

Update on a Staggered Multigrid Algorithm (with hope in 4D dimensions)

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Motivation

- Push to exascale enables increasingly accurate lattice calculations.
- Physical pion mass, finer lattices: critical slowing down.
 - ▶ MILC: $144^3 \times 288$, physical pion mass, single precision multi-mass solve to rel. resid. 10^{-6} : about 25,000 iterations.
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Multigrid methods can completely eliminate critical slowing down.

MG for the Wilson-Clover operator has a rich history:

[Phys.Rev.Lett. 105 (2010): R. Babich, J. Brannick, R.C. Brower, M.A. Clark, T.A. Manteuffel,
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And for Twisted Mass:

- On \mathcal{D} : [Phys. Rev. D 94, 114509: C. Alexandrou, S. Bacchio, J. Finkenrath, A. Frommer, K. Kahl, and M. Rottmann]
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Multigrid is needed for Staggered fermions!

Staggered State of the Art

Given the staggered \not{D} operator:

$$\begin{aligned} D_{xy} &= \sum_{\mu} \eta_{\mu}(x) \left[U_{\mu}^{\dagger}(x) \delta_{y+\mu,x} - U_{\mu}(x) \delta_{x+\mu,y} \right] + m \delta_{x,y} \\ &= iA + m\mathbf{I} \quad \leftarrow \text{anti-Hermitian} + \text{Hermitian piece} \end{aligned}$$

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The E/O preconditioned operator is Hermitian positive definite:
Rich theory of multigrid exists.

$$\left(m^2 - D_{eo} D_{oe} \right) \psi_e = m b_e - D_{eo} b_o; \quad \psi_o = \frac{1}{m} (b_o - D_{oe} \psi_e)$$

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Issues:

- Two link sparsity pattern is computationally inefficient.
- Further issue: Even and odd decouple in chiral limit.

Testing Environment for the Interacting Case

- Use a simpler model *with similar physics* for initial testing:
- Two-flavor Schwinger model in two dimensions:

$$\mathcal{L} = \frac{1}{2}F^2 - i \sum_{f=1,2} \bar{\psi}_f \gamma^\mu (\partial_\mu - igA_\mu) \psi_f + m \sum_{f=1,2} \bar{\psi}_f \psi_f$$

- ▶ Confinement
 - ▶ Chiral symmetry breaking
 - ▶ Vortices (2 dimensional “instantons”)
 - ▶ Topology
 - ▶ Two flavor theory has a “pion”-like state: $M_{gap}(m) = A_{gap} m^{2/3} g^{1/3}$
[Phys.Rev. D55 (1997)]
- Comparatively very inexpensive to look at large volumes: 256^2
 - Constant physics and box size: double L , half g (quadruple β), half m .

Standard story...

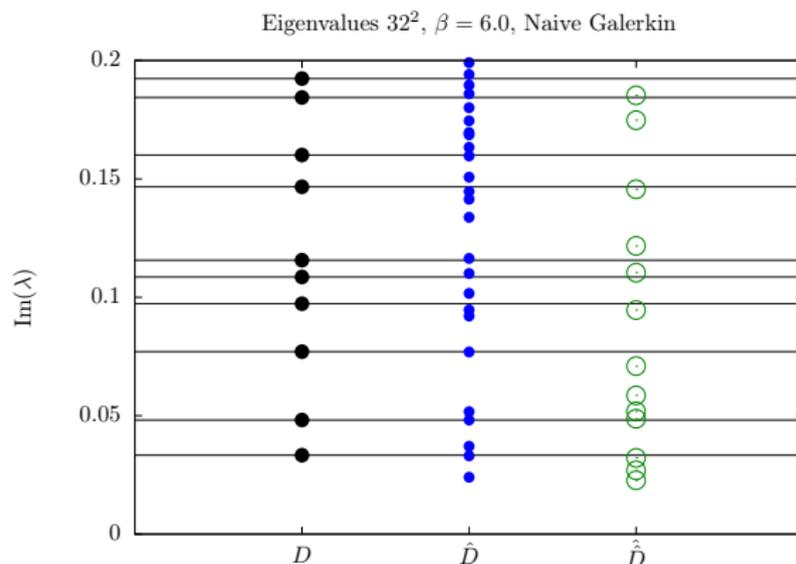
- Generate null vectors (I use BiCGstab-8 to a tolerance of 5×10^{-5} or maximum of 500 mat-vec).
 - ▶ For this talk: 8 null vectors.
- Block orthonormalize, Galerkin projection.
- Preserve chirality $(1 \pm \gamma_5) \rightarrow$ even/odd for staggered.
 - ▶ For this talk: 16 degrees of freedom per coarse site.

Galerkin Projection: Failure

Naïvely using what worked for Wilson fails for staggered.

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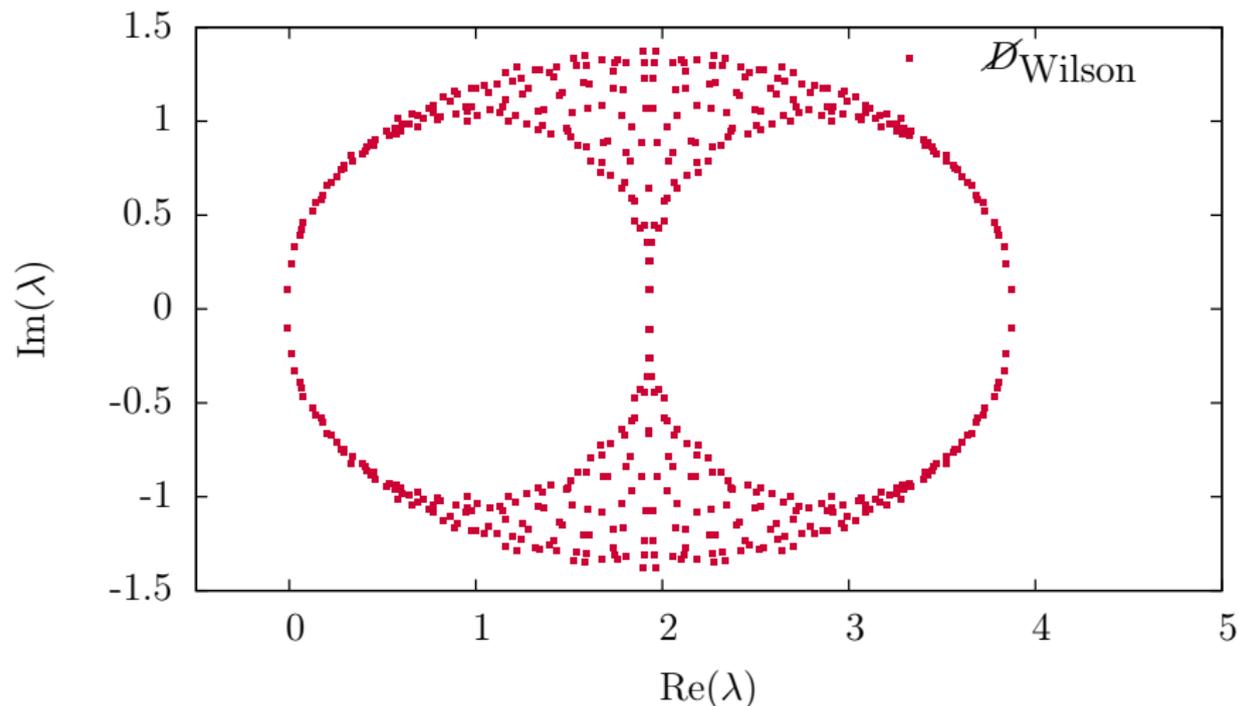
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Spurious low modes.

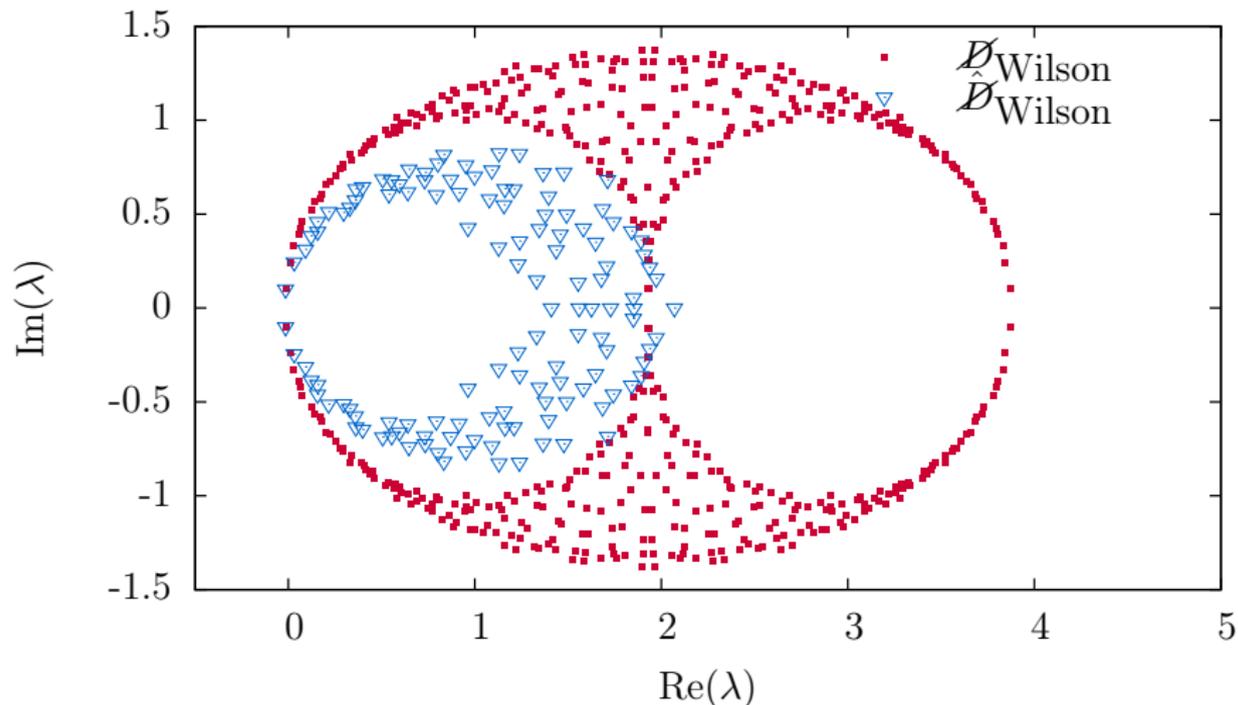
Why (?) Wilson Works

$32^2, \beta = 6.0, m = -0.07$

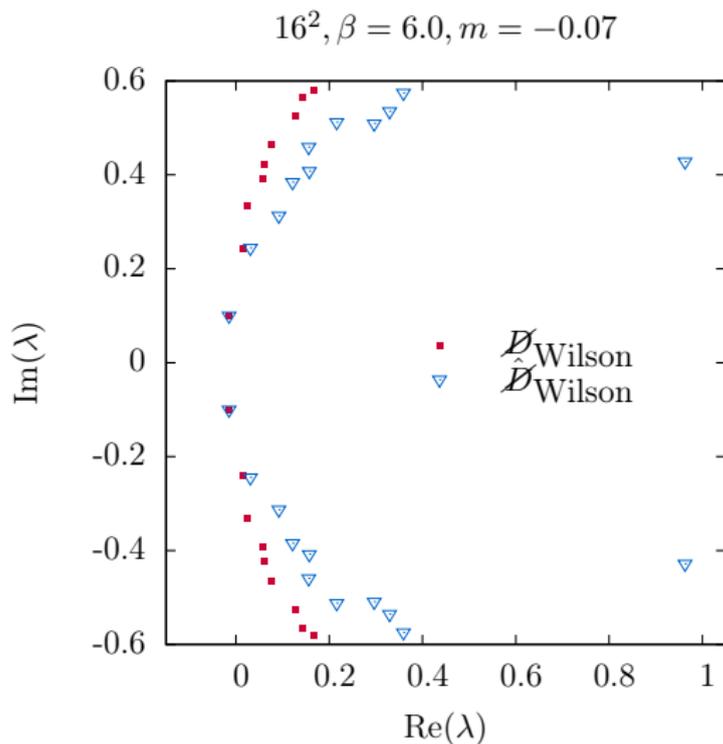


Why (?) Wilson Works

$16^2, \beta = 6.0, m = -0.07$



Why (?) Wilson Works



Kahler-Dirac Thinking

Staggered fermions distribute d fermions over 2^d sites.
Think of this as one “super-site” (or Kahler-Dirac block)

[arXiv:0509026: S. Durr]

$$\mathcal{S} = b^4 \sum_{X,\mu} \bar{q}(X) \left[\nabla_\mu (\gamma_\mu \otimes 1) - \frac{b}{2} \Delta_\mu (\gamma_5 \otimes \tau_\mu \tau_5) + m (1 \otimes 1) \right] q(X)$$

$$\equiv b^4 \sum_{X,\mu} \bar{q}(X) [\not{D} + m] q(X)$$

$$(\nabla_\mu q)(X) = \frac{q(X + b\hat{\mu}) - q(X - b\hat{\mu})}{2b}$$

$$(\Delta_\mu q)(X) = \frac{q(X + b\hat{\mu}) - 2q(X) + q(X - b\hat{\mu})}{b^2}$$

The free staggered operator can be transformed into this form via a
unitary transformation.

Super-site preconditioning

Unitary: spectrum is unchanged.

What if we solve the Kahler-Dirac block exactly?

$$B(m) = \begin{pmatrix} m & 0 & -\frac{1}{2} & -\frac{1}{2} \\ 0 & m & -\frac{1}{2} & \frac{1}{2} \\ \frac{1}{2} & \frac{1}{2} & m & 0 \\ \frac{1}{2} & -\frac{1}{2} & 0 & m \end{pmatrix}$$

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- Wilson-Clover: preconditioning with the inverse clover term.
- Domain Wall: 4D preconditioned operator.

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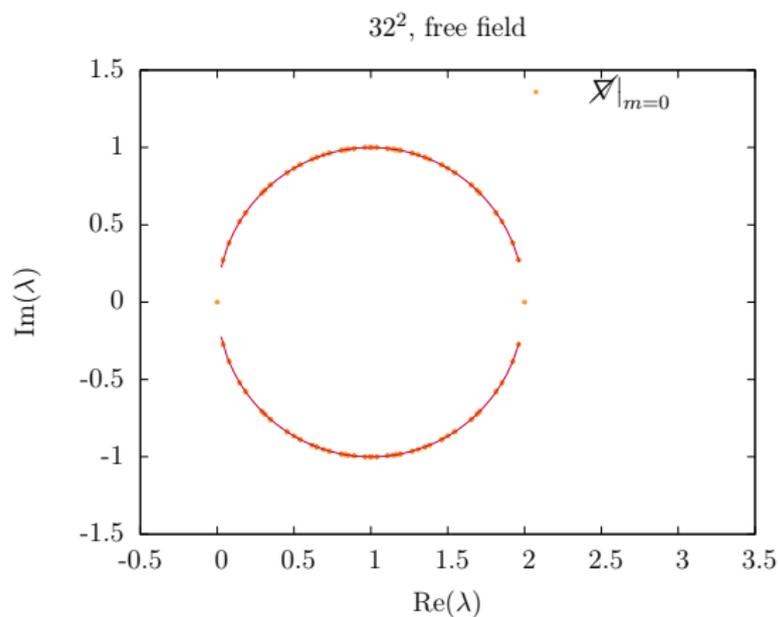
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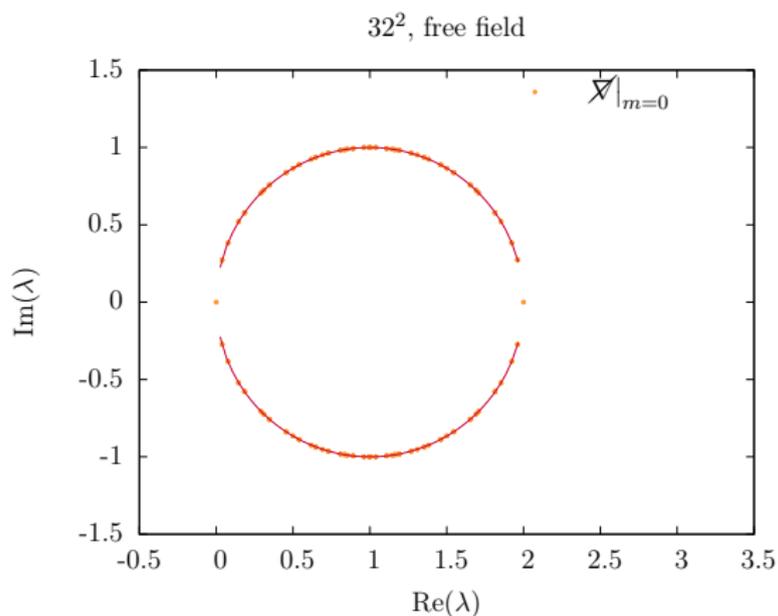
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$$\begin{aligned} \not{D}|_m &= [\not{D} + m] B(m)^{-1} \\ &= \left[\nabla_\mu (\gamma_\mu \otimes 1) - \frac{b}{2} \Delta_\mu (\gamma_5 \otimes \tau_\mu \tau_5) + m(1 \otimes 1) \right] \begin{pmatrix} m & 0 & -\frac{1}{2} & -\frac{1}{2} \\ 0 & m & -\frac{1}{2} & \frac{1}{2} \\ \frac{1}{2} & \frac{1}{2} & m & 0 \\ \frac{1}{2} & -\frac{1}{2} & 0 & m \end{pmatrix}^{-1} \end{aligned}$$

Preconditioned Op

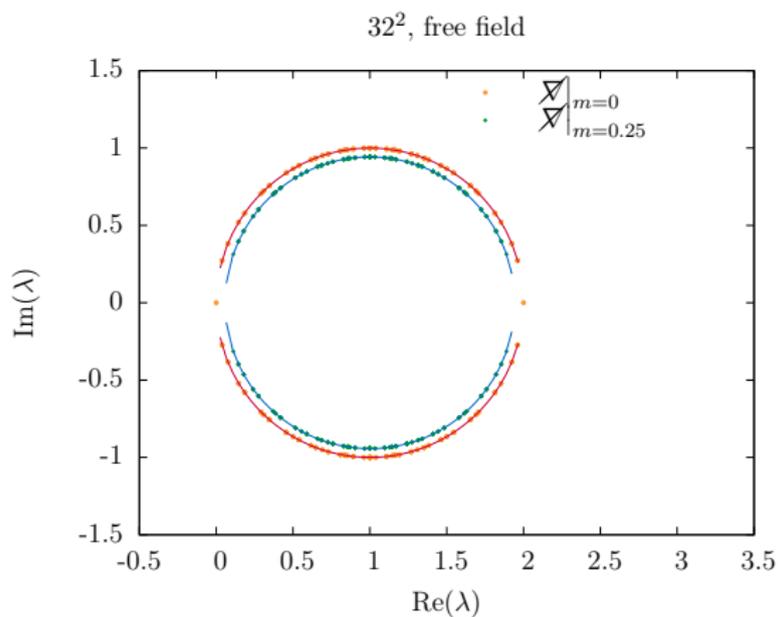


Preconditioned Op



$1 + U$, where U is a *unitary operator*. **This is an Overlap operator.**

Preconditioned Op



Massive case: $1 + \frac{1}{\sqrt{1+2m^2}}U$

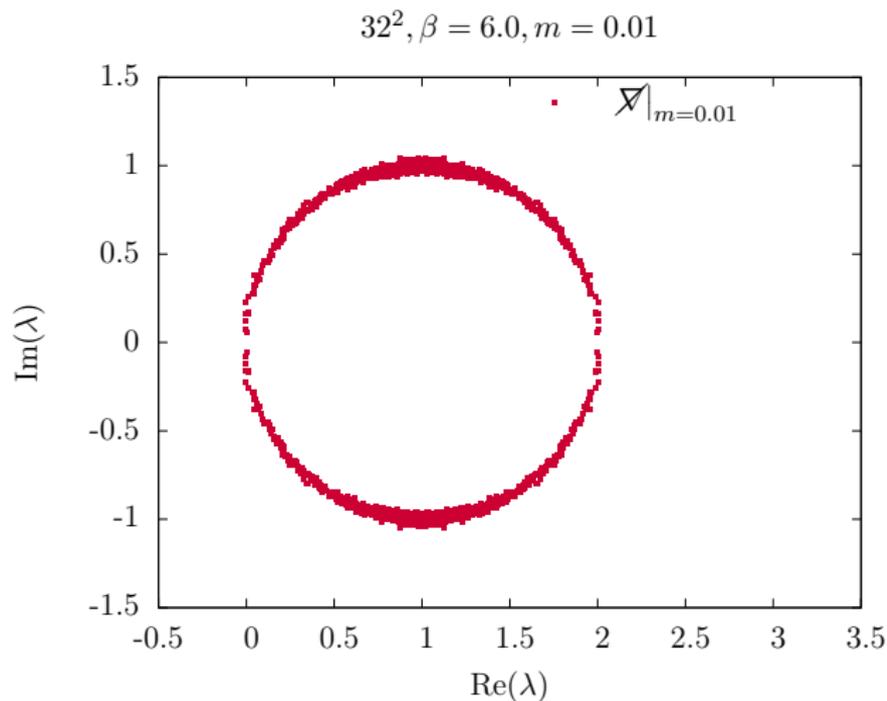
Interacting Case

For the interacting case, we can still perform the unitary transformation then solve the Kahler-Dirac block exactly.

It's now just gauge field \rightarrow coordinate dependent.

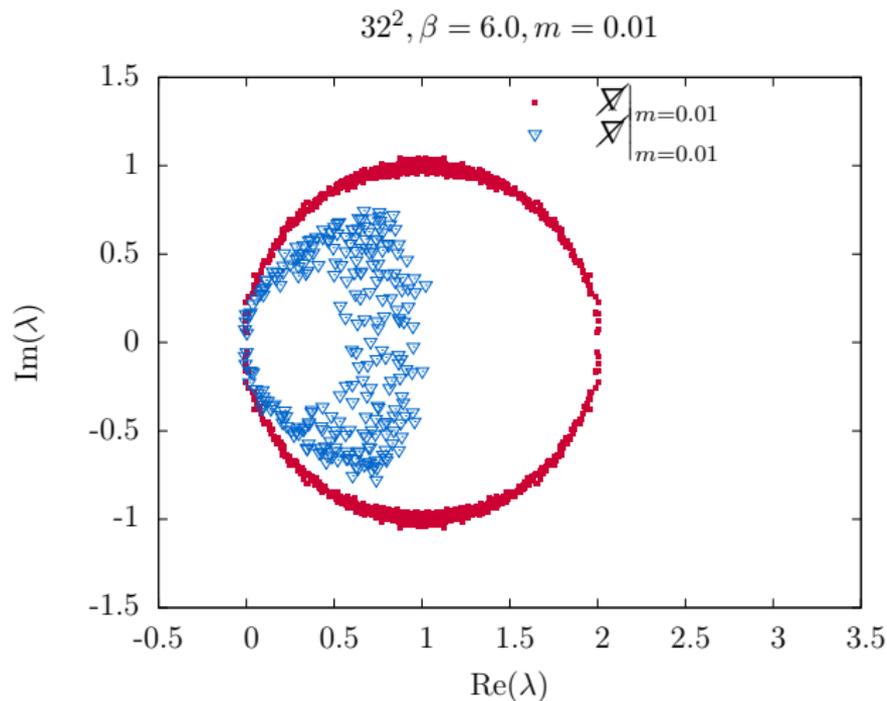
$$\begin{pmatrix} m & 0 & -\frac{1}{2}U_x(2\vec{n}) & -\frac{1}{2}U_y(2\vec{n}) \\ 0 & m & -\frac{1}{2}U_y^\dagger(2\vec{n} + \hat{x}) & \frac{1}{2}U_x^\dagger(2\vec{n} + \hat{y}) \\ \frac{1}{2}U_x^\dagger(2\vec{n}) & \frac{1}{2}U_y(2\vec{n} + \hat{x}) & m & 0 \\ \frac{1}{2}U_y^\dagger(2\vec{n}) & -\frac{1}{2}U_x(2\vec{n} + \hat{y}) & 0 & m \end{pmatrix}$$

Interacting Preconditioned Op



The circle “fuzzes”.

Interacting Preconditioned Op



The circle “fuzzes”... but the spectrum collapses well!

Our Multigrid Prescription

Remark: Our prescription isn't final—this just fixes our tests!

Our Multigrid Prescription

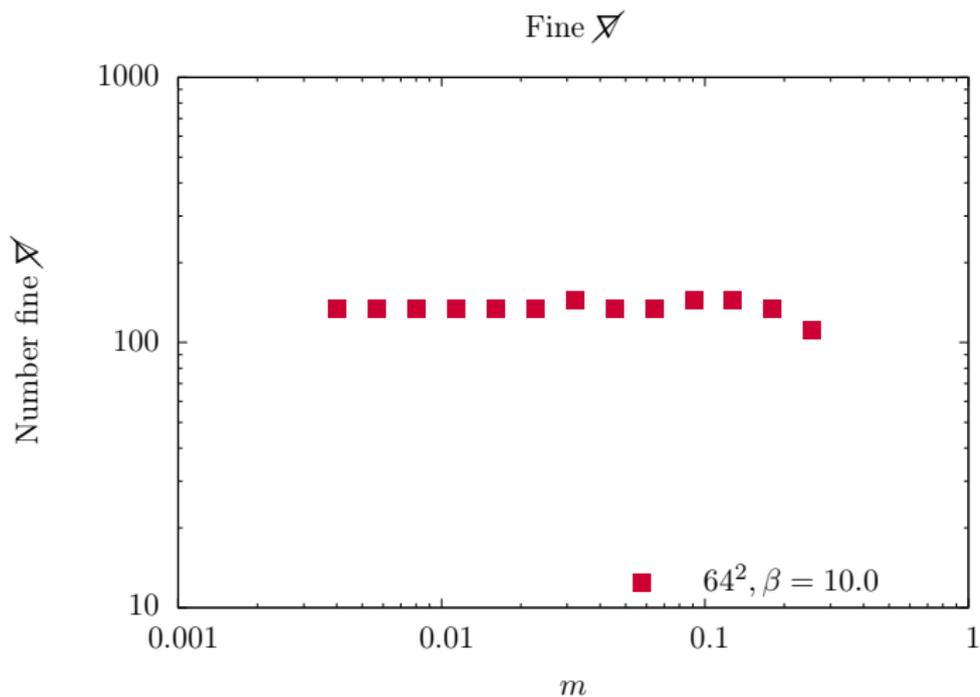
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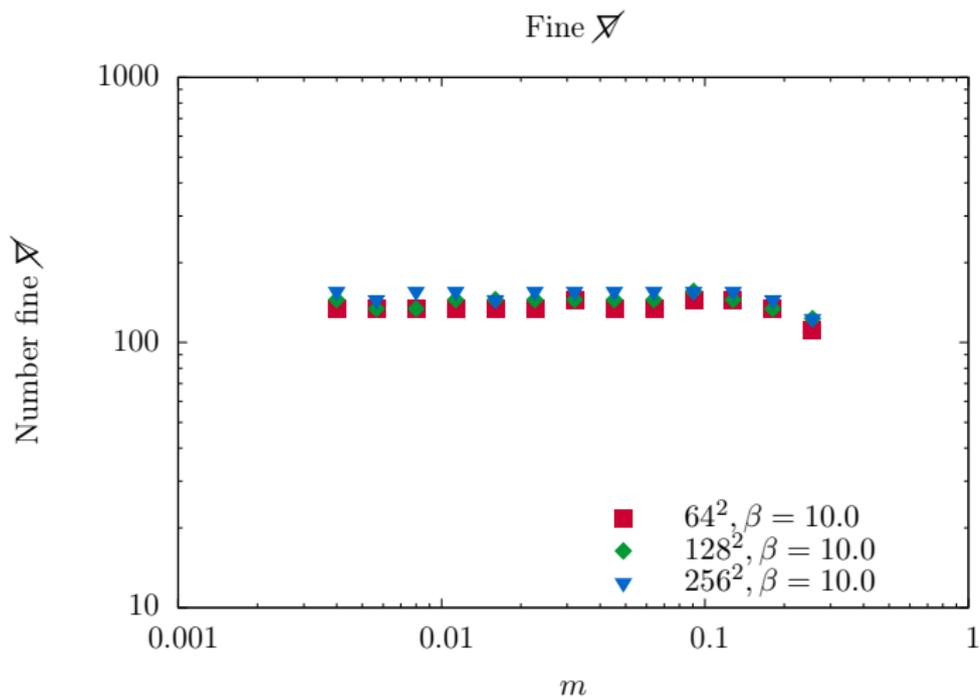
- Null vector generation:
 - ▶ Use BiCGStab-8
 - ▶ Generate random vector \vec{x}_0 .
 - ▶ Solve residual equation $\not{D}\vec{e} = \vec{r}(\equiv -A\vec{x}_0)$ to tolerance 5×10^{-5} , or to a maximum of 500 \not{D} .
 - ▶ Generate 8 null vectors \rightarrow 16 internal dof after chiral projecting

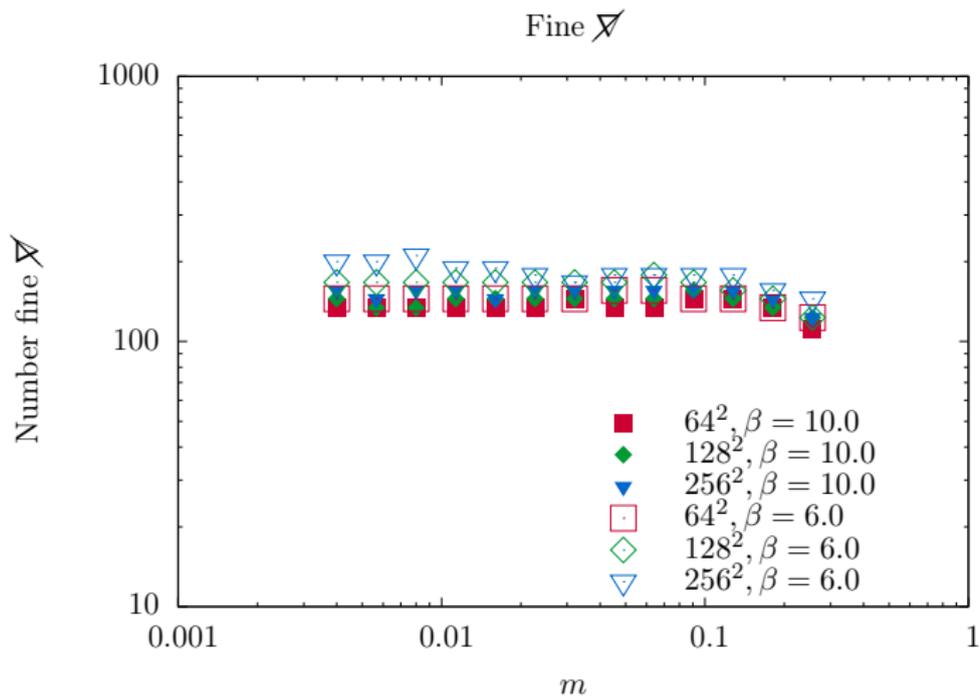
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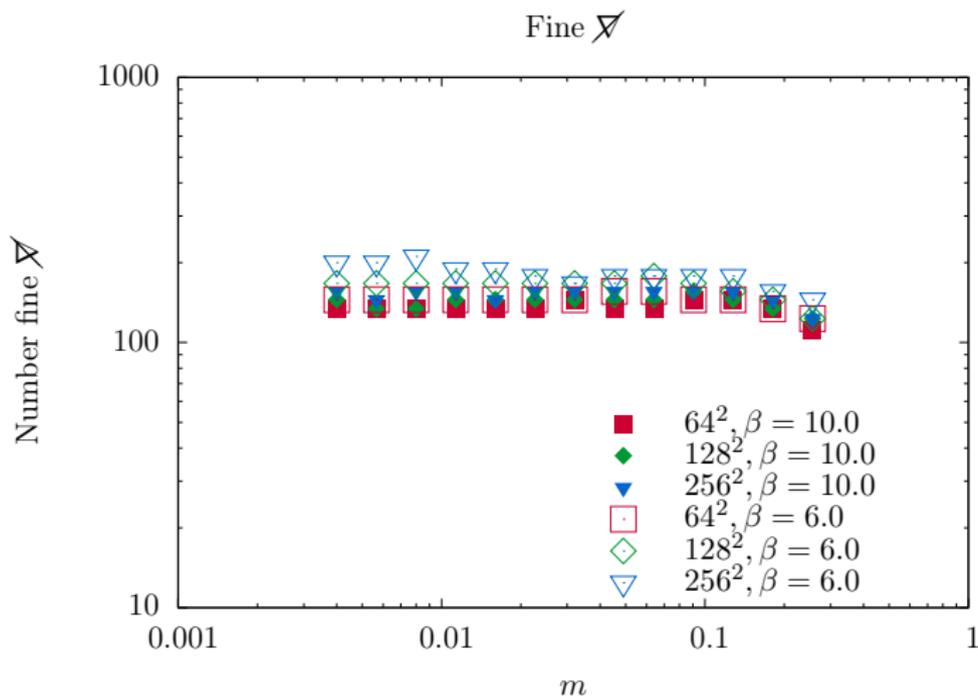
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- MG Solve:
 - ▶ Three level K -cycle
 - ▶ Outer solver: GCR(32) to tolerance 10^{-10}
 - ▶ Block size: 4×4 (e.g., $128^2 \rightarrow 32^2 \rightarrow 8^2$)
 - ▶ Pre- and Post-Smoother: 2 iterations of GCR
 - ▶ Inner solver: GCR(32) to tolerance 0.2



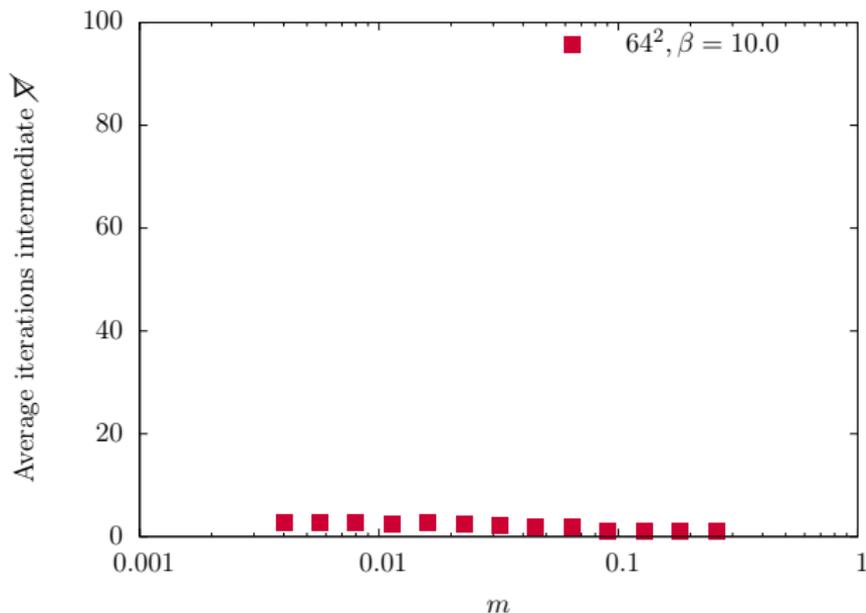




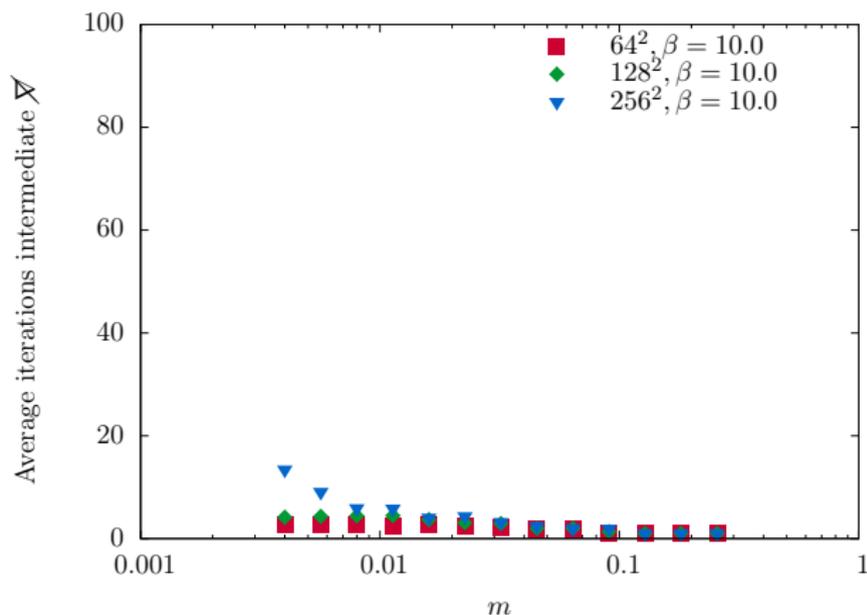


Relatively constant independent of mass, volume, β ...

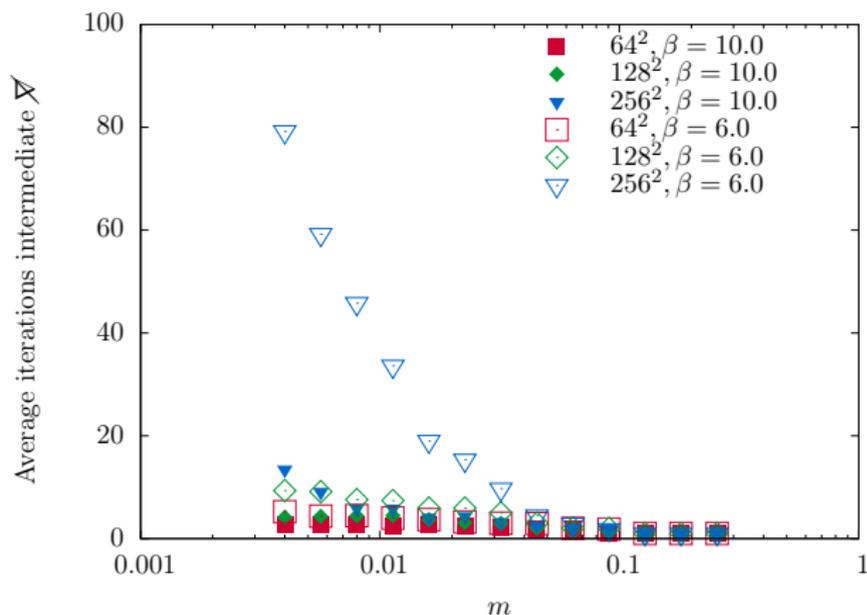
Recursive test: average iterations, intermediate ∇



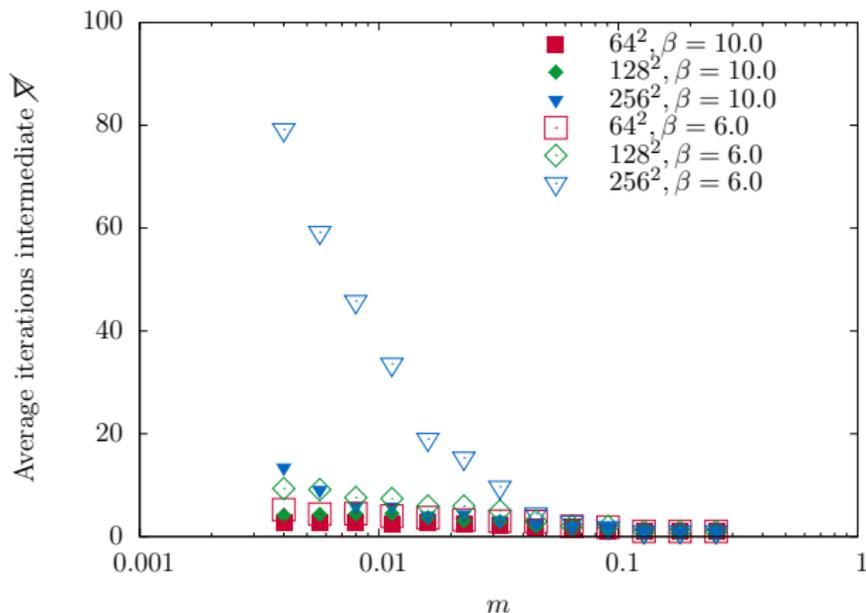
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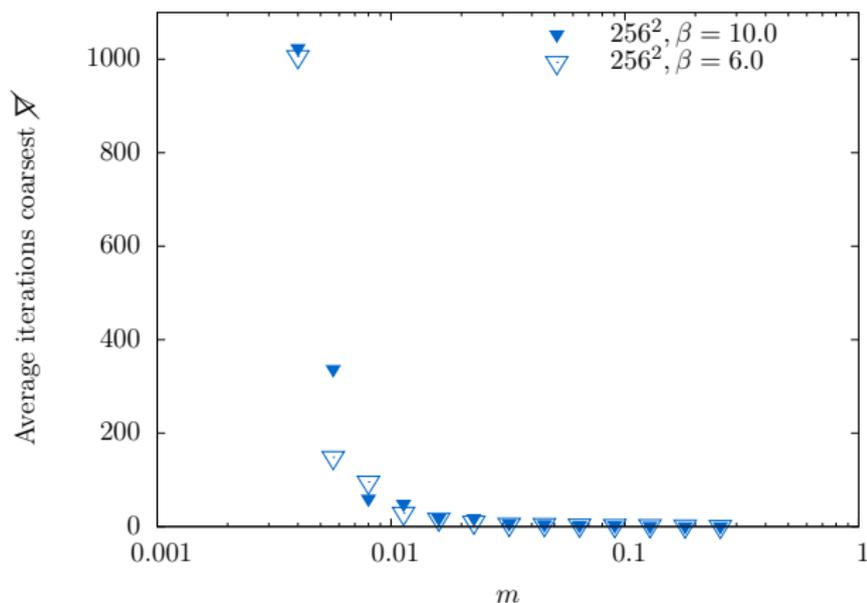


Recursive test: average iterations, intermediate ∇



For larger volumes, smaller β , the coarsest level does *not* precondition the intermediate level well.

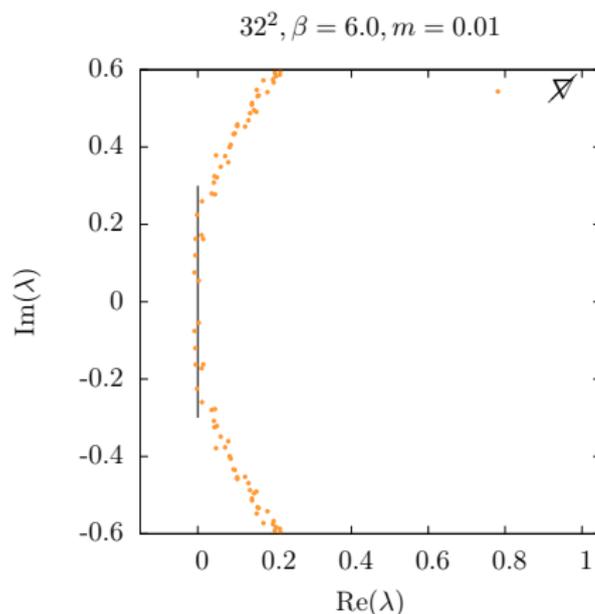
Average iterations, coarsest ∇



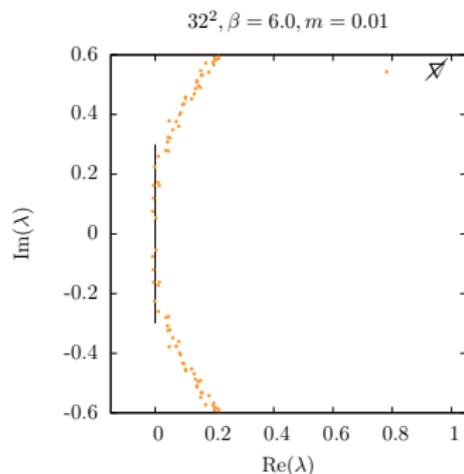
The coarsest ∇ is becoming very ill-conditioned... which is damaging the intermediate solve.

The crux of the issue...

The recursive algorithm is falling apart at larger volume, smaller mass, moreso at smaller β . A likely culprit? The “fuzzed” circle.



Discussion



- Larger volume, the spectrum becomes more dense.
- Smaller mass, the “fuzzed” circle expands.
- Smaller β , the circle becomes more fuzzed.

These all lead to more “exceptional eigenvalues”.

What happens with finer lattices at constant physics?
(double L , quadruple β , half m)

- $64^2, m = 0.01, \beta = 3.0$:
- $128^2, m = 0.005, \beta = 12.0$:
- $256^2, m = 0.0025, \beta = 48.0$:

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- $64^2, m = 0.01, \beta = 3.0$: 200 Fine ∇ , ~ 20 intermediate iterations, ~ 16 coarse iterations.
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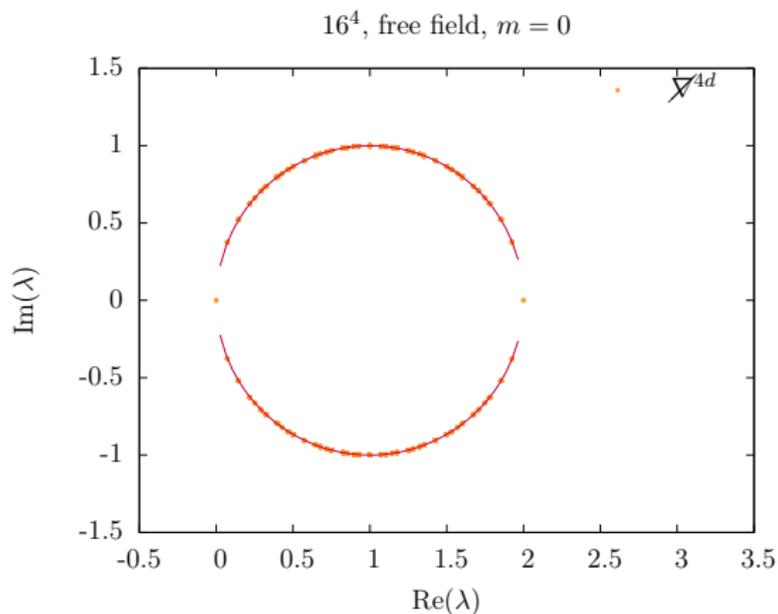
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Even though the coarsest level is becoming ill-conditioned, it becomes a better and better preconditioner.

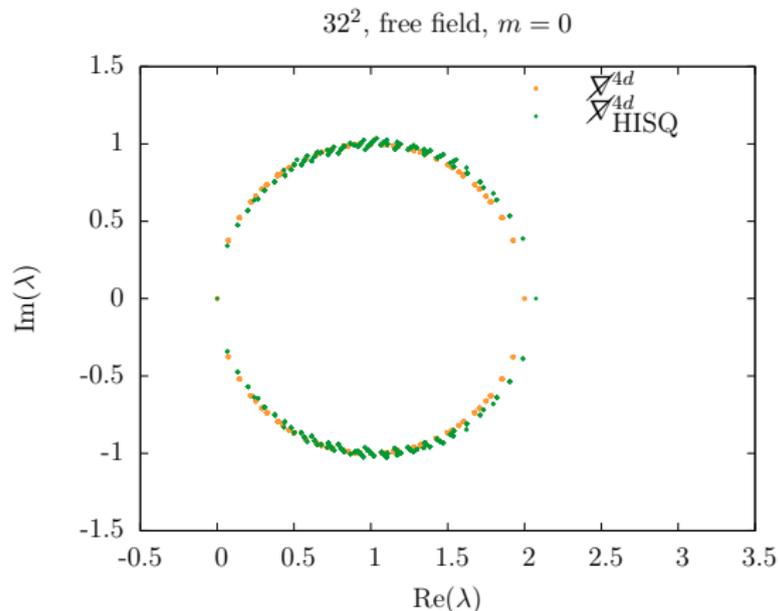
Looking forward: 4D



$$1 - \frac{1}{d} \left(\sum_{\mu} \cos(p_{\mu}) \right) \pm i \sqrt{1 - \frac{1}{d^2} \left(\sum_{\mu} \cos(p_{\mu}) \right)^2}$$

Circle of radius 1

Looking forward: 4D



HISQ term: no longer a circle, but circle-like, stretched to the right.

Overview

- Showed a *recursive, mostly* working MG algorithm for the two-flavor Schwinger model.
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Future Work

- Extend to the *interacting* 4 dimensional staggered operator
- Continue progress on an implementation in QUDA
- Experiment with alternative methods to generate null vectors
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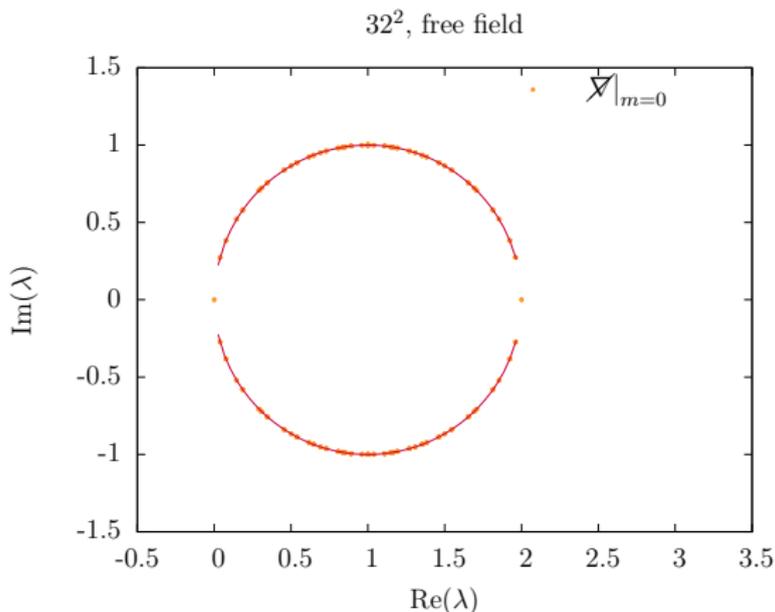
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Thank you!

BACKUP

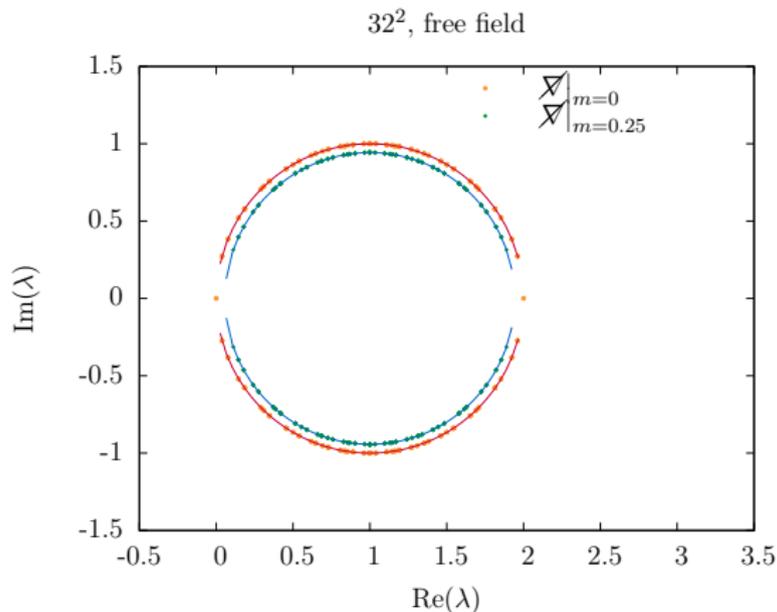
Preconditioned Op



$$1 - \frac{1}{2} \left(\cos(p_x) + \cos(p_y) \right) \pm i \sqrt{1 - \frac{1}{4} \left(\cos(p_x) + \cos(p_y) \right)^2}$$

Circle of radius 1

Preconditioned Op



$$1 - \frac{1}{1+2m^2} \left(\frac{1}{2} (\cos(p_x) + \cos(p_y)) \pm i \sqrt{1 + 2m^2 - \frac{1}{4} (\cos(p_x) + \cos(p_y))^2} \right)$$

Circle of radius $1/\sqrt{1+2m^2}$