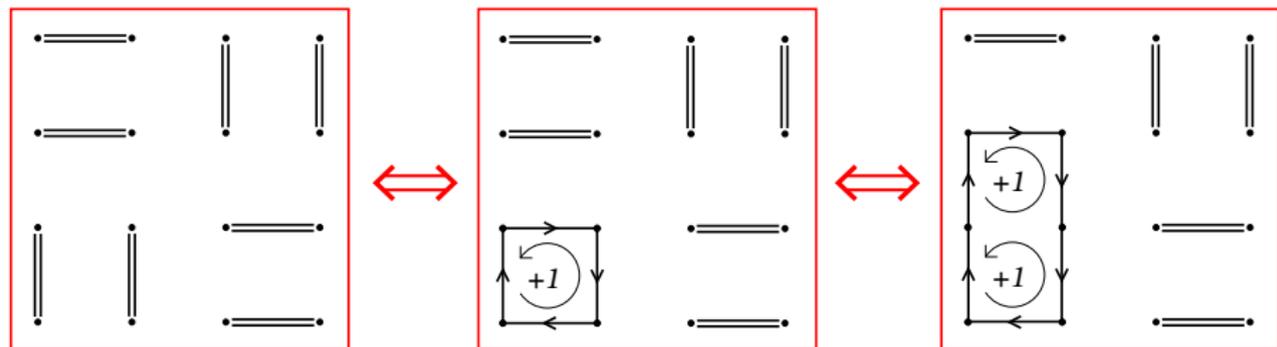
# Local updates - accepted with Metropolis steps



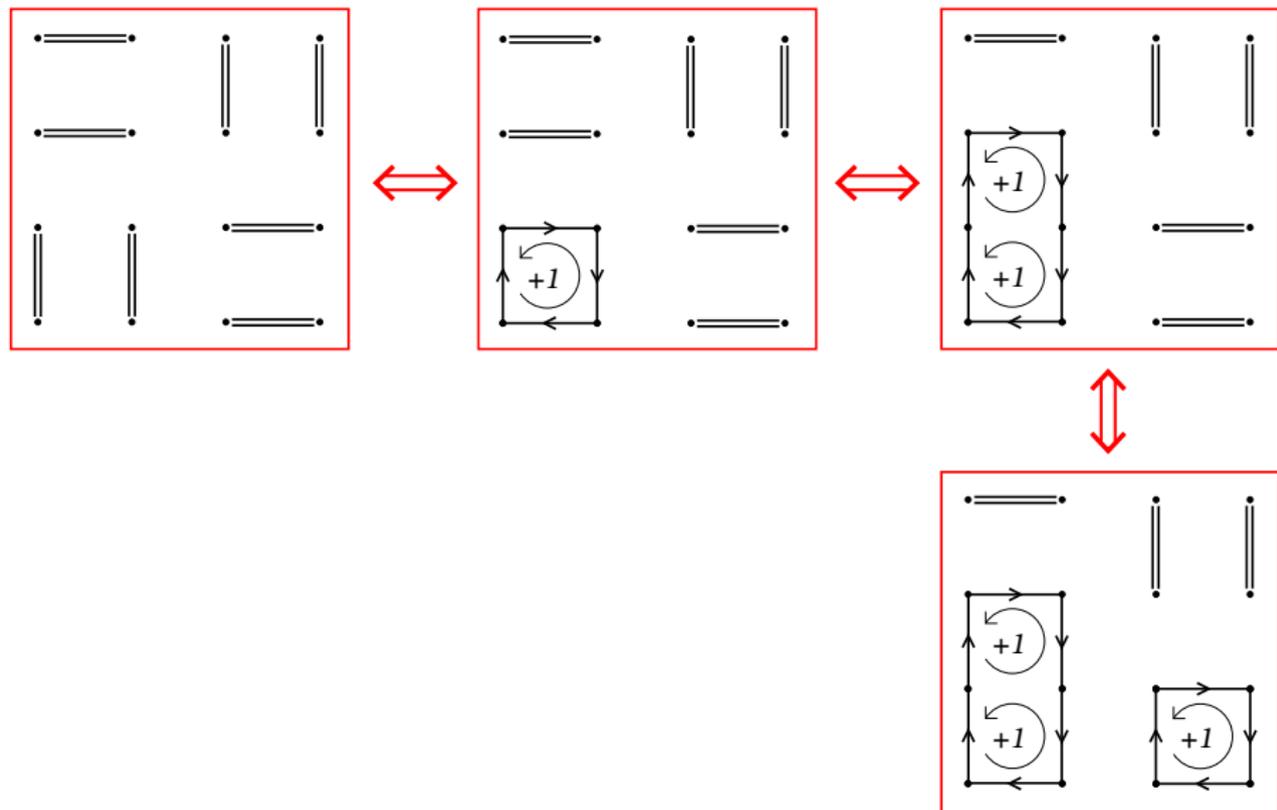
# Local updates - accepted with Metropolis steps



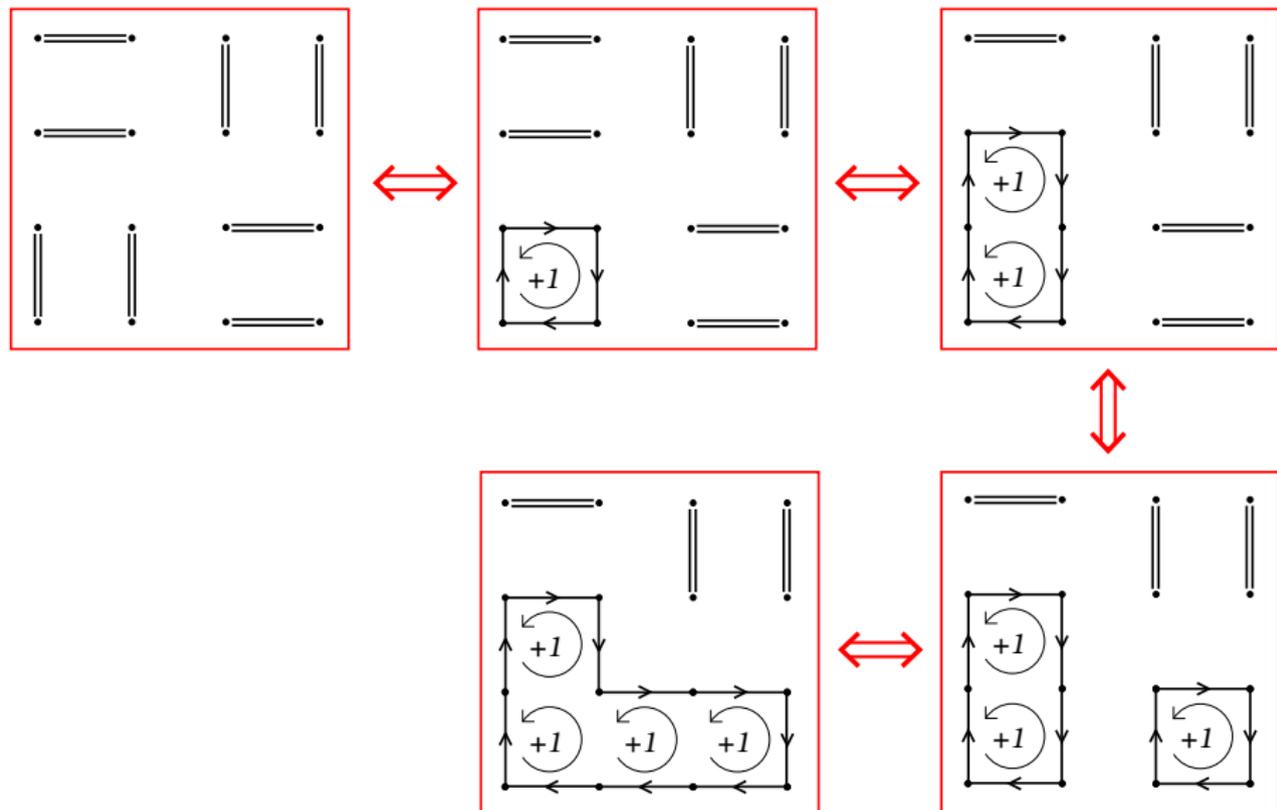
# Local updates - accepted with Metropolis steps



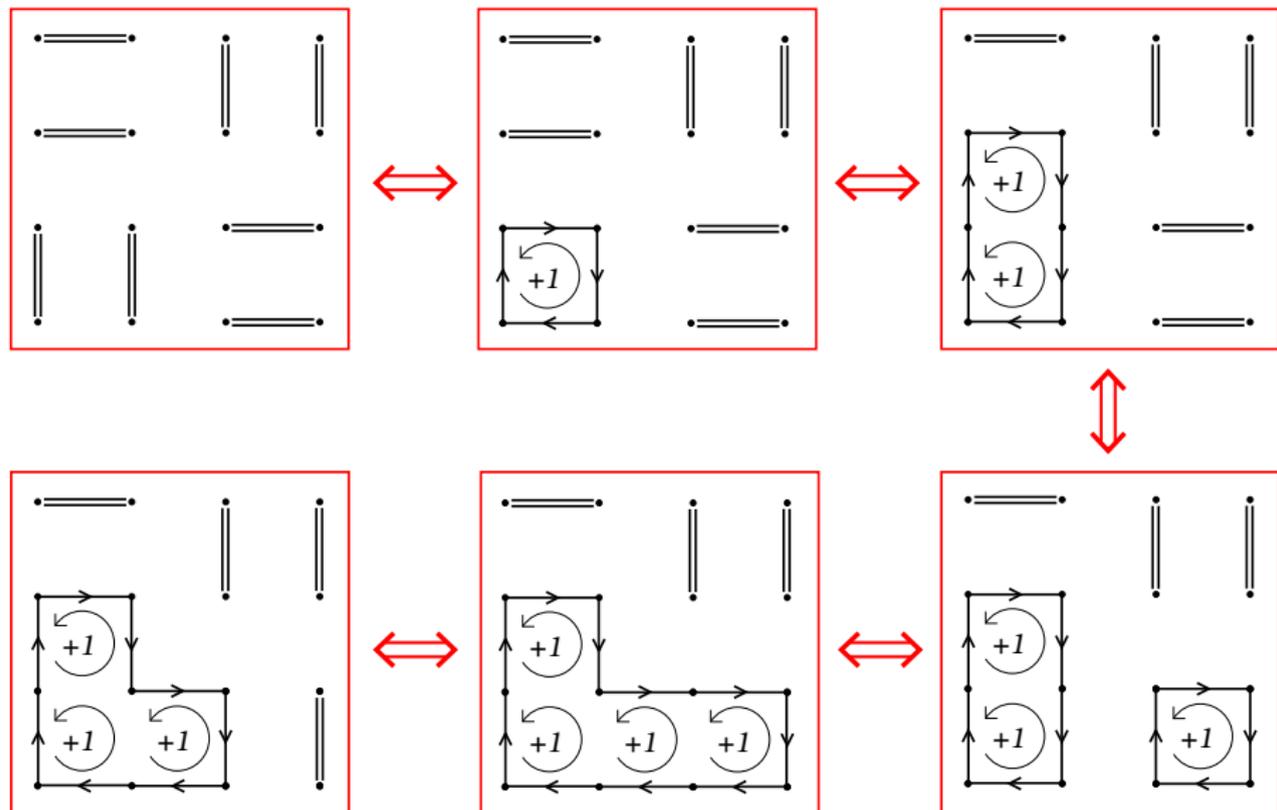
# Local updates - accepted with Metropolis steps



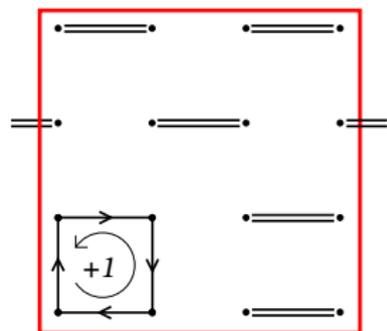
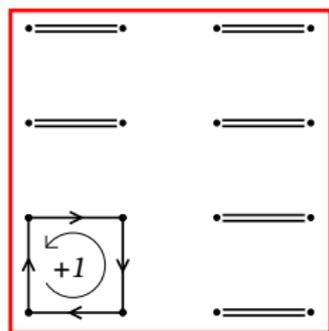
# Local updates - accepted with Metropolis steps



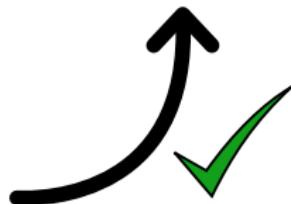
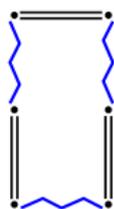
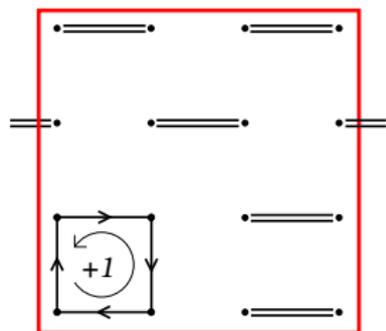
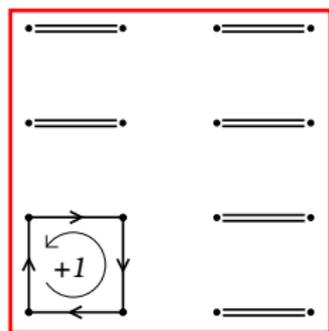
# Local updates - accepted with Metropolis steps



# Dimers



# Worm update for dimers



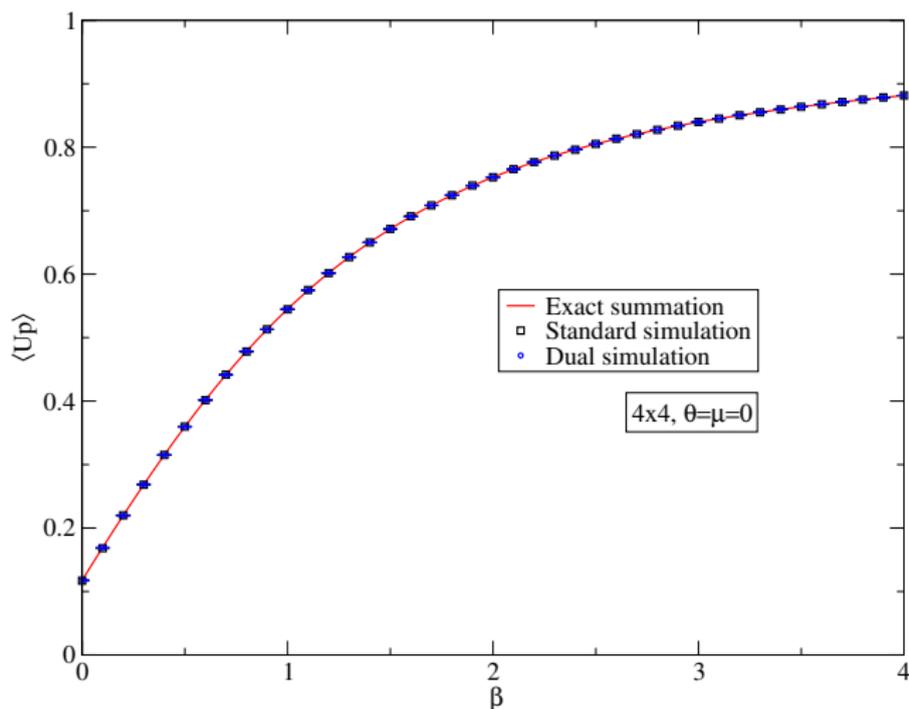
- Plaquette expectation value and susceptibility:

$$\langle U_p \rangle = \frac{1}{V} \frac{\partial}{\partial \beta} \ln(Z)$$
$$\chi_p = \frac{1}{V} \frac{\partial^2}{\partial \beta^2} \ln(Z)$$

- Topological charge density and susceptibility:

$$\langle q \rangle = - \frac{1}{V} \frac{\partial}{\partial \theta} \ln(Z)$$
$$\chi_q = - \frac{1}{V} \frac{\partial^2}{\partial \theta^2} \ln(Z)$$

# Testing the algorithm



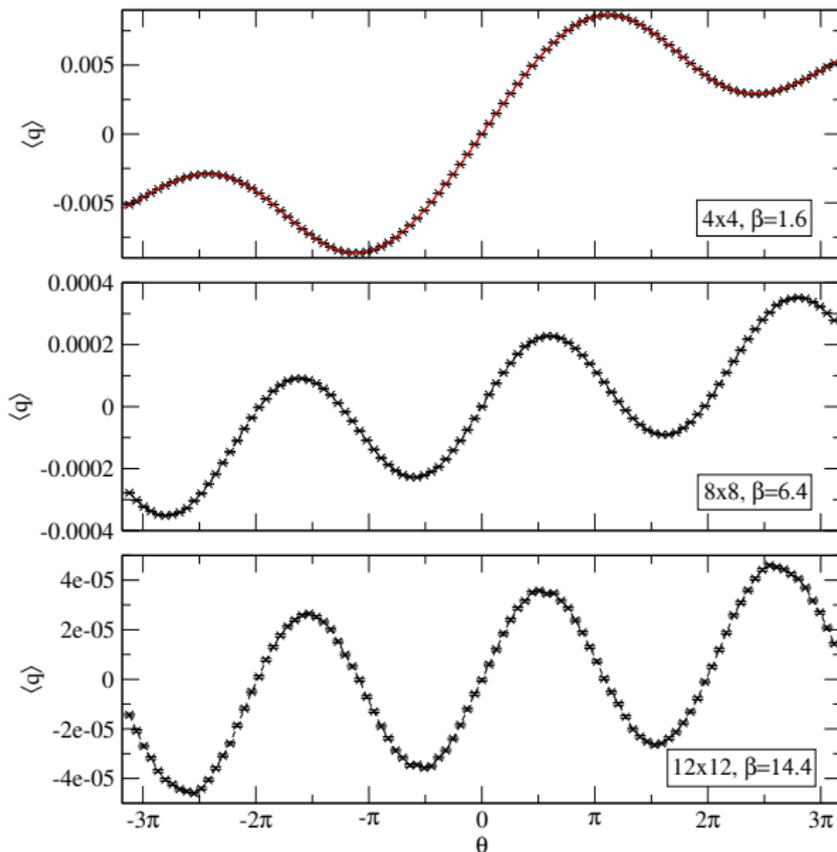
Testing the algorithm with standard simulation at  $\theta = 0$  and exact summation on a small lattice.

# Emerging $2\pi$ periodicity

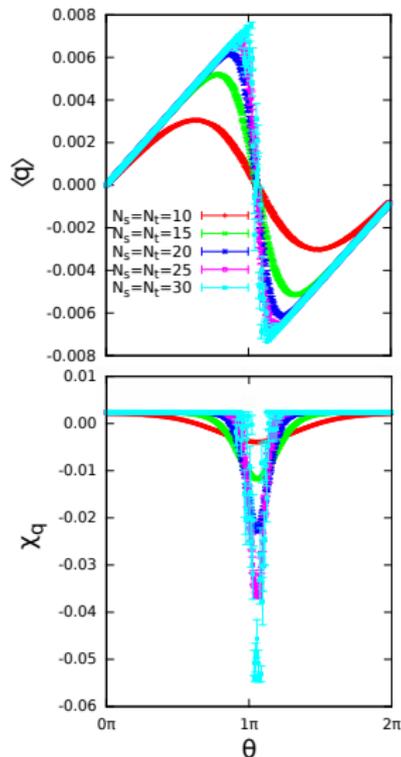
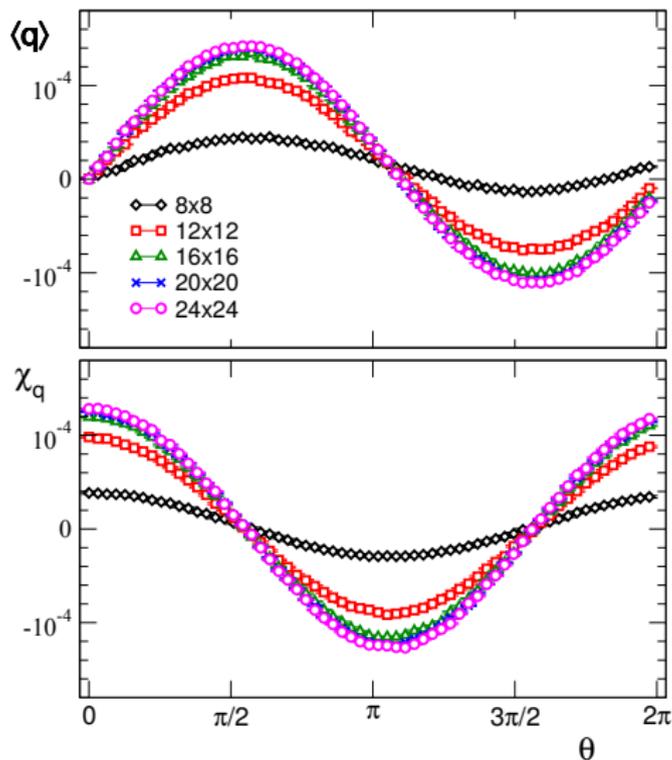
**Topological charge density:**

Approach to the continuum limit with  $\frac{\beta}{N_s N_t} = 0.1 = \text{const}$

Equivalent behavior for  $\langle U_p \rangle$ .



# Behavior near $\theta = \pi$



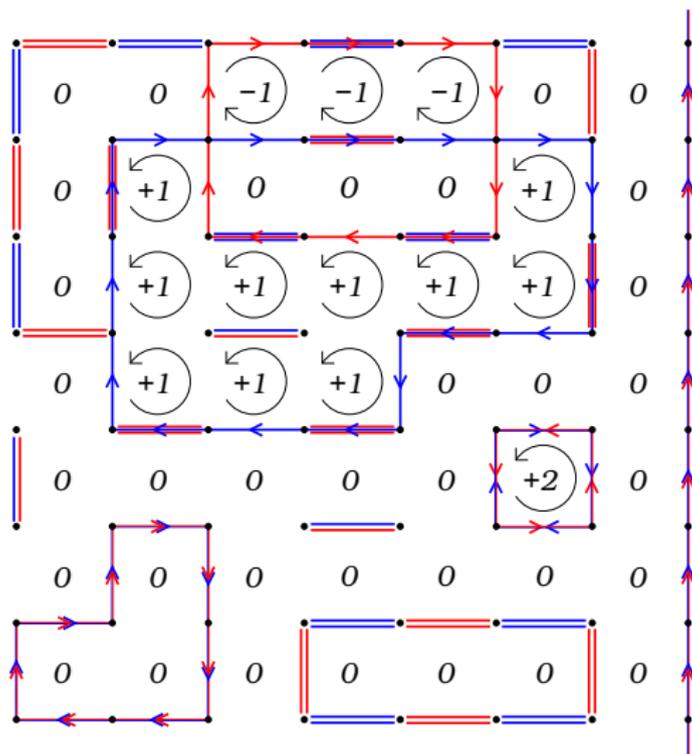
Comparison to scalar QED<sub>2</sub>: [C. Gattringer et al.: Phys. Rev. D 92 \(2015\)](#)

## Two flavor model with chemical potential

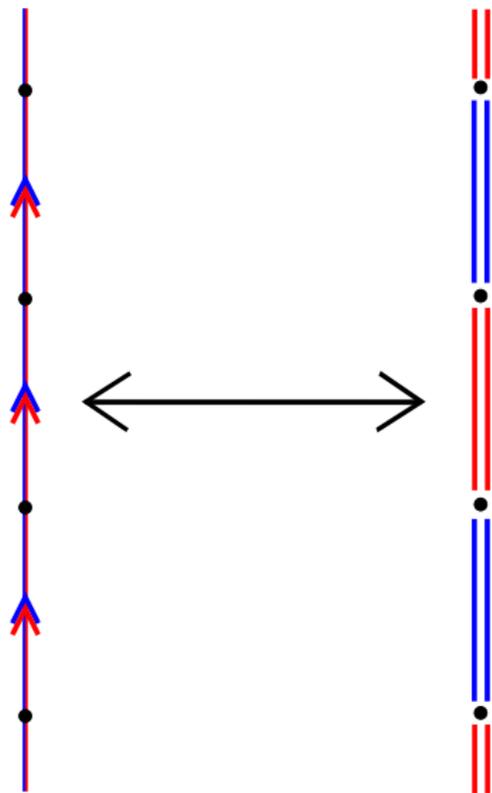
$$Z = \left(\frac{1}{2}\right)^{2V} \sum_{\{l, d, \bar{l}, \bar{d}, p\}} e^{\mu_\psi N_T W(l)} e^{\mu_\chi N_T W(\bar{l})} \prod_n \mathbb{I}_{|p(n)|} (2\sqrt{\eta\bar{\eta}}) \left(\sqrt{\frac{\eta}{\bar{\eta}}}\right)^{p(n)}$$

$$l_{n,\nu}, \bar{l}_{n,\nu} \in \{-1, 0, 1\}, \quad d_{n,\nu}, \bar{d}_{n,\nu} \in \{0, 1\}, \quad p_n \in (-\infty, \infty)$$

# Example of an admissible configuration

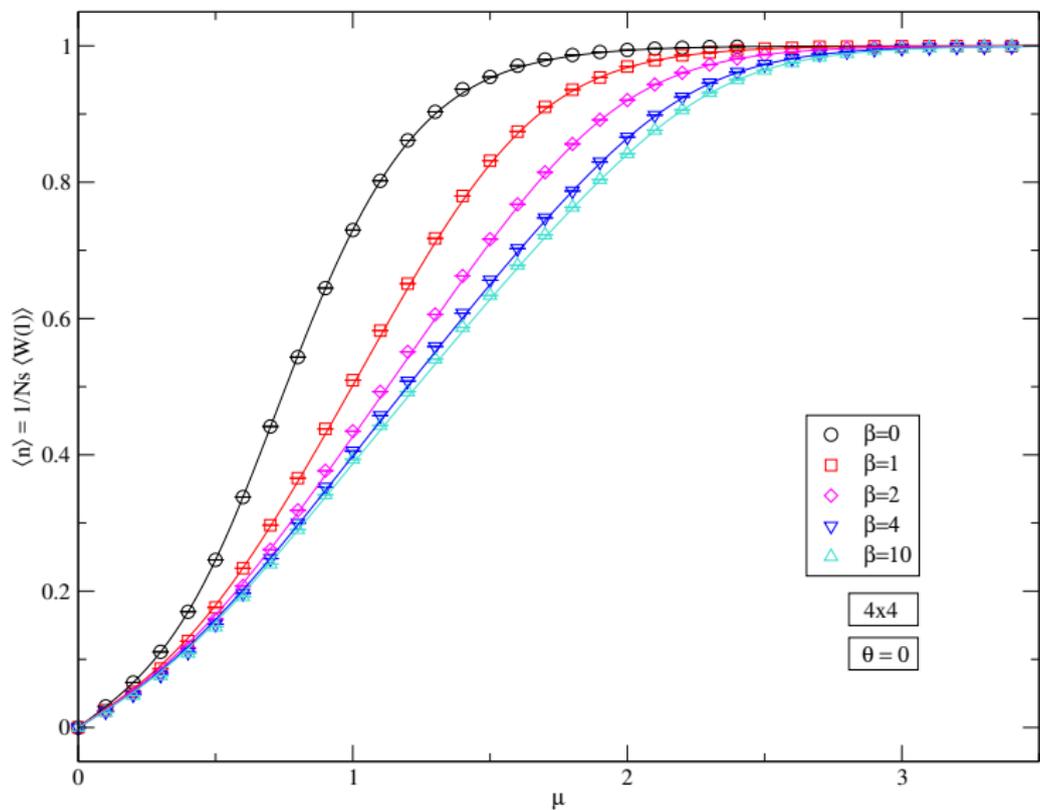


## Additional updates for the winding numbers



- The weight only changes for temporally winding structures.
- Challenge: topological nature of states at large  $\mu$  makes the system very stiff.

# Testing the updates against exact summation



# Summary

- First example of a dual simulation of a lattice field theory with gauge fields and fermions at arbitrary couplings, finite chemical potential and non-zero vacuum angle.
- Fermionic systems pose a challenge due to additional constraints from the Pauli principle.
- Mix of local and worm updates works very well for  $\theta > 0$ .
- Remaining  $\theta$  dependence due to staggered fermions.
- For large  $\mu$ , two-dimensional worldline models can be quite stiff due to topologically stabilized configurations.

Thank you for your attention! :-)

Massless quantum electrodynamics in 2D with a topological term:

$$S[A, \bar{\psi}, \psi] = S_G[A] + S_\psi[A, \bar{\psi}, \psi] + i\theta Q[A]$$

$$S_G[A] = \frac{\beta}{4} \int d^2x F_{\mu\nu}(x) F_{\mu\nu}(x)$$

$$S_\psi[A, \bar{\psi}, \psi] = \int d^2x \bar{\psi}(x) [\gamma_\mu (\partial_\mu + iA_\mu(x))] \psi(x)$$

$$Q[A] = \frac{1}{4\pi} \int d^2x \epsilon_{\mu\nu} F_{\mu\nu}(x)$$

Partition sum of the massless one-flavor model with staggered fermions and a topological term in conventional representation:

$$\begin{aligned}
 Z &= \int \mathcal{D}[U] \int \mathcal{D}[\bar{\psi}, \psi] e^{-S_G[U] - i\theta Q[U] - S_\psi[U, \bar{\psi}, \psi]}, \\
 S_\psi[U, \bar{\psi}, \psi] &= \frac{1}{2} \sum_{n, \nu} [\gamma_\nu(n) U_\nu(n) \bar{\psi}(n) \psi(n + \hat{\nu}) \\
 &\quad - \gamma_\nu(n) U_\nu(n - \hat{\nu})^{-1} \bar{\psi}(n) \psi(n - \hat{\nu})], \\
 S_G[U] &= \beta \sum_n \text{Re}[1 - U_p(n)], \\
 Q[U] &= \frac{1}{i4\pi} \sum_n [U_p(n) - U_p(n)^{-1}], \\
 \gamma_1(n) &= 1, \quad \gamma_2(n) = (-1)^{n_1}, \\
 U_p(n) &= U_1(n) U_2(n + \hat{1}) U_1(n + \hat{2})^{-1} U_2(n)^{-1}.
 \end{aligned}$$

# Expansion in $\beta$

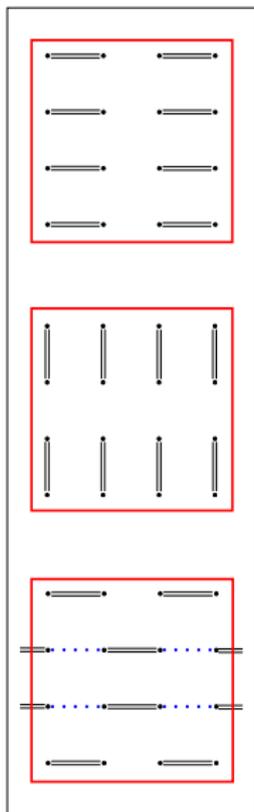
For the case of a vanishing vacuum angle the partition sum in dual representation becomes

$$Z = \sum_{\{l,d,p\}} \prod_n \mathbb{I}_{|\rho(n)|}(\beta)$$

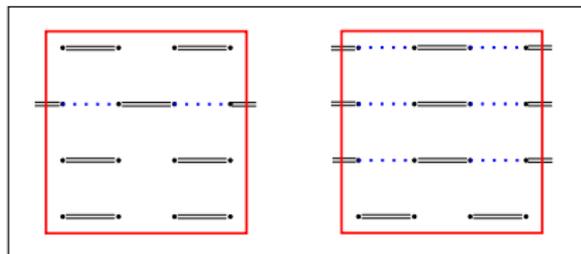
$\Rightarrow$  We expand in  $\beta \hat{=}$  expansion in the number of plaquettes

# Worm updates

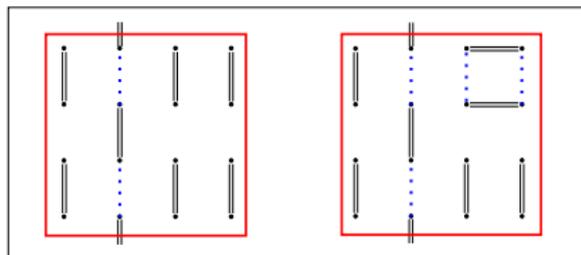
(e,e):



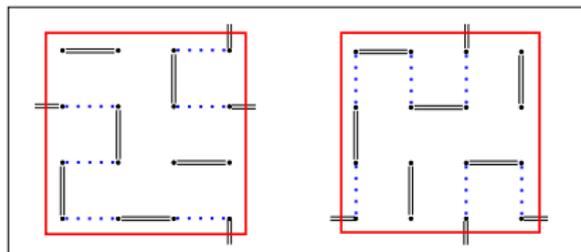
(o,e):



(e,o):



(o,o):



# The worm updates

⇒ We can test dimer (worm) algorithm by explicitly counting the number of configurations in the respective classes:

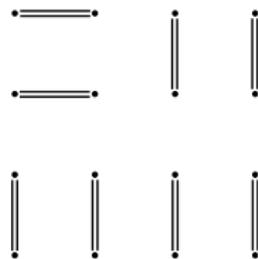
<i>Class</i>	<i>Number of configurations</i>	<i>Analytical ratio</i>	<i>Worm update</i>
(e,e)	136	0.5	$0.49979 \pm 0.000498781$
(e,o)	64	0.235294	$0.235595 \pm 0.000426918$
(o,e)	64	0.235294	$0.235127 \pm 0.00043796$
(o,o)	8	0.0294118	$0.029488 \pm 0.000171351$

(a) Comparison of the ratios on a  $4 \times 4$  lattice.

<i>Class</i>	<i>Number of configurations</i>	<i>Analytical ratio</i>	<i>Worm update</i>
(e,e)	155926656	0.5	$0.50041 \pm 0.000566953$
(e,o)	67108864	0.215194	$0.214472 \pm 0.000467769$
(o,e)	67108864	0.215194	$0.215176 \pm 0.000483039$
(o,o)	21708928	0.0696126	$0.069942 \pm 0.000280876$

(c) Comparison of the ratios on a  $8 \times 8$  lattice.

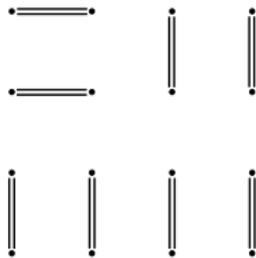
Zero-th order in  $\beta$ :  
272 dimer configurations  
 $\rightarrow 272 I_0(\beta)^{16}$



Zero-th order in  $\beta$ :

272 dimer configurations

$\rightarrow 272 I_0(\beta)^{16}$



First order in  $\beta$ :

32 dimer configurations

16 placements of the plaquette

2 orientations of the loop

$\rightarrow 1024 I_0(\beta)^{15} I_1(\beta)$

